



**Evidence of dilatant and non-dilatant damage processes in oolitic iron ore: P-wave velocity and acoustic emission analyses**

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Review

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# Evidence of dilatant and non-dilatant damage processes in oolitic iron ore: P-wave velocity and acoustic emission analyses

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## SUMMARY

Uniaxial and triaxial compression experiments were performed on oolitic iron ores to investigate damage processes. Most of these experiments included four indirect measurements of damage evolution, i.e. P-wave velocity and maximum amplitude received during pulse transmission experiments, elastic properties (apparent Young's modulus and apparent Poisson's ratio) and acoustic emission (AE) monitoring. The mechanical behaviour deduced from strain measurements is dilatant for some samples and non-dilatant for the other samples. However, variations in elastic properties indicate damage processes for all samples. AE source mechanism analysis shows two different microscopic damage processes: 1) for dilatant rock, the development of axial extensive microcracks as well as their interaction and coalescence lead to the formation of shear macroscopic discontinuities; 2) for non-dilatant oolitic iron ore, both compressive and shear micro-mechanisms take place and interact with macro-

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5  
6 19 scopic fractures. A particular consistency between the four types of measurements  
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8 20 employed was observed.

9  
10 **Key words:** damage – acoustic emission – dilatancy  
11

## 12 13 14 1 INTRODUCTION

15  
16 23 Mechanical loading of rocks induces irreversible microstructural changes such as pore crush-  
17  
18 24 ing, microcrack nucleation and growth and grain compaction, which are referred to as damage  
19  
20 25 processes. Such processes, when they occur suddenly, generate transient elastic waves known  
21  
22 26 as Acoustic Emissions (noted AE, Evans 1979; Lockner 1993). Thus, AE constitute a practical  
23  
24 27 tool for studying damage processes and have been widely used during uniaxial and triaxial  
25  
26 28 compression tests of rocks (Scholz 1968b; Lockner et al. 1991; Cox & Meredith 1993; Lock-  
27  
28 29 ner 1993; Zang et al. 1998; Lei et al. 2000; Amitrano 2003). From a macroscopic point of  
29  
30 30 view, damage processes induce physical changes, such as modification of elastic properties  
31  
32 31 (Lemaitre & Chaboche 1990). Damage processes thus induce wave velocity changes. Hence,  
33  
34 32 the measurement of such parameters during the mechanical loading of rocks constitutes an  
35  
36 33 indirect means for estimating damage evolution (Sayers et al. 1990).

37  
38 34 In the case of rocks, in the early stage of deformation, extensive cracks nucleate and prop-  
39  
40 35 agate parallel to the principal stress axis (Jaeger & Cook 1979; Reches & Lockner 1994).  
41  
42 36 This phenomenon is well known and termed dilatancy: the strain measured perpendicularly  
43  
44 37 to principal stress increases faster than the strain measured parallel to the principal stress.  
45  
46 38 Some works have observed comparable anisotropic modifications using wave velocity measure-  
47  
48 39 ments (Ayling et al. 1995; Stanchits et al. 2006). The onset of dilatancy is also related to the  
49  
50 40 appearance of AE sources (Scholz 1968a; Lockner et al. 1991; Lockner 1993). The orientation  
51  
52 41 of cracks can be inferred from AE source mechanism analysis (Zang et al. 1998). Wave veloc-  
53  
54 42 ity has also been used to monitor fracturing processes resulting from damage accumulation  
55  
56 43 (Yukutake 1989; Chow et al. 1995; Rao & Kusunose 1995; Schubnel et al. 2003). In some cases,  
57  
58 44 these measurements have been completed by amplitudes determined from pulse transmission  
59  
60 45 experiments (Rao & Ramana 1992; Zang et al. 2000).

46  
47 46 The stress level corresponding to dilatancy appears to dramatically vary with regard to the  
48  
49 47 microstructure of the rocks and the confining pressure conditions (Rawling et al. 2002; Katz  
50  
51 48 & Reches 2004). For porous rocks such as sandstone under high confining pressure, dilatancy  
52  
53 49 may not appear and may instead be replaced by compaction due to grain crushing (e.g.,  
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55 50 Bésuelle et al. 2000).

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The goal of this study was to acquire a better understanding of iron ore damage processes that have occurred in the context of mine ground collapses. Uniaxial and triaxial compression experiments were performed on samples coming from different strata of an iron ore mine (the Tressange mine) located in Lorraine (eastern France). For studying the damage evolution, an extensive acoustic (i.e. AE monitoring and P-wave velocity measurements) and mechanical instrumentation was used during eight compression experiments. Here we present the results from six experiments including AE monitoring, P-wave velocity and mechanical measurements and two experiments including only mechanical measurements and AE monitoring. The mechanical measurements show a particular pattern indicating dilatant and non-dilatant mechanical behaviours of samples studied. Two experiments performed on dilatant and non-dilatant iron ores are presented here in detail, a particularly extensive sensor array (with 18 transducers) was used to determine damage micromechanisms.

## 2 EXPERIMENTAL PROCEDURE

### 2.1 Rock Description

The ferriferous formation in Lorraine (eastern France), dating from the Aalenian, is made up of a succession of sedimentary beds with varying iron content. Iron ore has been extracted from beds with high iron content (30 to 35 %). These iron ore beds, known by miners as the "Green", "Black", "Brown", "Grey", "Main Yellow", "Savage Yellow" and "Red" beds, from the base (deepest bed) to the top of the deposit, are separated by intercalated marl beds. An iron ore bed is composed of a sequence of three lithological formations materialized in a littoral environment (Bubenicek 1961; Teyssen 1989). There are, from the base to the top of the bed, an argillaceous and carbonated siltstone, an oolitic iron ore and a coarse shelled limestone. The jointed oolites of the oolitic iron ores that interest us, ovoid or spherical grains with a diameter of less than 1 mm, are made of goethite. The oolitic cortex is constituted of concentric layers most probably formed during biochemical processes (Burkhalter 1995; Konhauser 1998). The oolitic iron ore also contains other elements such as carbonated chips of shell tests, quartz grains and pieces of phosphatic bones (Bubenicek 1961). Centimetric black marl inclusions are also present, especially in the oolitic iron ore of the brown-iron layer. The iron ore grains (oolites and others) are cemented by a mixture of siderite, calcite and phyllosilicates in various proportions (Grgic 2001). One of the phyllosilicates, very ferriferous and green coloured, known as chamosite (Bubenicek 1970), is localized around grains or fills the entire space between oolites.

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5  
6 83 Table 1 presents the eight compression tests performed. Two samples came from a sub-vertical  
7 84 borehole (GM0203) drilled in intercalated marls with sub-horizontal bedding between the grey-  
8  
9 85 iron layer and the brown-iron layer. The six iron ore samples came from two sub-horizontal  
10 86 boreholes drilled in oolitic iron ores. More precisely, these samples came from the G mine  
11 87 pillars of the grey-iron layer (GGP borehole) and from the brown-iron layer (BGP borehole)  
12 88 in the Tressange mine. The bedding is sub-horizontal at the GGP borehole location. In the  
13  
14 89 case of the BGP borehole location, a substantial reworking has disturbed the original sub-  
15  
16 90 horizontal bedding. The oolitic iron ores tested here are very porous rocks with a porosity  
17  
18 91 of about 30 % for both the grey-iron and brown-iron layers. The intercalated marls have a  
19  
20 92 porosity of about 15 %.

## 23 2.2 Experimental Setup

### 24 2.2.1 Sample Instrumentation

25  
26 94 Samples were cut into cylinders of 140 mm in length (L) by 70 mm in diameter (D). Particular  
27  
28 95 attention was paid to the parallelism of the ends of specimens. In order to measure the local  
29  
30 96 strains in the central part of the samples, six strain gauges of 20 mm in length, three in the axial  
31  
32 97 direction and three in the transversal direction, were placed on the sample surface. To ensure  
33  
34 98 efficient coupling of the strain gauges, a resin was used to cover the sample surface. Eighteen  
35  
36 99 piezoelectric transducers of 8 mm in diameter (Nano 30) from Physical Acoustic Corporation  
37  
38 100 (PAC) were coupled to the surface of the sample (Figure 1). Phenyl salicylate crystals glued  
39  
40 101 transducers on the lateral surface of the sample; silicon grease was used for the transducers on  
41  
42 102 both ends of the specimen. The spectral band-pass of the transducers ranges from 100 kHz to 1  
43  
44 103 MHz with a major resonant frequency at 300 kHz. Their piezoelectric ceramic component can  
45  
46 104 vibrate only in a longitudinal direction. Three transducers worked as transmitters and were  
47  
48 105 dedicated to P-wave velocity measurements. Fifteen therefore operated as receivers of acoustic  
49  
50 106 signals and constituted a mini-seismic network. Once the sample was fully instrumented, it  
51  
52 107 was placed between two stainless steel end-pieces, machined especially to accommodate two  
53  
54 108 transducers. The upper end-piece included a spherical seat to ensure a vertical position of the  
55  
56 109 whole system constituted by the rock sample and the end-pieces that was placed between the  
57  
58 110 machine platens. In order to reduce the end-effects of the contact between the rock specimen  
59  
60 111 and the machine platens, the sample was placed between steel end-pieces with a cross-section  
112  
113 112 equal to D. The ratio "length (L) over diameter (D)" of specimens was set to a value of 2 to  
114  
115 113 restrict the stress heterogeneity domain and to obtain stress homogeneity in the central part  
114  
115 114 of the sample (Brown & Brady 1985).

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### 116 2.2.2 Compression Machine

117 The axial load was applied by a servo-controlled loading frame. The 1000 kN capacity machine  
118 used was enslaved by a computer controlled hydraulic pump. A pressure sensor placed between  
119 the pump and the press piston allowed the load during the test to be measured. Two LVDT  
120 (Linear Variable Differential Transformers) located between the machine platens (Figure 2)  
121 measured the axial displacements during sample deformation.

### 122 2.2.3 Data Acquisition Systems

123 The mechanical data acquisition system comprised two distinct units. One was dedicated to  
124 strain measurements while the other controlled the machine via the valve of the hydraulic  
125 pump and registered the load and the displacements. All the mechanical measurements were  
126 recorded every 6 seconds.

127 The acoustic data acquisition system consisted of two chains (Figure 2). The first included  
128 the receivers connected to 40 dB pre-amplifiers (PAC) with a 50 kHz-1.8 MHz spectral band-  
129 pass and the analogical/digital card (Engineering Seismology Group Canada Inc., Hyperion  
130 system), which digitises the signals after pre-amplification at a sample frequency of 10 MHz  
131 and with 14-bit vertical resolution with an amplitude interval of  $\pm 2.5$  V for each channel.  
132 This chain also performed AE counting in a continuous manner. The acoustic data (AE  
133 event waveforms) were stored on the hard drive of a computer during the experiment when  
134 amplitude reached a value larger than a given threshold. The maximum frequency recording  
135 was about 5 events per second. After the determination of the noise maximum amplitude,  
136 the signal trigger was set to 50 mV. When signals of 7 channels of the sensor array exceeded  
137 the former threshold, the system triggered and the signal was recorded for a fixed duration  
138 of 409.6  $\mu$ s. The second acoustic chain enabled P-wave propagation velocity measurements. It  
139 consisted of a pulse generator (PAC) connected to the three transmitters via a switching box.  
140 The pulse generator generated a negative pulse with a maximum amplitude of about -360 V.  
141 During this type of measurement, AE monitoring was not performed.

