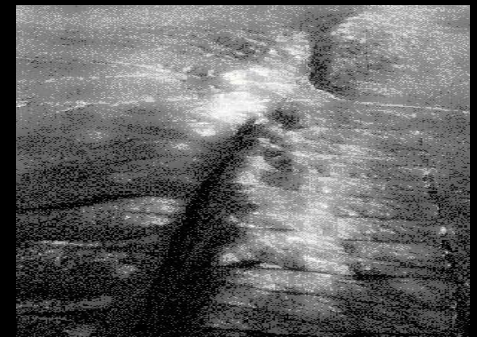
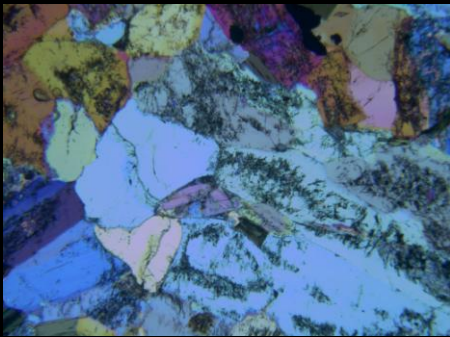


# Nonlinear elastic and plastic behavior of dense granular systems

Paul Johnson



# Mixed nonlinear dynamics

# associated information



## International Conference on Nonlinear Elasticity in Materials (ICNEM)

Santa Fe, New Mexico USA, July 2-7, 2017

<http://icnem.org/index.html>

## International Symposium on Nonlinear Acoustics (ISNA), Santa Fe, New Mexico July 8-13, 2018

LANL group Home page with links to publications:

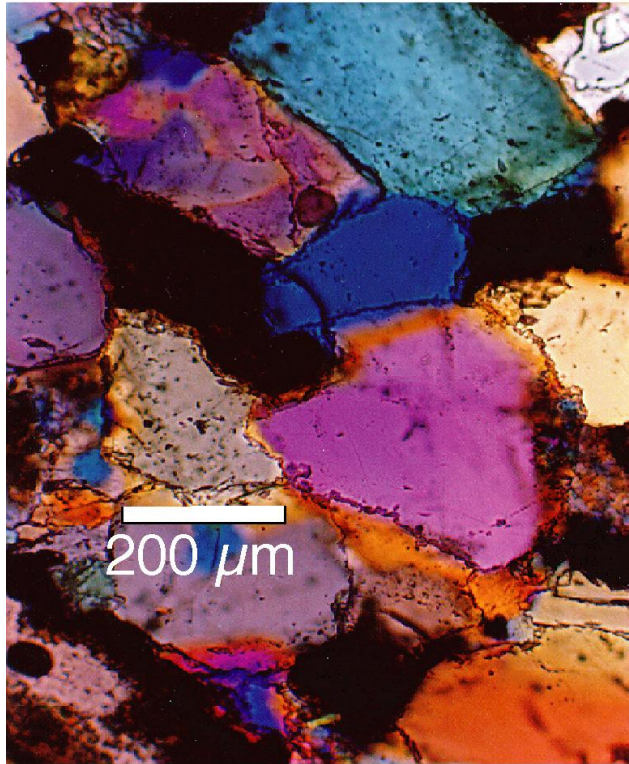
<http://geophys.lanl.gov/nonlinear/nonlinear.shtml>

## Contents

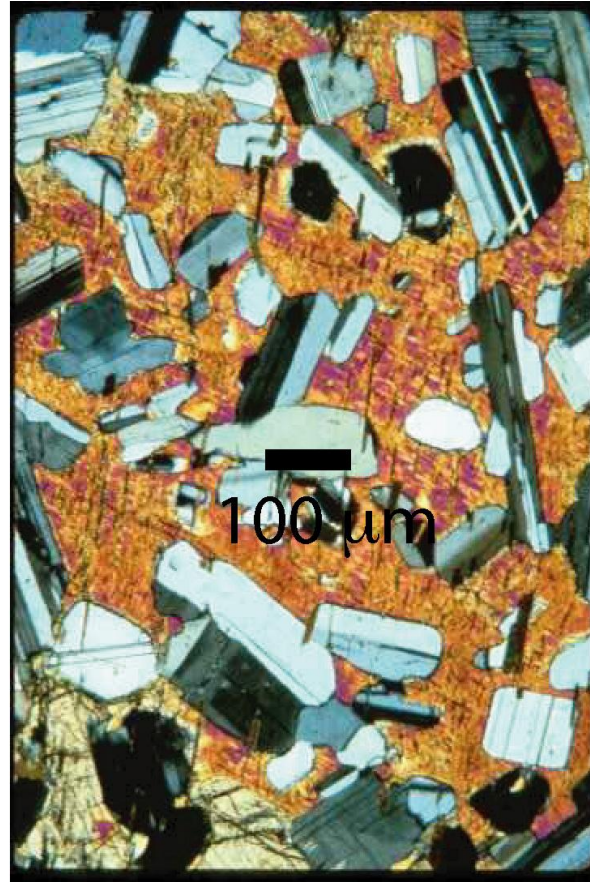
- Universal concept of nonequilibrium behavior
- Background—quasistatic and dynamic elasticity
- Nonequilibrium behavior in earth materials  
Simulation—soft modes?
- Theory
- Conclude



