#### **Physical interpretations**

Time reversal **Rays and stationary phase** 

Example: an even uncorrelated distribution of source (acoustic case) (Roux et al., 2004)

$$C_{12}(\tau) = \int_{-\infty}^{\infty} P(\mathbf{r}_1, t) P(\mathbf{r}_2, t + \tau) \mathrm{d} t. \quad \text{where,}$$

$$P(\mathbf{r}_i,t) = \int_0^\infty d\mathbf{r}_s \int_{-\infty}^t dt_s \ S(\mathbf{r}_s;t_s)g(\mathbf{r}_s,\mathbf{r}_i;t-ts); \ i=1,2.$$

- 1. Finite recording time  $T_r$
- 2. Uncorrelated noise sources:  $\langle S(\mathbf{r}_s; t_s) S(\widetilde{\mathbf{r}}_s; t_s) \rangle = Q^2 \upsilon \delta(\mathbf{r}_s \widetilde{\mathbf{r}}_s) \delta(t_s, \widetilde{t}_s)$
- 3. Change variable:  $\hat{t} = t t_s$

$$C_{12}(\tau) = Q^2 \upsilon T_r \int_0^\infty d\mathbf{r}_s \int_0^\infty d\hat{\mathbf{t}} g(\mathbf{r}_s, \mathbf{r}_1; \hat{\mathbf{t}}) g(\mathbf{r}_s, \mathbf{r}_2; \hat{\mathbf{t}} + \tau).$$

Spatial sum Cross-correlation for each noise source

# Noise cross-correlation: Free space

Noise sources yielding constant time-delay  $\tau$ , lay on same Hyperbola



Isotropic distribution of uncorrelated random noise sources

# C, dC/dt, band-limited signal





$$G(\mathbf{x}_{B}, \mathbf{x}_{A}, t) + G(\mathbf{x}_{B}, \mathbf{x}_{A}, -t) \approx \frac{2}{\rho c} \int_{S} G(\mathbf{x}_{B}, \mathbf{x}, t) * G(\mathbf{x}_{A}, \mathbf{x}, -t) d^{2}\mathbf{x}$$



$$G(\mathbf{x}_{B}, \mathbf{x}_{A}, t) + G(\mathbf{x}_{B}, \mathbf{x}_{A}, -t) \approx \frac{2}{\rho c} \int_{S} G(\mathbf{x}_{B}, \mathbf{x}, t) * G(\mathbf{x}_{A}, \mathbf{x}, -t) d^{2}\mathbf{x}$$



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$$G(\mathbf{x}_{B}, \mathbf{x}_{A}, t) + G(\mathbf{x}_{B}, \mathbf{x}_{A}, -t) \approx \frac{2}{\rho c} \int_{S} G(\mathbf{x}_{B}, \mathbf{x}, t) * G(\mathbf{x}_{A}, \mathbf{x}, -t) d^{2} \mathbf{x}$$

#### Correlations, Green function and equipartition Elasticity

# P-SV case

# **Green function in 2D**

$$G_{ij} = \frac{i}{4\rho\omega^2} \left\{ -\delta_{ij}k^2 H_0^{(2)}(kr) + \frac{\partial^2}{\partial x_i \partial x_j} \left[ H_0^{(2)}(qr) - H_0^{(2)}(kr) \right] \delta_{lj} \right\}$$
$$G_{ij}(P,Q) = \frac{-i}{8\rho} \left\{ A \delta_{ij} - B \left( 2\gamma_i \gamma_j - \delta_{ij} \right) \right\} \qquad \gamma_j = \frac{x_j - \xi_j}{r}$$



Hankel functions

 $q = \frac{\omega}{\alpha}$   $k = \frac{\omega}{\beta}$   $\alpha = \sqrt{\frac{\lambda + 2\mu}{\rho}}$   $\beta = \sqrt{\frac{\mu}{\rho}}$  r = |P,Q|