## 142 2.3 Test Procedure

143 Uniaxial tests were carried out at a  $10^{-5} \text{ s}^{-1}$  constant longitudinal strain rate. The loading  
144 consisted in successive loading-relaxation-unloading cycles, with increasing axial load until  
145 failure (Figure 3). The relaxation phase allows the release of delayed elastic strain and there-  
146 fore an appropriate determination of the elastic modulus during the unloading phase with

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147 reduced hysteresis.

148 P-wave propagation velocity measurements were performed in the axial direction (ray T1R15,  
149 Figure 1) and in two transversal directions (rays T2R2 and T3R10, Figure 1), before every  
150 axial load cycle, and then at two points of each cycle (Figure 3), at the end of each relaxation  
151 phase (points  $a_i$ ,  $i = 1$  to  $N_{cycles}$ ,  $N_{cycles}$  is equal to the cycle number) and during the fol-  
152 lowing loading phase at the same stress level (points  $b_i$ ). The recording of the AE activity was  
153 performed during loading and cycles, unless while P-wave measurements were being carried  
154 out.

## 155 **2.4 Data Analysis Methods**

### 156 *2.4.1 Mechanical Analysis*

157 Longitudinal ( $\epsilon_L$ ), transversal ( $\epsilon_T$ ) and volumetric ( $\epsilon_V = \epsilon_L + 2\epsilon_T$ ) strains were plotted versus  
158 axial stress ( $\sigma_1$ ). The displacements measured between the machine platens were also used to  
159 monitor the behaviour associated with the macro-failure. Indeed, strain gauges cannot record  
160 the strains during this phase due to macroscopic fractures destroying the strain gauges. The  
161 apparent Young's modulus ( $E_{app}$ ) and the apparent Poisson's ratio ( $\nu_{app}$ ) have been estimated  
162 by linear regressions performed at each unloading phase of the cycles.

$$163 \quad E_{app} = \frac{\sigma_1}{\epsilon_L} \quad (1)$$

$$164 \quad \nu_{app} = -\frac{\epsilon_T}{\epsilon_L} \quad (2)$$

165 Where  $\sigma > 0$  is compression and  $\epsilon > 0$  denotes shortening.

### 166 *2.4.2 Acoustic Analysis*

167 The signals recorded during the pulsing phase were analysed in order to determine the P-wave  
168 velocity. The arrival time ( $t_P$ ) was manually determined (so-called picking) from the signal  
169 waveforms. Each pulse generates an electromagnetic signal in the system that is present on all  
170 the waveforms recorded during the P-wave velocity measurements. The time of this artefact  
171 allowed identification of the origin time ( $t_0$ ) of the pulse. Thus, P-wave velocity ( $V_p$ ) is deduced  
172 from the time-of-flight ( $t_P - t_0$ ) of the P-wave and the distance between the receiver R and  
173 the transmitter T ( $d_{RT}$ ).

$$174 \quad V_p = \frac{d_{RT}}{t_P - t_0} \quad (3)$$

175 The greatest uncertainty in  $V_p$  determination depends on the accuracy in picking the arrival  
176 time of the P-wave. This is estimated to be less than 2 %.



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177 Maximum amplitude of the signal received during the P-wave velocity measurements allowed  
178 us to monitor evolution of attenuation of the medium during compression tests.

179 Concerning AE event signals, the first arrival time and the end of the signal were automatically  
180 picked using a Short Time Average/Long Time Average algorithm (Allen 1978). Thus, an  
181 event is recognized between the first arrival time and the end picking or end of the waveform.  
182 The energy of AE event signals was determined between such pickings as follows:

$$183 \quad E = \frac{1}{N_k} \sum_{k=1}^{n_k} \left( \sum_j A_{jk}^2 \cdot \Delta t \right) \quad (4)$$

184 where E is the signal energy of one AE event (in  $V^2 \cdot s$ ),  $N_k$  the quantity of channels where  
185 the event has been identified,  $A_j$  the amplitude values at each sampling time between P-wave  
186 pick and end-pick in Volts (V), and  $\Delta t$  the sampling pitch. We considered that the signal  
187 energy received at the sensors was proportional to the source energy released (Evans 1979;  
188 Lockner 1993). The transducer polarisation was checked beforehand. It is negatively polarised  
189 (i.e. negative oscillation corresponds to compressional first motion). Then, the polarity of the  
190 first pulse amplitude, which is right after the P-wave pick, was studied for each waveform  
191 in order to assess the first motion characteristics of each event and thus AE event source  
192 mechanisms (Sato et al. 1990; Lei et al. 1992; Zang et al. 1998). If the majority of channels  
193 have compressional (see Figure 4) or dilatational first motions, the events can result from  
194 tensile sources or compressive sources respectively. In other instances, the sources can be  
195 associated with a shear mechanism. Signal to noise ratios (SNR) were used to check the  
196 accuracy in picking the P-wave arrival time and thus the first motion determination, and  
197 were determined as follows as a function of the middle noise:

$$198 \quad SNR = A_{first} \cdot \left( \frac{1}{N} \cdot \sum_N A_{noise} \right)^{-1} \quad (5)$$

199 where  $A_{first}$  is the maximum amplitude of the first pulse right after P-pick,  $A_{noise}$  the ampli-  
200 tude of the signal before P-wave pick and N the number of amplitude points before P-wave  
201 pick. If the SNR is less than 2.5, the channel is not taken into account in the compressive first  
202 motion ratio calculation.

### 203 3 EXPERIMENTAL RESULTS

#### 204 3.1 Mechanical Behaviour

205 Figure 5 shows the mechanical behaviours of samples coming from the grey-iron layer (sam-  
206 ple GGP1-A2, see Table 1), the brown-iron layer (BGP2) and from different facies of the



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6 207 intercalated marl layer (samples GM0203S1 and GM0203S2). Dilatancy is clearly observed  
7 208 on the volumetric strain curves of GGP1-A2 and GM0203S1 (Figures 5a and 5b), the rela-  
8  
9 209 tive volumetric strain becomes negative as compression goes on. For BGP2 and GM0203S2  
10 210 samples (Figures 5c and 5d), the volumetric strain curves do not show any dilatancy, the  
11 211 relative volumetric strain remains positive with compression. The mechanical behaviours of  
12 212 uniaxially compressed GGP and BGP samples including extensive instrumentation (Table 1  
13  
14 213 and Figure 1) are detailed below. Figures 6a and 6b illustrate the mechanical behaviour of the  
15  
16 214 GGP sample. In 6b, six loading-relaxation-unloading cycles were performed in the linear part  
17  
18 215 of the differential stress-displacement curve. Longitudinal, transversal and volumetric strains  
19 216 were plotted versus differential stress (axial stress minus confining pressure) in Figure 6a.  
20  
21 217 Four main phases can be identified (Jaeger & Cook 1979). Firstly, the phase where the curve  
22 218 is upwardly concave, which is attributed to the closure of pre-existent cracks; secondly, the  
23  
24 219 linear part of the curve related to pure elasticity; thirdly, the behaviour becomes non-linear,  
25 220 the curve reaches a maximum ( $\sigma_M$ , about 31 MPa) corresponding to the initiation of the  
26  
27 221 macro-failure that grew and sheared during the fourth phase. Several stress thresholds were  
28 222 graphically identified (Table 2):  $\sigma_{CC}$  determined on the stress-displacement curve, is the initial  
29 223 crack closure threshold. The other thresholds came from stress-strain curves.  $\sigma_C$ ,  $\sigma_{LL}$  and  $\sigma_D$   
30 224 are respectively the cracking initiation determined on the transversal stress-strain curve, the  
31 225 loss of linearity of the longitudinal stress-strain curve and the dilatancy threshold observed  
32 226 on the volumetric strain curve. The peak stress  $\sigma_M$  is the maximum stress determined on the  
33  
34 227 stress-displacement curves. The appearance of non-linear behaviour on the lateral strain curve  
35 228 from the threshold  $\sigma_C$  is attributed to the initiation of the cracking, while the  $\sigma_{LL}$  threshold  
36  
37 229 corresponds to the linearity loss observed on longitudinal behaviour; it is attributed to the  
38 230 coalescence of microcracks. The dilatancy appears at the  $\sigma_D$  stress threshold about 24 MPa,  
39 231 volumetric strain increases indeed clearly from  $\sigma_D$ . The cycles present linear and reversible  
40 232 unloading and re-loading phases, with a very slight hysteresis. It can be seen that each unload-  
41 233 ing phase was performed after the relaxation phase allowing, thereby the release of delayed  
42  
43 234 elastic strain. Unloading phases were also performed until the half of the preceding maximum  
44 235 stress. Non-linearity could be seen if the unloading phases would be performed until a null  
45  
46 236 stress. It may be observed slightly during the first and second cycles of BGP compression  
47 237 (Figure 5d). Linear regressions were performed along each unloading phase for all cycles in  
48 238 order to determine the apparent Young's modulus ( $E_{app}$ ) and apparent Poisson's ratio ( $\nu_{app}$ ).  
49 239 They are plotted versus stress in Figures 7a and 7c. The stress range from which the linear  
50 240 regression was performed is indicated for each modulus value by horizontal bars. Apparent

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241 Young's modulus variations are lesser than the error estimate and can be considered as constant, whereas the apparent Poisson's ratio clearly increases with stress.

242

243 The mechanical behaviour of the BGP sample is shown in Figures 6c and 6d. Stress is plotted

244 as a function of transversal, longitudinal and volumetric strains on the one hand and as a

245 function of displacements on the other hand. The BGP mechanical behaviour observed in

246 Figure 6 illustrates the four phases described for the GGP sample. However, for BGP the

247 fourth phase related to macro-failure growth was not fully observed, since the loss of sample

248 strength occurred too quickly to be monitored by the machine. The kink of the transversal

249 strain-stress curve near the peak stress (point k, Figure 6c) corresponds to the failure

250 of transversal strain gauges due to the macroscopic fracture growth (Figure 11). The pre-

251 existing crack closure phase was more pronounced for BGP than for GGP (Figure 6) and

252 BGP strength was higher than GGP strength (see Table 2). The  $\sigma_C$  and  $\sigma_{LL}$  thresholds were

253 easily identified for BGP. Dilatancy was not observed for the BGP sample. The seven cycles

254 show linear and reversible unloading and re-loading phases with a very slight hysteresis. As

255 for GGP compression, apparent Young's modulus variations during BGP loading (Figure 7b)

256 are lesser than error estimate and should be considered as constant. Concerning  $\nu_{app}$  (Figure

257 7d), important errors are associated with the measures performed during the three first cycles

258 before the  $\sigma_{CC}$  threshold, so no trend can be discerned. After the  $\sigma_{CC}$  threshold, it increases

259 with stress. Table 3 shows the mean values of the mechanical characteristics for each tested

260 sample.

### 3.2 P Wave Velocities And Attenuation

261

262 The middle values of P-wave velocities determined along the axial ray (T1R15 ray, see Figure

263 1) and transversal rays (T2R2 and T3R10, see Figure 1) for each sample are given in Table 3.