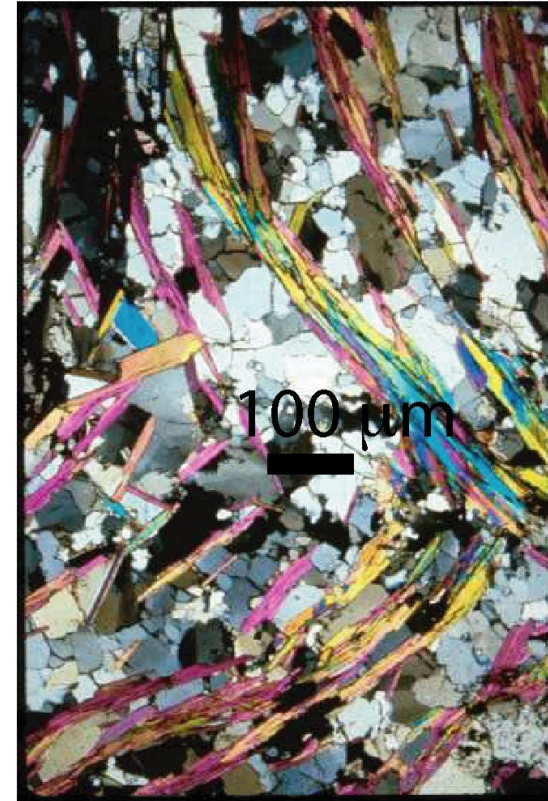
Rocks are complex granular materials, 'sintered' together