#### THE 2D VECTOR CASE

 $\bigcap$ 

$$\beta^2 \frac{\partial^2 u_i}{\partial x_j \partial x_j} + (\alpha^2 - \beta^2) \frac{\partial^2 u_j}{\partial x_i \partial x_j} = \frac{\partial^2 u_i}{\partial t^2}$$

Summation of P and S plane waves:

$$u_{l}(\mathbf{x},\omega,t) = P(\omega,\phi)n_{l}\exp(-i\frac{\omega}{\alpha}x_{j}n_{j}) + S(\omega,\psi)m_{l}^{*}\exp(-i\frac{\omega}{\beta}x_{j}m_{j})$$
  
Forrelation:

$$u_{l}(\mathbf{y})u_{s}(\mathbf{x}) = (P^{2}n_{l}n_{s} + SP^{*}m_{l}n_{s})\exp(\mathbf{i}kr\cos[\phi - \theta]) + (S^{2}m_{l}m_{s}^{*} + PS^{*}n_{l}m_{s}^{*})\exp(\mathbf{i}kr\cos[\psi - \theta])$$

$$u_1(\mathbf{y})u_1^*(\mathbf{x}) = (P^2 \cos^2 \phi - SP^* \sin \psi \cos \phi) \exp(\mathbf{i}qr \cos[\phi - \theta]) + (S^2 \sin^2 \psi + PS^* \cos \phi \sin \psi) \exp(\mathbf{i}kr \cos[\psi - \theta])$$

Azimuthal average:

$$\langle \bullet \rangle = \frac{1}{4\pi^2} \int_0^{2\pi} d\phi \int_0^{2\pi} \bullet d\psi$$

$$\left\langle u_1(\mathbf{y})u_1^*(\mathbf{x})\right\rangle = \frac{P^2\alpha^2}{2} \frac{J_0(qr)}{\alpha^2} + \frac{S^2\beta^2}{2} \frac{J_0(kr)}{\beta^2} - \frac{P^2\alpha^2}{2} \frac{J_2(qr)}{\alpha^2} \cos 2\theta + \frac{S^2\beta^2}{2} \frac{J_0(kr)}{\beta^2} \cos 2\theta \right\rangle$$

with  $P^2 \alpha^2 = \varepsilon S^2 \beta^2$ 

$$\langle u_i(\mathbf{y})u_j^*(\mathbf{x})\rangle = \frac{S^2\beta^2}{2} \{A\delta_{ij} - B(2\gamma_i\gamma_j - \delta_{ij})\}$$

$$A = \varepsilon \frac{J_0(qr)}{\alpha^2} + \frac{J_0(kr)}{\beta^2} \text{ and } B = \varepsilon \frac{J_2(qr)}{\alpha^2} - \frac{J_2(kr)}{\beta^2}$$

$$P^2 \alpha^2 = \varepsilon S^2 \beta^2$$

## Equipartition (*ε*=1):

$$\left\langle u_{i}(\mathbf{y})u_{j}^{*}(\mathbf{x})\right\rangle = \frac{S^{2}\beta^{2}}{2} \left\{ A\delta_{ij} - B(2\gamma_{i}\gamma_{j} - \delta_{ij}) \right\}$$

$$A = \varepsilon \frac{J_{0}(qr)}{\alpha^{2}} + \frac{J_{0}(kr)}{\beta^{2}} \text{ and } B = \varepsilon \frac{J_{2}(qr)}{\alpha^{2}} - \frac{J_{2}(kr)}{\beta^{2}}$$

$$\downarrow$$

$$\downarrow$$

$$\left\langle u_{i}(\mathbf{y},\boldsymbol{\omega})u_{j}^{*}(\mathbf{x},\boldsymbol{\omega})\right\rangle = -8E_{S}k^{-2}\operatorname{Im}\left[G_{ij}(\mathbf{x},\mathbf{y},\boldsymbol{\omega})\right]$$

Formally, same result in 3D (Sánchez-Sesma and Campillo, 2006)