264 The axial velocities appear to be slightly faster than the transversal velocities. The difference

265 between the two transversal velocities of the BGP sample is not significant regarding velocity

266 uncertainty, for GGP it can be attributed to the heterogeneities. The P-wave velocities of the

267 GGP sample are faster than for the BGP sample. The evolution of P-wave velocity of the sam-

268 ples during uniaxial compression and that of maximum amplitude of signals received during

269 these measurements performed at both points  $a_i$  and  $b_i$  (see Figure 3) of each cycle are shown

270 in Figure 8a and 8b for GGP and in Figures 8c and 8d for BGP. The attenuation was estimated

271 through the maximum amplitude of the received signal. Normalized velocity ( $Vp/Vp_{(\sigma=0)}$ )

272 and maximum amplitude ( $Am/Am_{(\sigma=0)}$ ) are plotted versus stress. Stress thresholds are also

273 given on the X-axis. On the whole, whatever velocity or maximum amplitude is considered,

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6 274 the values at the end of the relaxation phase (points  $a_i$ ) before unloading and values at an  
7 275 equivalent stress level during the re-loading phase (points  $b_i$ ) are very close. This indicates  
8  
9 276 that velocity and maximum amplitude do not vary during an unloading-reloading cycle.

10 277 As regards the GGP sample, the axial velocity measured along the ray T1R15 seems to be  
11  
12 278 constant during the test, while transversal velocities (along T2R2 and T3R10 rays) decrease  
13  
14 279 after cracking initiation ( $\sigma_C$  threshold) has been reached. Maximum amplitude variations are  
15  
16 280 a little different, especially in the axial direction, where maximum amplitude increases until  
17  
18 281 the  $\sigma_C$  threshold is reached and then becomes constant. In transversal directions, the maxi-  
19  
20 282 mum amplitude evolution is similar to those of transversal velocities.

21 283 Figures 8c and 8d show the evolution of normalized velocity and maximum amplitude for the  
22  
23 284 BGP sample. Axial velocity increases slightly until cracking initiation ( $\sigma_C$ ), while transversal  
24  
25 285 velocities appear roughly constant during stress build-up. Concerning maximum amplitude,  
26  
27 286 whatever direction is considered, it seems to decrease after stress reached the  $\sigma_{LL}$  threshold.

### 28 287 **3.3 Acoustic Emissions**

#### 29 288 *3.3.1 Acoustic Activity*

30  
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32  
33 289 The number of AE events identified for each sample compression test is given in Table 3.  
34  
35 290 A set of 112 AE events were identified (P-wave picking performed) for the GGP sample  
36  
37 291 and 160 for the BGP sample. Figures 9a and 9c superimpose AE activity (cumulative AE  
38  
39 292 events) plotted versus displacements and stress-displacement curve respectively for the GGP  
40  
41 293 and BGP samples. The cumulative energy of AE events during compression tests are also  
42  
43 294 plotted in Figures 9b and 9d. It should be observed that, as for the GGP or BGP samples,  
44  
45 295 loading-relaxation-unloading cycles are aseismic. AE is not observed unless stress has reached  
46  
47 296 the previous maximum stress value revealing the Kaiser effect (Kaiser 1950; Goodman 1963).  
48  
49 297 Concerning the GGP sample, AE activity is very low before stress has reached  $\sigma_{CC}$ . It increases  
50  
51 298 after that stress level in a monotonous manner until peak stress, except between  $\sigma_{LL}$  and the  
52  
53 299 onset of the non-linear pre-peak phase of the stress-displacement curve where it decreases  
54  
55 300 slightly. During macro-failure growth beyond the strength peak ( $\sigma_M$ ) and during macro-  
56  
57 301 discontinuity shearing, the AE activity increase is stronger and includes the most energetic  
58  
59 302 AE events. With regard to the BGP sample (Figures 9c and 9d), AE activity is higher than  
60  
303 for the GGP sample during the pre-existing crack closure phase (prior to the  $\sigma_{CC}$  threshold);  
304  
305 304 it shows a notable decrease during the linear phase of the overall mechanical behavior (linear  
306  
305 part of the stress-displacement curve). The energy level of these AE events is low. AE activity  
306  
306 increases strongly during the non-linear part of the differential stress-displacement curve then,

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the more stress approaches  $\sigma_M$ , the more it becomes stronger and results in very energetic AE events.

### 3.3.2 First Motion Analysis

Compressive first motion ratios were determined for each AE event. Figure 10a shows its evolution during the GGP uniaxial compression test (see the stress thresholds on the X-axis). The first phase of the mechanical behaviour delimited by the  $\sigma_{CC}$  threshold includes events with a low compressive first motion ratio (less than 30 %) corresponding to compressive mechanisms and events with more than 70 % of compressive first motion related to extensive mechanisms. During the second phase, the linear part of the behaviour between  $\sigma_{CC}$  and  $\sigma_C$  (superimposed on  $\sigma_D$  in this figure) the majority of events includes a high compressive first motion ratio due to pure extensive mechanisms and also, in lesser proportion, a medium one caused by shear mechanisms. After the cracking initiation threshold ( $\sigma_C$  superimposed on  $\sigma_D$ ), a mixed first motion dominated with a medium compressive first motion ratio, hence shear is the most important mechanism.

Concerning the BGP sample, the compressive first motion ratios are plotted in Figure 10b. Compression is associated with the high compressive first motion ratio blank that marked the absence of pure extensive mechanisms. BGP events that indeed include mainly a low (less than 30 %) and medium (between 30 % and 70 %) compressive first motion ratio, are distributed respectively during the first phase of mechanical behaviour (prior to  $\sigma_{CC}$ ) and after cracking initiation (from the  $\sigma_C$  threshold). During the first phase, AE source mechanisms are mostly compressive. After  $\sigma_C$  threshold (superimposed on  $\sigma_{LL}$  threshold in this figure), shear mechanisms are dominant.

## 4 DISCUSSION

The first stage of the mechanical behaviour of GGP (Figures 6a and 6b) i.e. before the  $\sigma_{CC}$  threshold, where the concavity of the curve is upward, expresses a reduction in the axial compressibility of the rock. This behaviour is generally attributed to the closure of narrow pre-existing microcracks (Paterson 1978), the largest dimension of which is sub-orthogonal to the axial stress direction ( $\sigma_1$ ). This hypothesis is reinforced on the one hand by the observation, during this first stage, of an increase in the P-wave velocity ( $V_p$ ) and the maximum amplitude ( $A_m$ ) measured along the axial ray T1R15 (Figures 8a and 8b) and, on the other hand, by purely compressive sources (0 % of compressive first motions, Figure 10a). In the

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338 case of the BGP sample, this first stage is also distinguished (Figures 6c and 6d). The assump-  
339 tion of microcrack closure is confirmed by purely compressive sources (Figure 10b) observed  
340 prior to the  $\sigma_{CC}$  threshold, but P-wave velocity ( $V_p$ ) and maximum amplitude ( $A_m$ ) varia-  
341 tions (Figures 8c and 8d) appear to not be affected by the microcrack closure. Pre-existing  
342 microcrack closure can be considered as a more or less reversible process that occurs only  
343 during the first phase of the mechanical behaviour i.e. when upward concavity of the stress-  
344 displacement curve is observed; according to Brace et al. (1966), pre-existing microcracks are  
345 being closed when the curve becomes linear. It can be seen that the non-linearity of the first  
346 phase of mechanical behaviour is only observed on stress-displacement curves both for GGP  
347 and BGP compressions and not on stress-strain curves. This may be explain by the nature of  
348 the measurements. Displacements were measured between the press platens (Figure 2). The  
349 displacement measurement is a measure of the axial strains that occur in the entirety of the  
350 sample and not only strains that take place in a central part of the sample. Indeed, longitu-  
351 dinal strain gauges are 20 mm long and measure strains along one-seventh of the longitudinal  
352 dimension of the sample. The closure of narrow pre-existing microcracks sub-orthogonal to  
353 the compression direction involves reduced strains that can be only clearly visible on the  
354 stress-displacement curves embracing the all closure processes.

355 During compression of very porous rocks such as oolitic iron ores, we expected pore crushing  
356 (Paterson 1978) that could also produce AE. Such an irreversible process (damage) could oc-  
357 cur at low stress levels, especially if the pores are non-spherical inter-granular voids as in the  
358 rocks studied here. Grain rotations and translations could also occur (Kranz 1983) involving  
359 AE. Pore crushing is a pure compressive mechanism, whereas grain rotations and translations  
360 are related to shear mechanisms. These damage processes could also occur during the first  
361 stage of the mechanical behaviour and after the  $\sigma_{CC}$  threshold.

362 Several loading-relaxation-unloading cycles were performed in the linear part of the displacement-  
363 stress curves, six cycles for GGP and seven during BGP compression. Unloading phases, which  
364 correspond to the elastic response of the material, appear linear and reversible for both tests,  
365 even just after the loss of linearity ( $\sigma_{LL}$ ) of the overall stress-longitudinal strain curves (Fig-  
366 ure 6a and 6c), suggesting elastic deformation; these phases are mostly aseismic (Figures 9a  
367 and 9c). Indeed, we have seen that AE events can result from inelastic processes such as the  
368 creation of microcracks, pore crushing. Furthermore,  $V_p$  and  $A_m$  measured before unloading  
369 (points  $a_i$ , Figure 8) and after re-loading (points  $b_i$ ) are similar. Between cycles, a significant  
370 acoustic activity is recorded (Figures 9a and 9c), showing an inelastic process of deformation.  
371 Moreover, permanent strains are also visible when the unloading phase is linearly extrapo-

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lated to a zero stress value (Figures 6a and 6c). Hence, the mechanical behaviour observed during uniaxial compression of GGP or BGP appears as a superposition of two deformation processes. One results from reversible processes such as elastic deformation of oolites and cement. The second embraces irreversible processes such as the initiation and interaction of microcracks, pore crushing, grain rotations and translations.

Variations in the apparent Poisson's ratios ( $\nu_{app}$ ) during the GGP and BGP compression tests (Figures 7c and 7d) suggest damage processes. Concerning GGP compression, dilatancy observed on the volumetric strain-stress curve (Figure 6a) is attributed to the initiation of microcracks in extension, the largest dimension of which is sub-parallel to  $\sigma_1$  (Brace et al. 1966; Jaeger & Cook 1979). This hypothesis of anisotropic damage characterized by such extensive microcracks is reinforced by the variations in  $V_p$  and  $A_m$  (Figures 8a and 8b). Indeed, the axial velocity along the T1R15 ray is constant whereas the transversal velocities (rays T2R2 and T3R10) decrease strongly after the  $\sigma_C$  threshold, marking the opening of new voids such as microcracks (Goodman 1980). Oolitic iron ores are very porous sedimentary rocks and therefore pore crushing and readjustment of grains (ovoid oolites) should play an important role in irreversible deformation processes. However, as shown in Figure 10a, there is a small number of compressive AE sources registered (0 % of compressive first motions), AE sources before macroscopic manifestation of dilatancy ( $\sigma_D$ ) have mostly pure extensive mechanism (100 % of compressive first motions), suggesting that the dominant damage process is more opening of new voids (microcracks) than compaction. This is mirrored in macroscopic measurements: there is no axial hardening ( $E_{app}$  seems to be constant, Figure 7a) and dilatancy is clearly observed (Figure 6a). Beyond the dilatancy onset, shear mechanisms are mostly registered (Figure 10a, from 30 % to 70 % of compressive first motions), in addition they correspond to very energetic AE events (Figure 9b) involving larger sources. Shear sources could result from the coalescence of axial extensive microcracks observed during the previous mechanical behaviour phase. Such a damage process has been observed in crystalline rocks under compression (e.g., Moore & Lockner 1995; Lei et al. 2000). Coalescence of extensive axial microcracks should form shear discontinuities that grow and lead to macroscopic failure (Figure 11). Reches & Lockner (1994) have modelled the propagation of such a fracture. It forms in a region of high extensive microcrack density. Then it grows through the coalescence of extensive microcracks and propagates with the creation of extensive microcracks at its tip. For BGP compression, the apparent Poisson's ratio increases beyond a stress value of 7 MPa, (Figures 7d), suggesting a damage process that takes place after the  $\sigma_{CC}$  threshold. Most of the AE events registered up to  $\sigma_C$  have a compressive mechanism (Figure 10b, from 0 % to 30



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6 406 % of compressive first motions) suggesting pre-existing microcrack closure and pore crushing.  
7 407 Such processes lead to compaction in a direction sub-orthogonal to  $\sigma_1$  that is mirrored partly  
8 408 in the macroscopic behaviour with positive volumetric strains during compression (Figure  
9 409 6c) and total absence of dilatancy. This particular non-dilatant macroscopic behaviour was  
10 410 also observed during triaxial compression at low confining pressure of 5 MPa (Figure 5c).  
11 411 Associated damage process does not significantly affect the P-wave velocity and maximum  
12 412 amplitude (Figures 8c and 8d) and also the apparent Young's modulus, which seems to be  
13 413 constant (Figure 7b). Beyond the  $\sigma_C$  threshold, shear mechanisms are dominant until failure  
14 414 (Figure 10b, between 30 % and 70 % of compressive first motions). As suggested by Es-  
15 415 cartin et al. (1997) in their study of non-dilatant brittle deformation of serpentinites through  
16 416 mechanical measurements and direct observations, shear mechanisms can involve dilatancy  
17 417 localized on irregularities on the fracture plane. But this localized dilatancy does not affect  
18 418 macroscopic mechanical behaviour as dilatancy induced by extensive axial microcracks. In  
19 419 sedimentary porous rocks, shear mechanisms should result from slides along inter-granular  
20 420 discontinuities and also from grain rotations and translations. AE events that have shear  
21 421 mechanisms are very energetic, as shown in Figure 9d, especially during the pre-peak phase  
22 422 where the stress-displacement curve becomes non-linear. We assume that such AE events were  
23 423 produced by sources larger than grain size i.e. along inter-granular discontinuities. These inter-  
24 424 granular discontinuities with shear mechanisms do not affect elastic wave propagation (axial  
25 425 and transversal velocities are constant, Figure 8c). They should coalesce to form macroscopic  
26 426 fractures leading to failure (Figure 12).

27 427 During GGP compression, we observed stress-induced velocity anisotropy through the de-  
28 428 crease of transversal velocities (Figure 8a) and attenuation variations in axial and transversal  
29 429 directions (Figure 8b). This behaviour is associated on the one hand with pre-existing mi-  
30 430 crocrack closure and, to a less extent, with pore crushing and with the opening of new voids  
31 431 such as extensive axial microcracks, on the other hand (Sayers et al. 1990). For GGP, where  
32 432 the bedding planes are subparallel to the compression direction, such an inherent anisotropy  
33 433 due to the formation of the rock influences deformation and damage processes (Chow et al.  
34 434 1995; Gatelier et al. 2002). Bedding planes may constitute weakness planes subparallel to the  
35 435 compression direction in the GGP sample. Opening of extensive axial microcracks can occur  
36 436 on these planes. During GGP compression, stress-induced velocity anisotropy and dilatancy  
37 437 may be related to anisotropic damage processes (with opening of extensive microcracks, the  
38 438 largest dimensions of which are sub-parallel to  $\sigma_1$ ). Concerning BGP compression, the mea-  
39 439 surements do not present any evidence of anisotropy. Since the bedding planes of this stratum



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440 have been reworked, the BGP sample does not contain clear inherent anisotropy.  
441 The presented results coming from independent measurements seem consistent for both sam-  
442 ple compressions. Additional experiments (i.e. further uniaxial compressive tests for both  
443 grey-iron layer and brown-iron layer rocks) need to be performed to reinforce our hypothesis  
444 concerning the micro-mechanisms deduced from AE analysis and to state that the acousto-  
445 mechanical behaviours observed are representative for the rock types tested. The reproducibil-  
446 ity of the acoustical observations should be statistically evaluated especially in the case of the  
447 heterogeneous sedimentary ferriferrous formations studied here.

## 448 5 CONCLUSIONS AND PERSPECTIVES

449 Most of the experiments performed provided four indirect measurements of damage evolu-  
450 tion, i.e. the P-wave velocity and maximum amplitude received during the pulse transmission  
451 experiment, elastic properties (apparent Young's modulus and apparent Poisson's ratio) and  
452 AE monitoring. Variations in elastic properties are consequences of microstructural changes  
453 (damage processes) that occur during compression. Dilatancy of the GGP sample, variations  
454 in the P-wave velocity and maximum amplitude and the analysis of AE source mechanisms  
455 lead to the hypothesis of the development of extensive axial microcracks, their interaction  
456 and coalescence constitute shear fractures that split the sample. For BGP non-dilatant rock,  
457 damage processes inferred from AE source mechanism analysis involve microscopic compres-  
458 sive and shear mechanisms that interact and lead to macroscopic fractures. Figure 13 shows  
459 a summary of the main interpretations. In both cases, GGP or BGP compression tests, mi-  
460 croscopic processes deduced from AE analysis are consistent with macroscopic mechanical  
461 behaviour, P-wave velocity and maximum amplitude changes.

462 In order to complete these results, direct microstructural observations should reinforce the  
463 hypothesis concerning development of extensive axial microcracks observed during the GGP  
464 compression. Concerning BGP, since shear mechanisms are usually propagated in sedimen-  
465 tary rocks along the grain boundaries (Kranz 1983), direct observation could be difficult.  
466 Localization of AE sources and focal mechanisms deduced from first motion analysis could  
467 provide, as for seismic events, planes of possible shearing and hence the orientation of shear  
468 discontinuities.

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*Evidence of dilatant and non-dilatant damage processes in oolitic iron ore: P-wave velocity and acoustic emission***Table 1.** Synopsis of the laboratory experiments performed.

Sample	Sedimentary layer	Measurement type	Acoustic sensors	Channels for triggering	Pc (MPa)	$(\sigma_1 - Pc)_{max}$ (MPa)	AE events recorded
GGP1-A1	grey-iron	MA	8	1	0	27	1030
GGP1-A2	grey-iron	MA	8	1	0	28.4	1233
GGP2-2	grey-iron	MAVT	11	4	10	37.7	102
GGP	grey-iron	MAV	18	7	0	31.4	177
BGP	brown-iron	MAV	18	7	0	38.9	214
BGP2	brown-iron	MAVT	11	4	5	54.1	149
GM0203S1	intercalated marl	MAV	11	4	0	14.7	675
GM0203S2	intercalated marl	MAV	11	4	0	54.3	473

"MAVT" type corresponds to a Triaxial ("T") compression test including Mechanical ("M"), Acoustical ("A", AE monitoring) and P-wave Velocity ("V") measurements, "MAV" type includes the same measurements during uniaxial compression. "MA" type is for Mechanical ("M") and Acoustical ("A") measurements performed during uniaxial compressive test. If signals exceeded the trigger threshold (generally set to 50 mV) on at least the number of "Channels for triggering", the acoustical acquisition system triggered and the signal was recorded. The uniaxial compressive tests included six strain gauges whereas triaxial ones only four. Pc is the confining pressure.

**Table 2.** Stress thresholds ( $\sigma_{CC}$ ,  $\sigma_C$ ,  $\sigma_{LL}$ ,  $\sigma_D$ ,  $\sigma_M$ ) for each sample.

Sample	$\sigma_{CC}$ (MPa)	$\sigma_C$ (MPa)	$\sigma_{LL}$ (MPa)	$\sigma_D$ (MPa)	$\sigma_M$ (MPa)
GGP	5	11.4	18.6	24.2	31.4
BGP	5	11	15	-	38.9

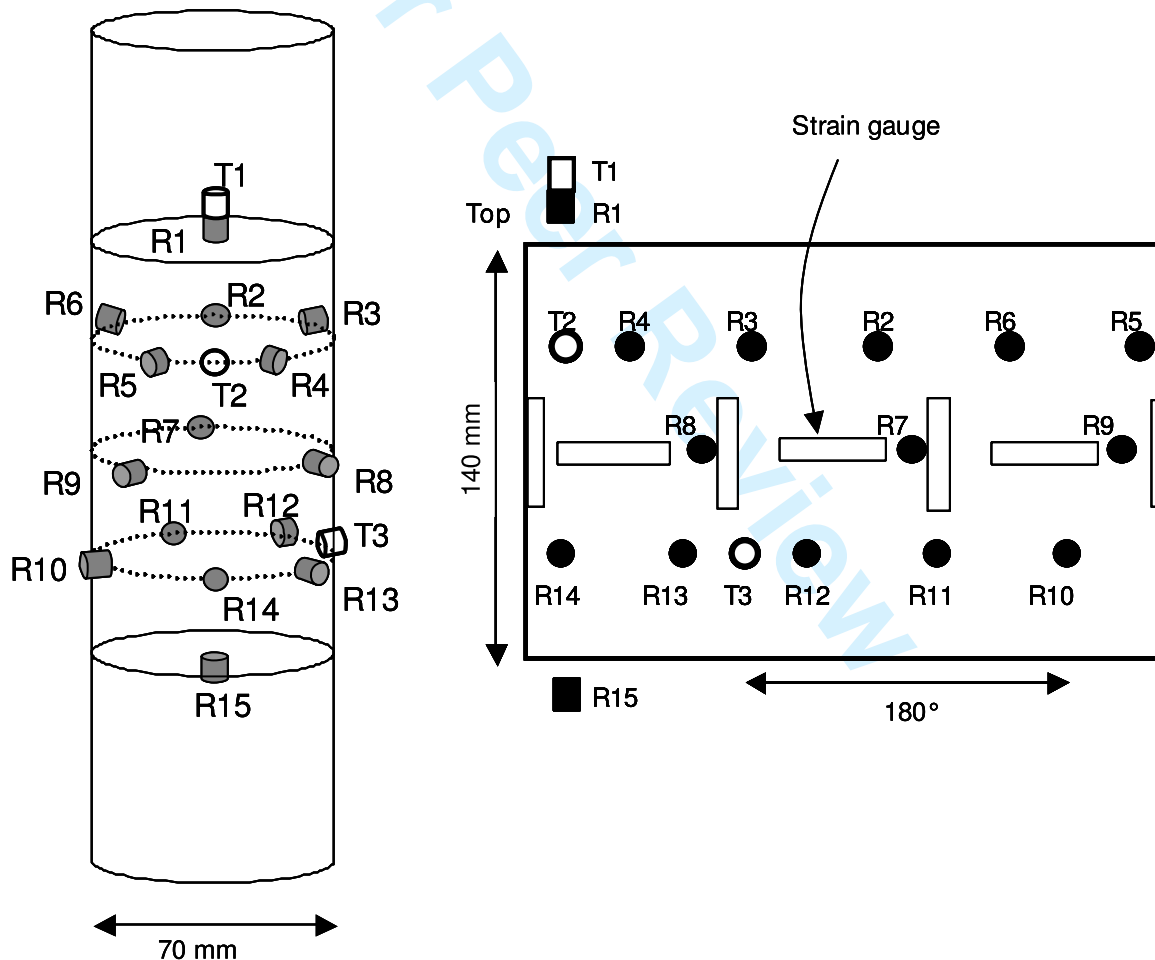
$\sigma_{CC}$  determined on the stress-displacement curve is the initial crack closure threshold. The other thresholds are deduced from stress-strain curves,  $\sigma_C$ ,  $\sigma_{LL}$  and  $\sigma_D$  are respectively the cracking initiation determined on the transversal strain-stress curve, the loss of linearity of the longitudinal strain-stress curve and the dilatancy threshold observed on the volumetric strain curve. The failure stress  $\sigma_M$  is the maximum stress determined on the stress-displacement curve.