#### Correlations, Green function and equipartition Locally heterogeneous body

#### SH waves in a medium with a cylindrical cavity

Analytical solution for the 2D SH Green function:



$$G_{22}(\mathbf{x}, \mathbf{y}; \omega) = v^{0} + v^{d} = \frac{1}{4i\mu} \left\{ H_{0}^{(2)}(kR) + \sum_{n=0}^{\infty} \varepsilon_{n} A_{n} H_{n}^{(2)}(kd) H_{n}^{(2)}(kr) \cos n\theta \right\}.$$

$$A_{n} = -\frac{J_{n}(ka) J_{n}^{'}(qa) - \xi J_{n}^{'}(ka) J_{n}(qa)}{H_{n}^{(2)}(ka) J_{n}^{'}(qa) - \xi H_{n}^{(2)'}(ka) J_{n}(qa)}. \quad \text{with} \quad \xi = \frac{\mu_{E} k}{\mu_{\Gamma} q} = \frac{\rho_{E} \beta_{E}}{\rho_{\Gamma} \beta_{\Gamma}}$$

Consider an incident plane wave:

 $v(\mathbf{x}, \omega, t) = F(\omega, \psi) \exp(-i k x_j n_j) \exp(i \omega t),$ 



Expansion (polar coordinates):

$$v^{\theta}(\mathbf{x}, \omega) = F(\omega) \sum_{n=0}^{\infty} \varepsilon_n \mathbf{i}^n J_n(kr) \cos n(\psi - \theta)$$
, and

and the diffracted field

$$v^{d}(\mathbf{x}, \omega) = F(\omega) \sum_{n=0}^{\infty} \varepsilon_{n} \mathbf{i}^{n} A_{n} H_{n}^{(2)}(kr) \cos n(\psi - \theta)$$

The incident plane wave and the diffracted field reduce to

$$v(r, \theta; \psi) = F(\omega) \sum_{n=0}^{\infty} V_n(r, \omega) \cos n(\theta - \psi)$$

where

$$V_n(r,\omega) = \mathbf{i}^n \varepsilon_n \left[ J_n(kr) - A_n H_n^{(2)}(kr) \right],$$

Correlation:

$$v(\mathbf{y}, \boldsymbol{\omega})v^*(\mathbf{x}, \boldsymbol{\omega}) = F^2(\boldsymbol{\omega}) \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} V_n(d, \boldsymbol{\omega}) V_m^*(r, \boldsymbol{\omega}) \cos n \, \psi \cos m(\psi - \theta)$$

azimuthal average:equipartition of the incident field

$$\langle v(\mathbf{y}, \omega) v^*(\mathbf{x}, \omega) \rangle = F^2(\omega) \sum_{m=0}^{\infty} \frac{1}{\varepsilon_m} V_m(d, \omega) V_m^*(r, \omega) \cos m\theta$$

$$\frac{1}{\varepsilon_m}V_m(d,\omega)V_m^*(r,\omega) = \frac{\varepsilon_m}{D_m^2}(N_mY_m(kd) - M_mJ_m(kd))(N_mY_m(kr) - M_mJ_m(kr)).$$

Noting that:

$$\operatorname{Im}[G_{22}(\mathbf{x},\mathbf{y};\omega)] = \frac{-1}{4\mu} \sum_{m=0}^{\infty} \frac{\varepsilon_m}{D_m^2} (N_m Y_m(kd) - M_m J_m(kd)) (N_m Y_m(kr) - M_m J_m(kr)) \cos m\theta$$



the exact relation:

$$\langle v(\mathbf{y}, \boldsymbol{\omega}) v^{*}(\mathbf{x}, \boldsymbol{\omega}) \rangle = -8E_{SH}k^{-2} \operatorname{Im}[G_{22}(\mathbf{x}, \mathbf{y}, \boldsymbol{\omega})].$$

including near filed, scattering, resonances,.....

one day of seismic record



Networks with continuous recording = huge amount of noise data!