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**Table 3.** Mechanical and acoustical characteristics and P-wave velocity for the two samples.

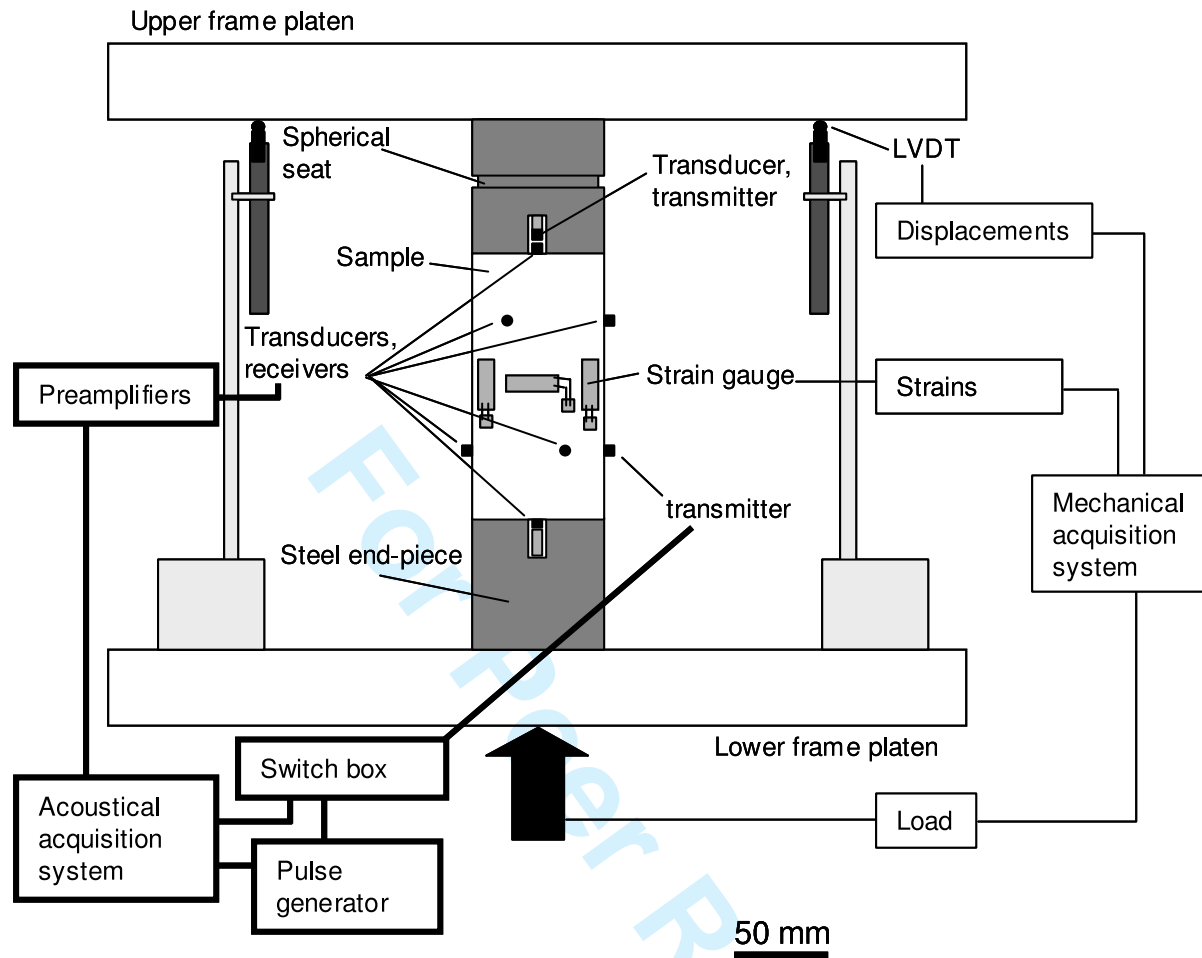
Sample	$E_{app}$ (GPa)	$\nu_{app}$	$V_{p_{axial}}$ m/s	$V_{p_{trans}}^{top}$ m/s	$V_{p_{trans}}^{bottom}$ m/s	AE events	$b_e$
GGP	29.2	0.28	4171	3981	4143	112	0.63
BGP	19.9	0.11	3160	3058	3031	160	0.51

$E_{app}$  is the apparent Young's modulus,  $\nu_{app}$  the apparent Poisson's ratio.  $V_{p_{axial}}$  is measured along T1R15 ray,  $V_{p_{trans}}^{top}$  along T2R2 transversal ray and  $V_{p_{trans}}^{bottom}$  along T3R10 transversal ray. An AE event is identified between P-pick and end-pick on at least 8 channels.  $b_e$  value is estimated from AE energy distribution ( $\log(N) = c - b_e \log(E)$ , where  $N$  is the number of AE events that have signal energy greater than or equal to  $E$ ,  $c$  and  $b_e$  are constants, and  $E$  the AE event signal energy, see equation 4).



**Figure 1.** Sample instrumentation scheme. Strain gauges (rectangles in the central part of the sample) and piezoelectric transducers are glued on the sample surface. T transducers are transmitters and R are receivers.

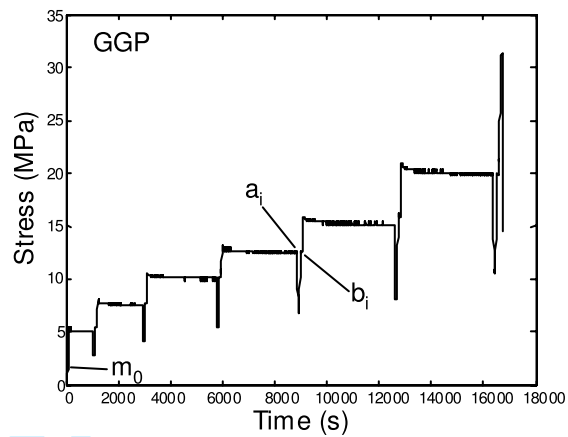
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**Figure 2.** Experimental system scheme. Mechanical acquisition system includes six strain gauges glued on the lateral surface of the sample and two LVDT placed between the press platens. The acoustical acquisition system is constituted by transducers (3 transmitters and 15 receivers), which are coupled to the sample surface, preamplifiers, acquisition analogical/digital card, and also, for the pulse transmission experiment, a pulse generator and a switch box.

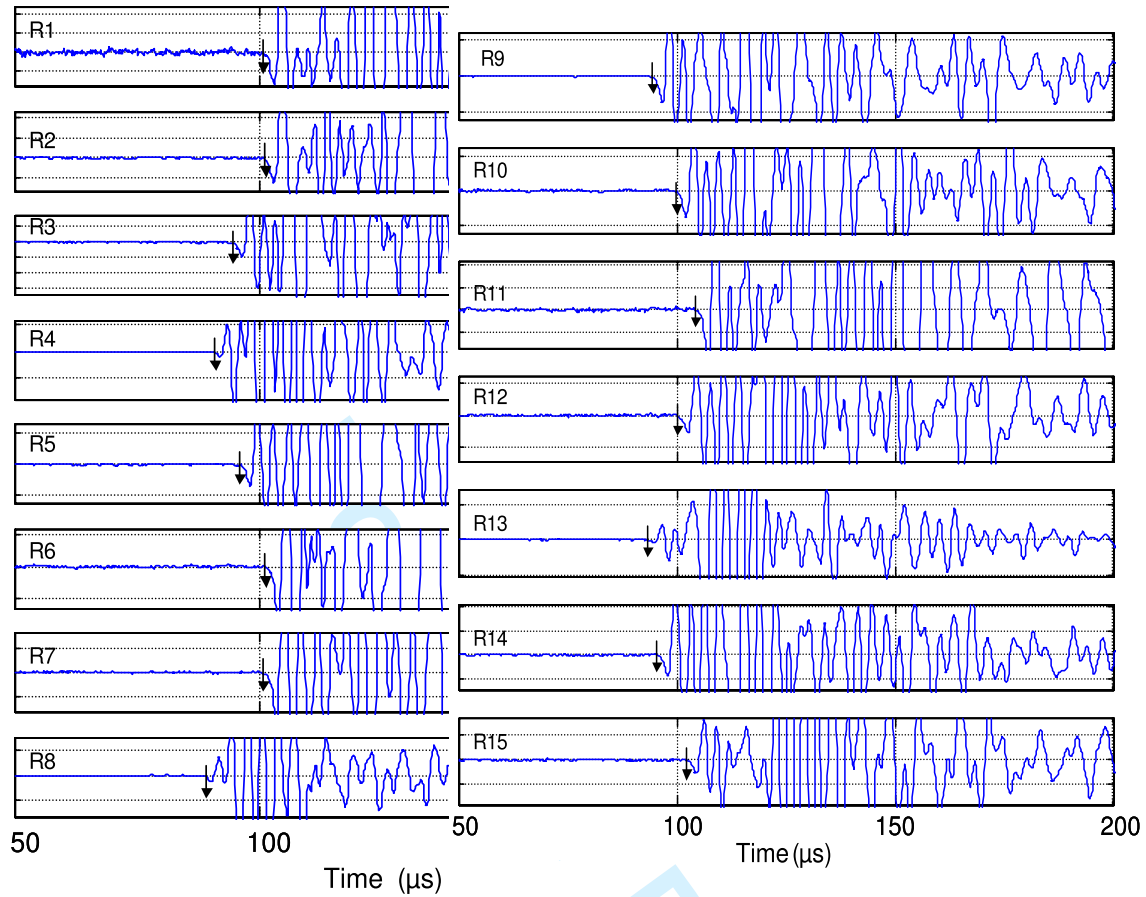


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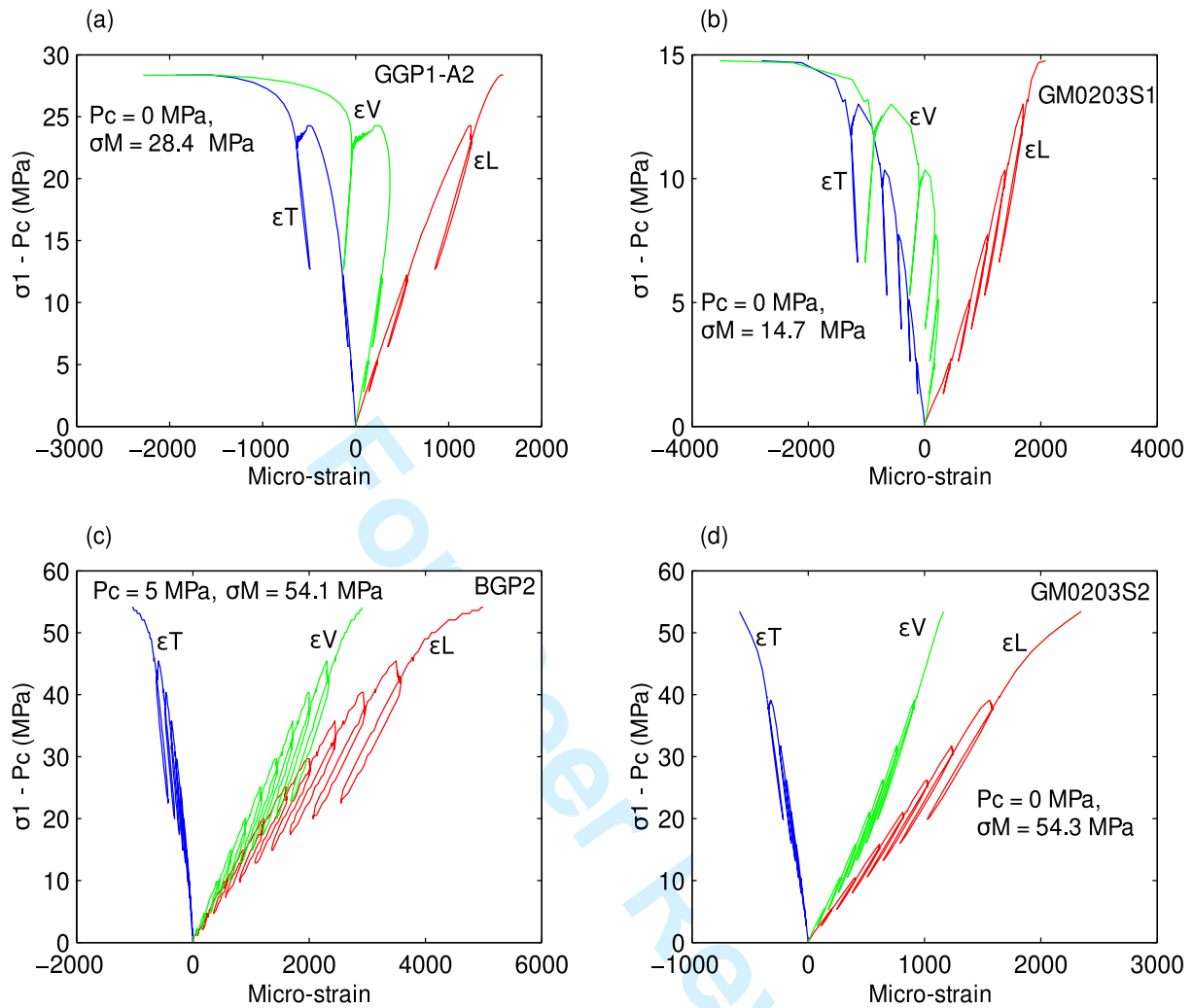
**Figure 3.** Loading history with loading-relaxation-unloading cycles. The stress is plotted versus time. Velocity measurements (in axial and transversal directions) were performed at points  $m_0$ ,  $a_i$  and  $b_i$  ( $i$ : cycle number).

Evidence of dilatant and non-dilatant damage processes in oolitic iron ore: P-wave velocity and acoustic emission



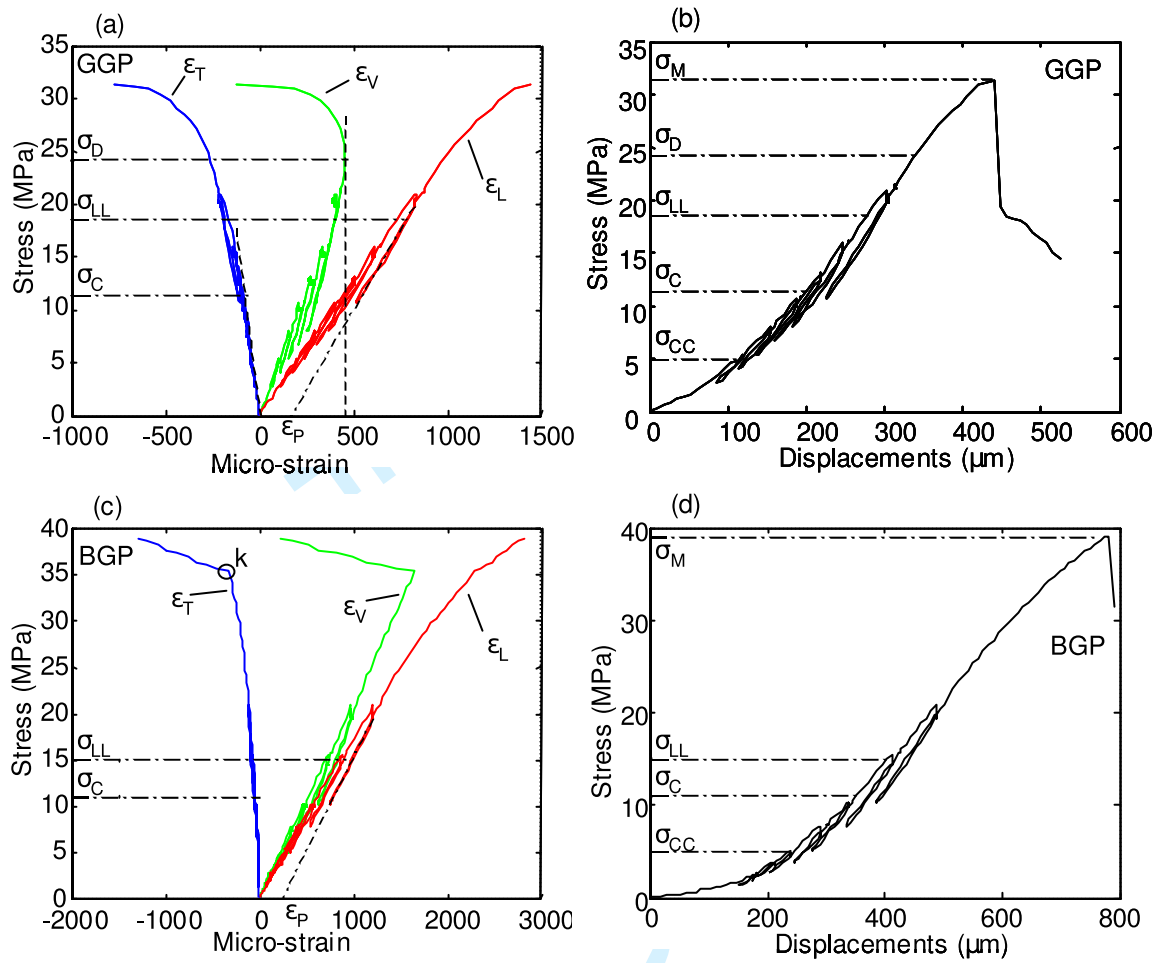
**Figure 4.** Enlargement of first motions of signal traces on 15 receivers (R1 to R15) of an AE event recorded during GGP compression test. The transducer polarisation is negatively polarised i.e. negative oscillation corresponds to compressional first motion. First motions showed here are all compressional. Hence, this event can result from a tensile source.

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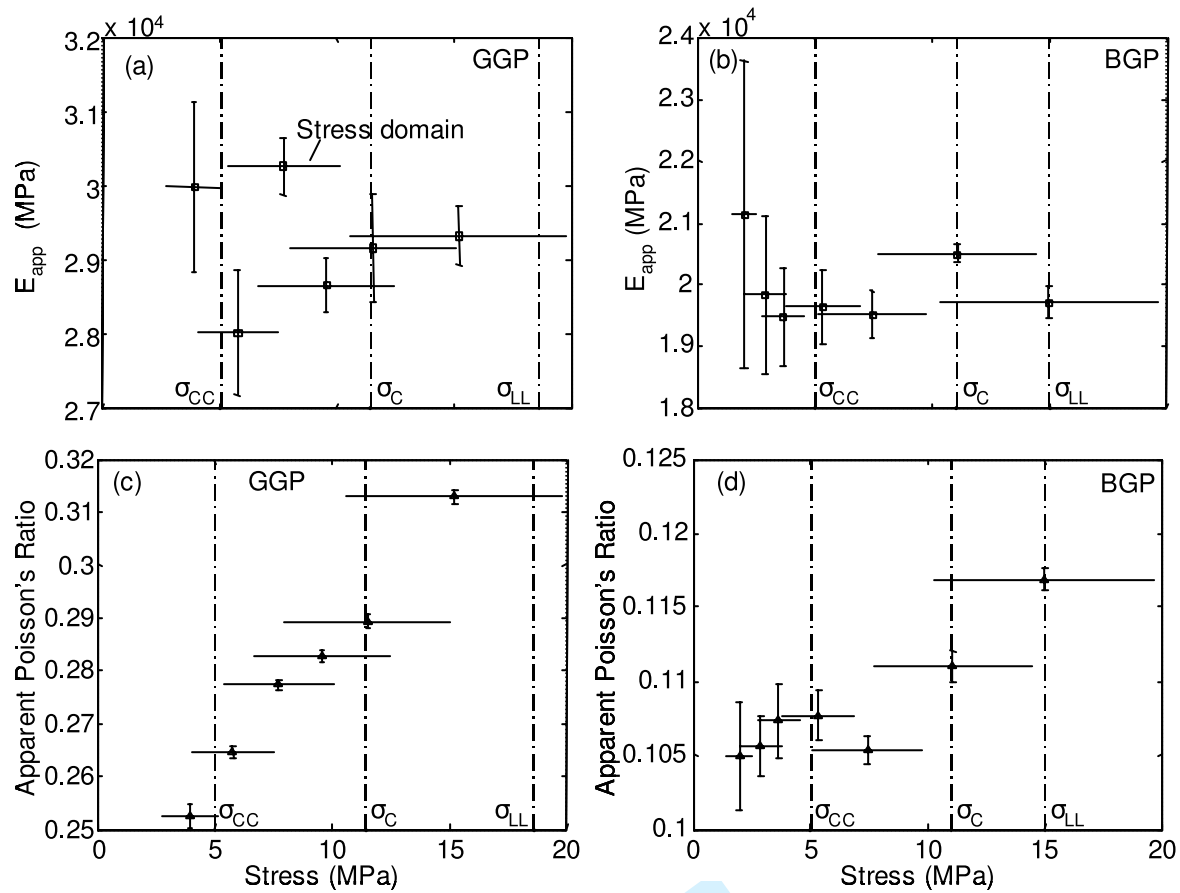
**Figure 5.** Mechanical behaviours of samples coming from different levels of the iron ore and intercalated marl beds. a) sample GGP1-A2 from the grey-iron layer; c) sample BGP2 from the brown-iron layer; b) and d) samples GM0203S1 and GM0203S2 from different facies of the intercalated marl layer. The differential stress ( $\sigma_1 - P_c$ , with  $P_c$  the confining pressure) is plotted versus transversal, longitudinal and volumetric strains, respectively,  $\epsilon_T$ ,  $\epsilon_L$  and  $\epsilon_V$  measured using strain gauges.  $\sigma_M$  is the maximum stress value.