## Measurement Method Applied to a Pair of North American Stations







## 1. Raw data (January 18, 2002)





#### 1. Raw data (January 18, 2002)









#### 2. Filtered seismograms (0.01-0.025 Hz)







time (s)

## 3. One-bit normalization









#### 3. One-bit normalization







3. One-bit normalization



# Cross-correlations of seismic noise: ANMO - CCM

(from Shapiro and Campillo, GRL, 2004)

30 days of vertical motion



### Dispersion analysis







Origin of the seismic noise

#### Isotropic distribution of sources: symmetric cross-correlation







#### Anisotropic distribution of sources: asymmetric cross-correlation



# Tracking the origin of the seismic noise















Winter 10–20s

Summer 10-20s



#### Apparent origin of the noise



winter

summer



#### Average sea wave height





# Imaging California with earthquake data



## Trajets déduits du bruit (~3000 paires)



High resolution velocity map obtained from noise (Rayleigh 7.5 s)



Shapiro, Campillo, Stehly et Ritzwoller, 2005

High resolution velocity map obtained from noise (Rayleigh 7.5 s)



High resolution velocity map obtained from noise (Rayleigh 7.5 s)











## Comparison noise correlation vs earthquake data

18 s cross-correlation



18 s global surface-wave tomography





-1.5

6 8 Range (km)

The 'correlation relation' (under the hypothesis of random sources or multiple scattering):

$$\partial_{\tau}C_{AB}(\tau) \propto G^{+}(A, B, \tau) - G^{-}(A, B, -\tau)$$

Average correlation of fields in A and B

1600

ΖZ

Green function between A and B (for positive and negative times)

Superposition of G(t) and G(-t).....



Rayleigh waves across USArray (from P. Boué, UJF)

#### Surface wave imaging with ambient seismic noise..... it works!

Map of Rayleigh wave group velocity in California



Shapiro et al., 2005.

The Moho beneath the Alps



Stehly et al. ,2009

Indeed, the quality depends on the distribution of 'noise' sources:

-analysis of specific noise characteristics

-the technique is robust for smooth azimuthal distributions of noise intensity

-clock errors detection and correction by time symmetry

-use of information redundancy:

9 component correlation tensor

positive and negative time vs. time symmetry

#### Surface wave imaging with ambient seismic noise..... it works!

Map of Rayleigh wave group velocity in California



The Moho beneath the Alps



Stehly et al. GJI 2009

Shapiro et al. Science 2005.

Ambient noise imaging complements the traditional approaches.

It is particularly useful for the surface wave tomography of the crust and upper mantle.

Can we use the noise for body wave imaging as well?

The smaller is the contribution of a specific arrival to the Earth response, the more difficult is its reconstruction.

It was worth trying....

The search for body waves : crustal propagation (e.g. Draganov et al., 2009; Zhang et al., 2010, Poli et al.,

#### → Earth's mantle discontinuities from ambient seismic noise



#### Monitoring the elastic properties of the rocks



Direct waves are sensitive to noise source distribution (relative errors small enough for tomography ( $\leq 1\%$ ) but too large for monitoring (goal  $\approx 10^{-4}$ )

Stability of the 'coda' of the noise correlations = frozen distribution of scatterers

We can construct virtual seismograms between stations pairs from noise records.

They contain the information about structures, but also all the complexity of actual seismograms (e.g. Weaver and Lobkis 2004)



Specifically they contain the scattered waves (coda waves). This is attested by the fact that we can also construct 'virtual' seismograms from the correlation of noise based virtual seismograms

 $\rightarrow$  C<sup>3</sup> method (Stehly et al., 2008; Garnier et al., 2011)

 $\rightarrow$  can even be iterated in C<sup>5</sup>.. (Froment et al., 2011)

→ long travel times = strong

sensitivity to changes

Detecting a change of seismic speed: coda waves

Comparing a trace with a reference under the assumption of an homogeneous change



The 'doublet' method: moving window cross spectral analysis (Poupinet et al., 1984; Snieder 2002,

Variation induced by the Parkfield earthquake



Alternatively, the stretching method is a global optimization technique

#### Application Piton de la Fournaise





Application to Parkfield (*Brenguier et al. 2008*) Short period sensors / Processing in the period 1-10s





**HRSN** network

→ GPS trend→ tremor activity