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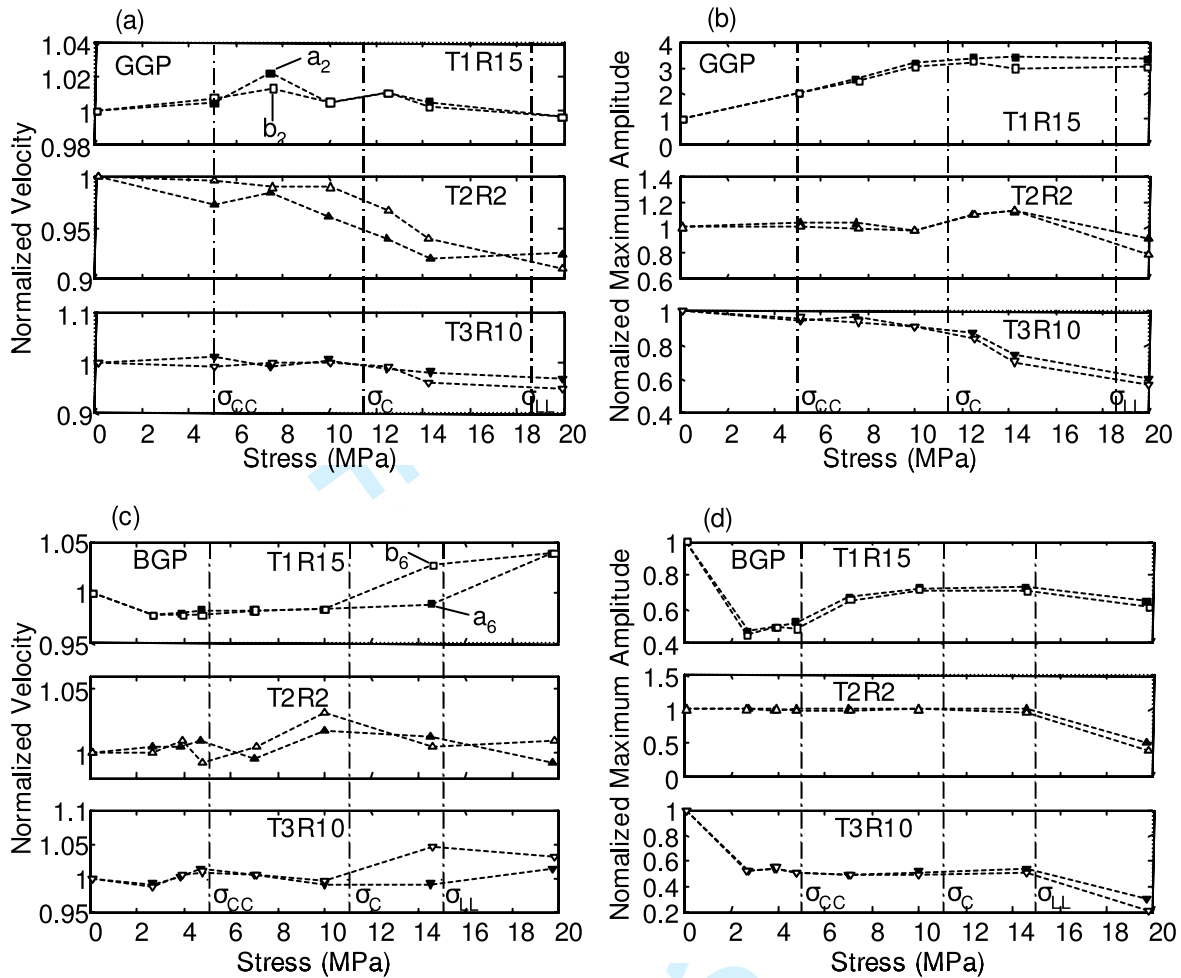
**Figure 6.** Mechanical behaviour of uniaxially compressed GGP (a and b) and BGP (c and d) samples. The stress is plotted versus transversal, longitudinal and volumetric strains, respectively,  $\epsilon_T$ ,  $\epsilon_L$  and  $\epsilon_V$  measured using strain gauges. Stress is also plotted as a function of displacements measured between the press platens. Stress thresholds were graphically determined.  $\sigma_{CC}$  determined on the stress-displacement curve is the initial crack closure threshold. The other thresholds came from stress-strain curves.  $\sigma_C$ ,  $\sigma_{LL}$  and  $\sigma_D$  are respectively the cracking initiation determined on the transversal strain-stress curve, the loss of linearity of the longitudinal strain-stress curve and the dilatancy threshold observed on the volumetric strain-stress curve. The peak stress  $\sigma_M$  is the maximum stress value determined on the stress-displacement curves. In c), the point k on the transversal strain-stress curve corresponds to the failure of transversal strain gauges that is mirrored in the volumetric strain-stress curve.

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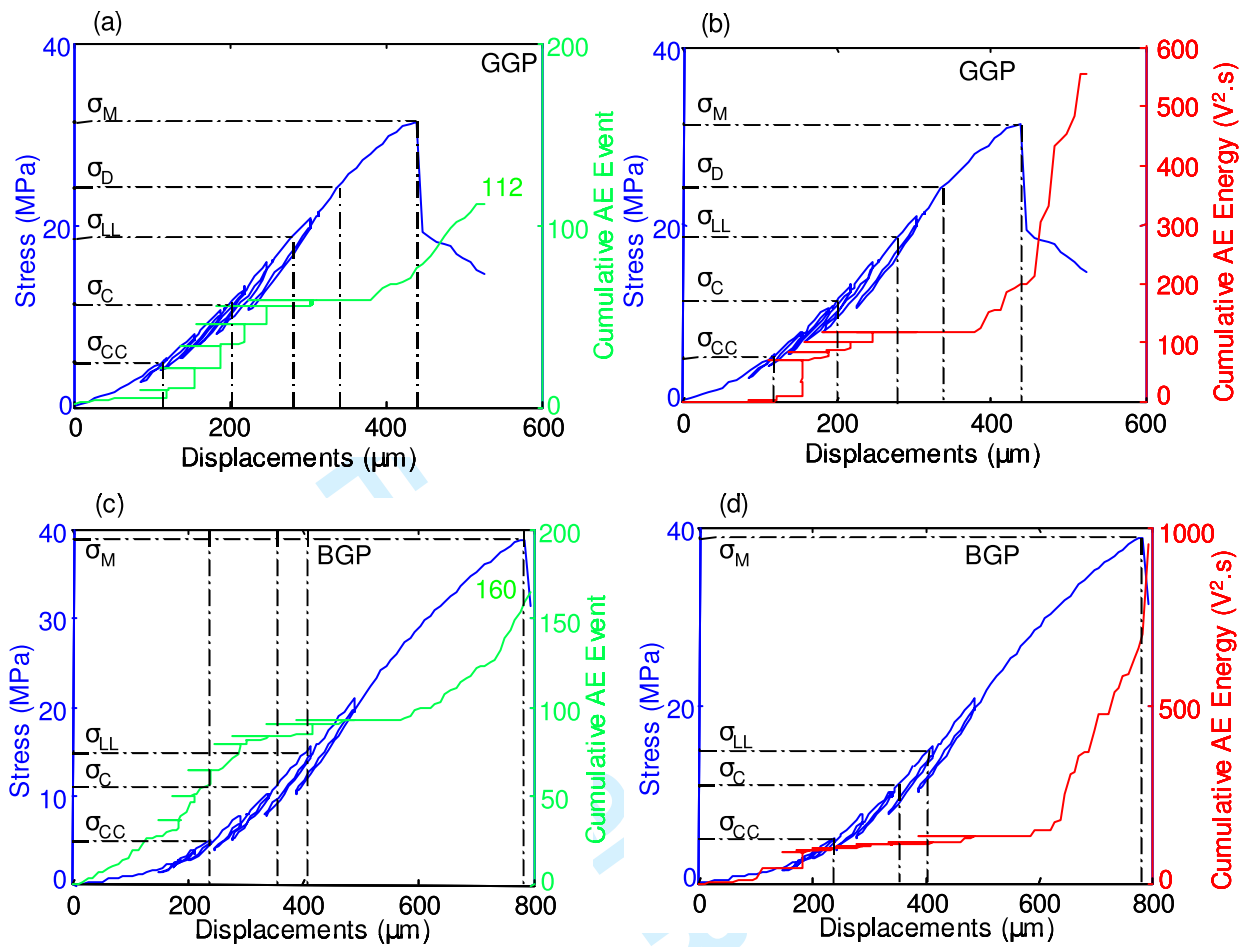
**Figure 7.** GGP elastic properties plotted versus stress: apparent Young's modulus ( $E_{app}$ , a), and apparent Poisson's ratio (c). BGP elastic properties plotted versus stress: apparent Young's modulus ( $E_{app}$ , b) and apparent Poisson's ratio (d). Each horizontal line represents the stress domain of each unloading phase where linear regression was performed to determine the modulus value and associated error (vertical segment). Stress thresholds  $\sigma_{CC}$ ,  $\sigma_C$  and  $\sigma_{LL}$ , are reported on the X-axis.

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**Figure 8.** Normalized velocity ( $V_p/V_{p(\sigma_1=0)}$ , a) and normalized maximum amplitude ( $Am/Am_{(\sigma_1=0)}$ , b) of the signal received during velocity measurements performed in compression of GGP. ( $V_p/V_{p(\sigma_1=0)}$ ) and ( $Am/Am_{(\sigma_1=0)}$ ) are plotted versus differential stress for axial ray (T1R15) and transversal rays (T2R2 and T3R10). Velocity and maximum amplitude are both represented for measurement points  $a_i$  (end of relaxation phases, black markers) and points  $b_i$  (stress level equal to this  $a_i$  during the loading phases, white markers). Stress thresholds  $\sigma_{CC}$ ,  $\sigma_C$  and  $\sigma_{LL}$  are reported on the X-axis. Normalized velocity (c) and normalized maximum amplitude (d) of the signal received during velocity measurements performed in compression of BGP.

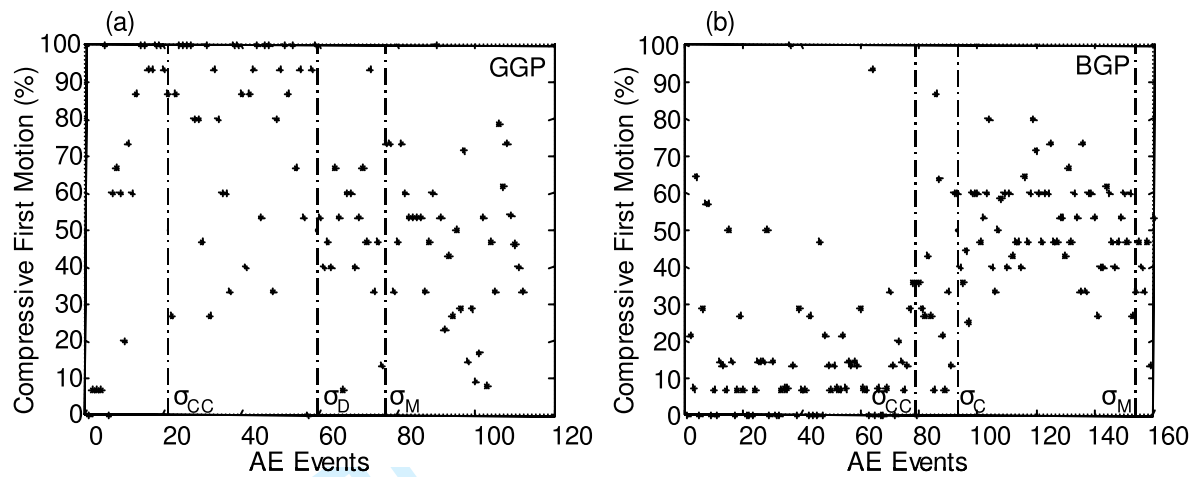
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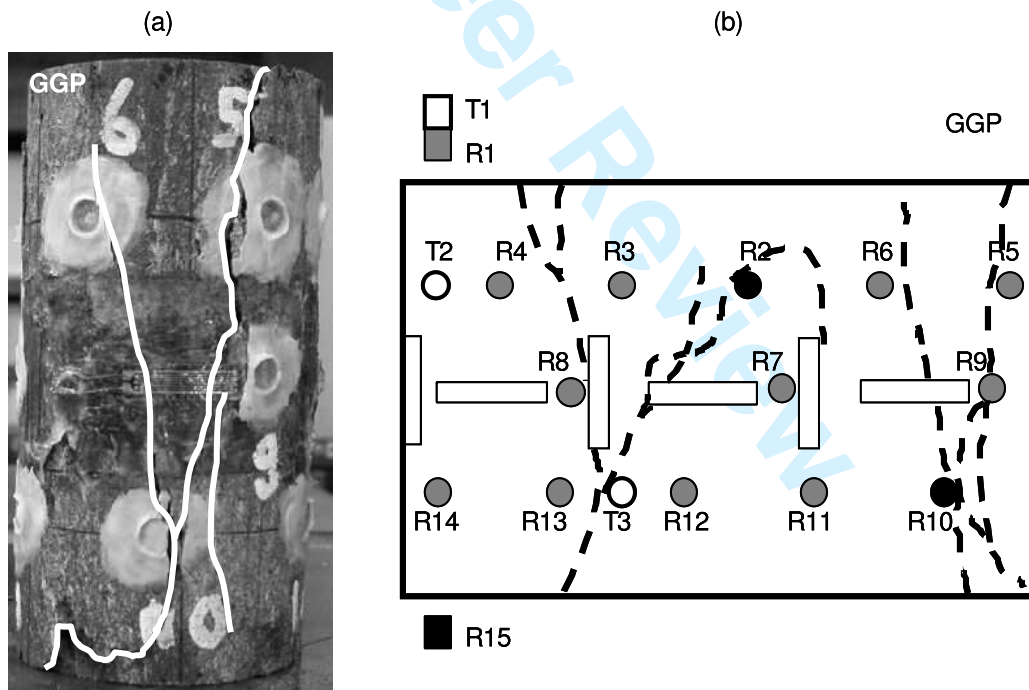
**Figure 9.** (a) Mechanical behaviour and acoustic activity during uniaxial compression of GGP. (b) Stress-displacement curve and AE energy during uniaxial compression of GGP. (c) Mechanical behaviour and acoustic activity during uniaxial compression of BGP. (d) Stress-displacement curve and AE energy for the BGP sample. Stress thresholds are reported on the Y-axis.



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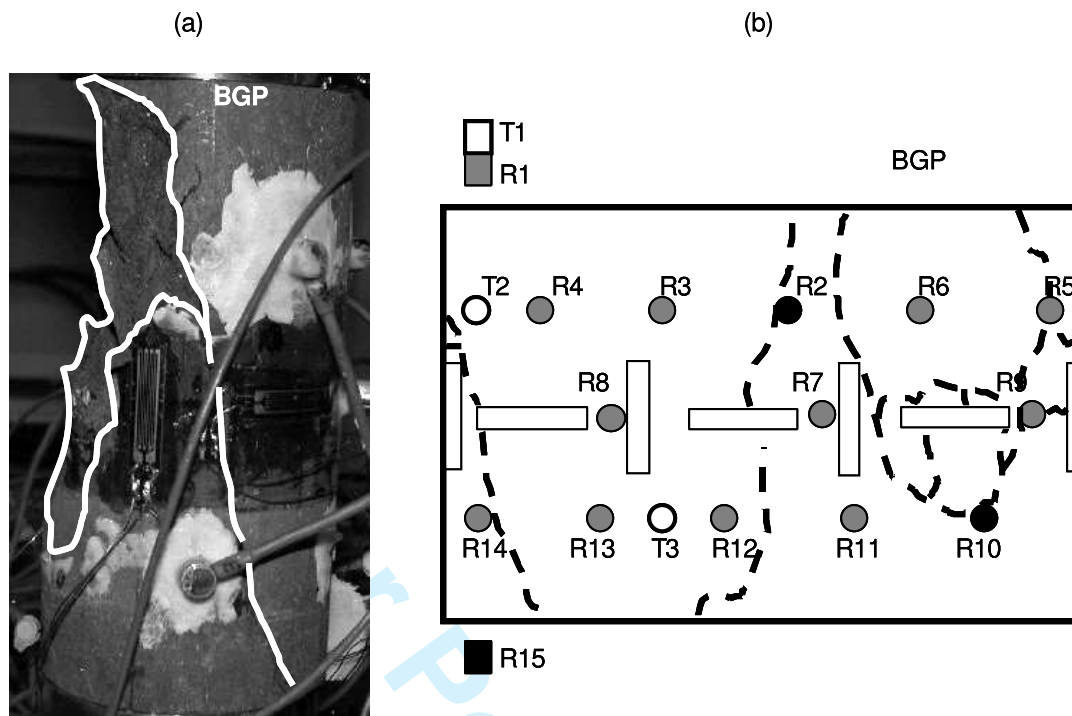


**Figure 10.** Compressive first motion ratio for each AE event recorded during GGP uniaxial compression (a) and during BGP uniaxial compression (b). Stress thresholds are reported on the X-axis. 100 % of compressive first motions correspond to pure extensive mechanism whereas 0 % of compressive first motions are related to pure compressive mechanism. Shear mechanisms induce AE events with compressive first motion ratio between 30 % and 70 %.



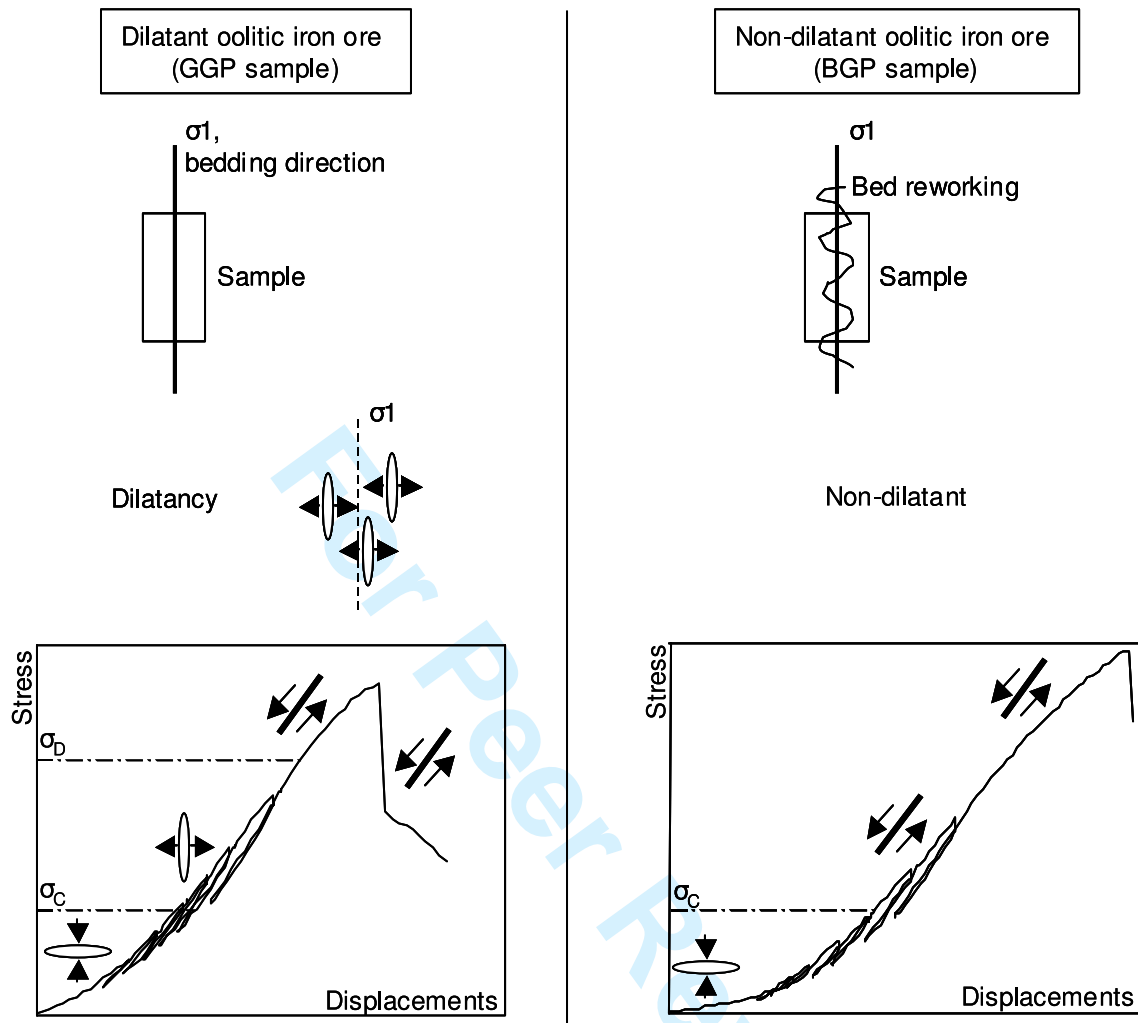
**Figure 11.** (a) Picture of the broken GGP sample showing irregular longitudinal splitting. (b) GGP sample scheme with strain gauges (rectangles in the central part of the sample) and piezoelectric transducer locations (circles on the lateral surface of the sample, rectangles at the both ends). The dashed lines represent the final macroscopic fracture traces.

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30 **Figure 12.** (a) Picture of the broken BGP sample. (b) BGP sample scheme with strain gauges (rect-  
31 angles in the central part of the sample) and piezoelectric transducer locations (circles on the lateral  
32 surface of the sample, rectangles at the both sample ends). The dashed lines represent the final macro-  
33 scopic fracture traces.  
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Evidence of dilatant and non-dilatant damage processes in oolitic iron ore: P-wave velocity and acoustic emission



**Figure 13.** Diagram showing possible damage processes for dilatant and non-dilatant oolitic iron ores (respectively, GGP and BGP samples). Both sample deformations include damage processes. For the dilatant GGP sample, tensile cracking parallel to the  $\sigma_1$  direction develops prior to the  $\sigma_D$  threshold. These extensive microcracks interact to form shear cracking. For the BGP non-dilatant sample, shear cracking develops beyond the  $\sigma_C$  threshold.