Sandstone

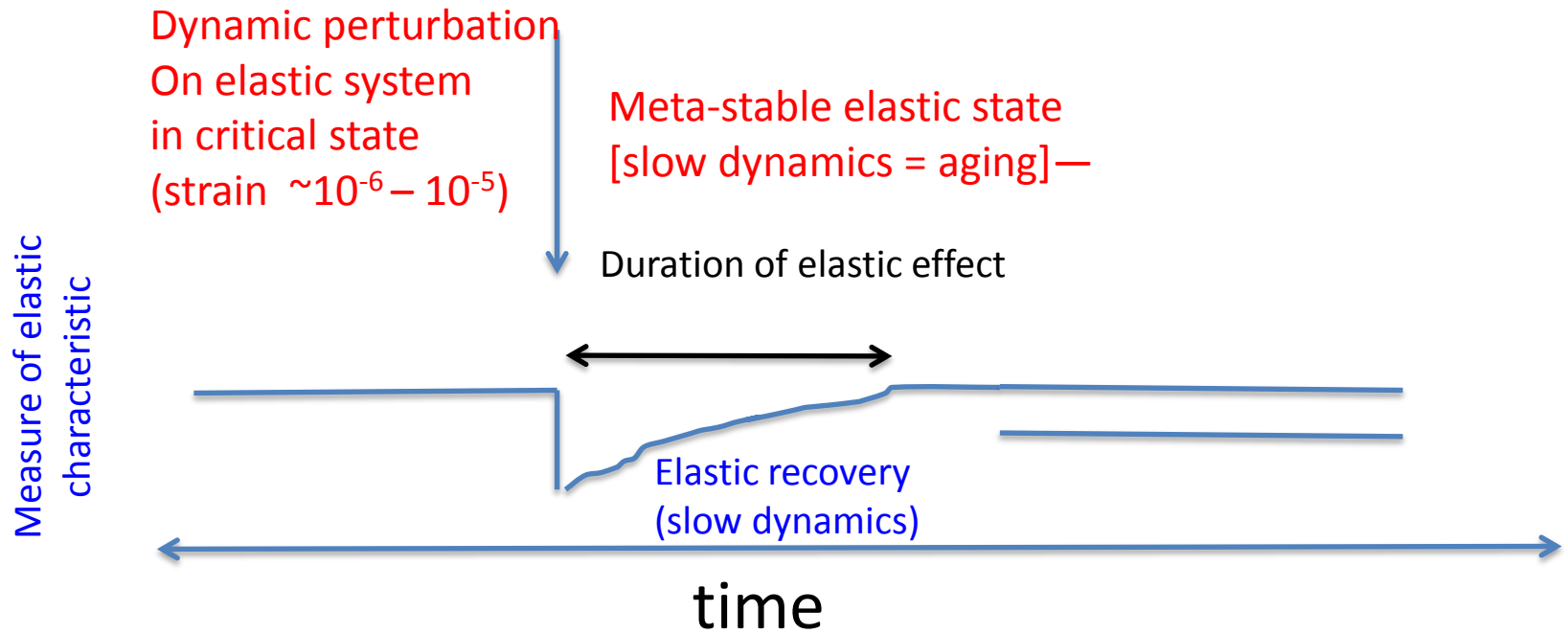


Anorthosite



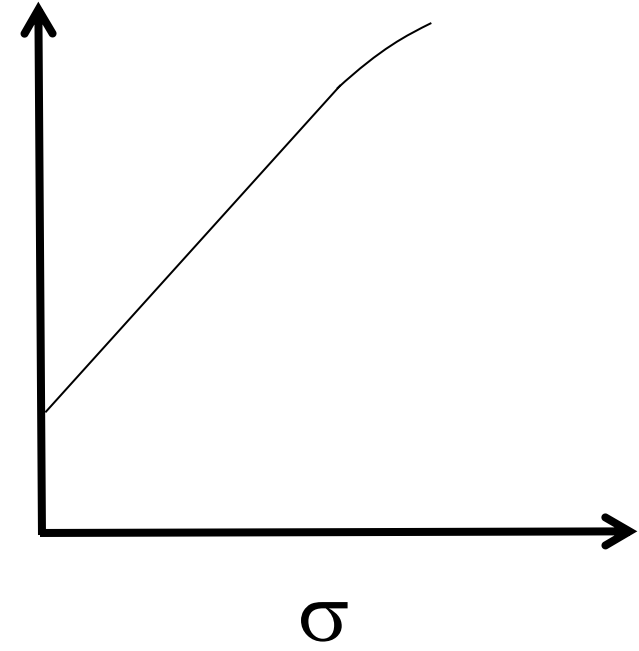
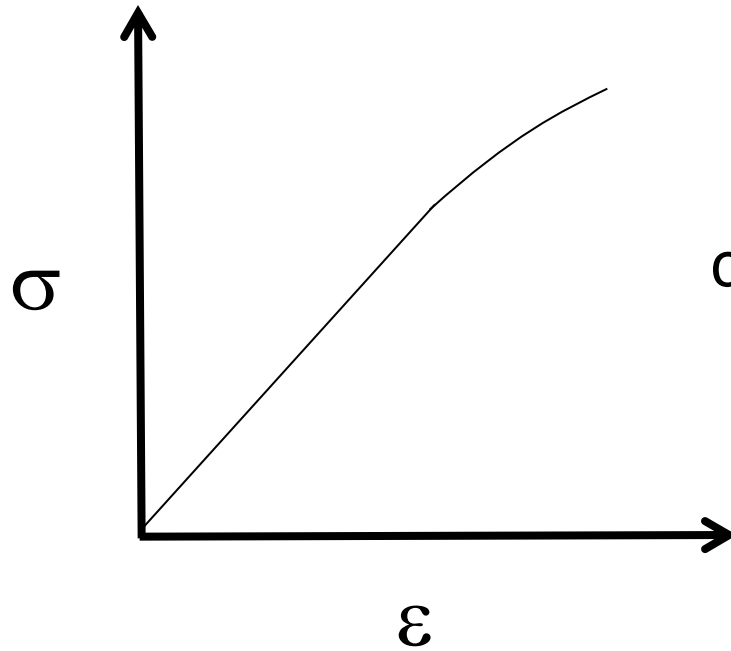
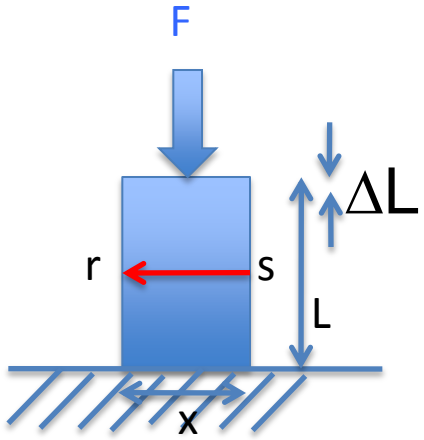
Quartz-mica schist

# All exhibit nonlinear fast and slow dynamics



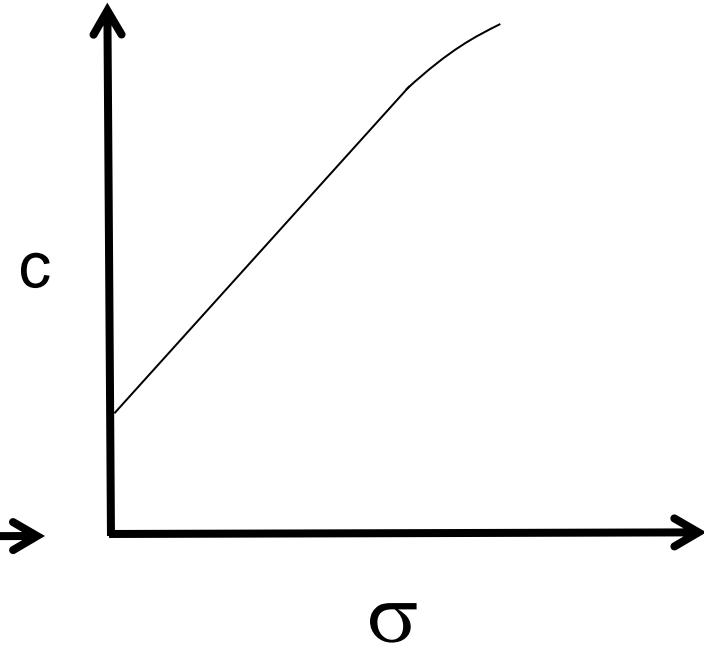
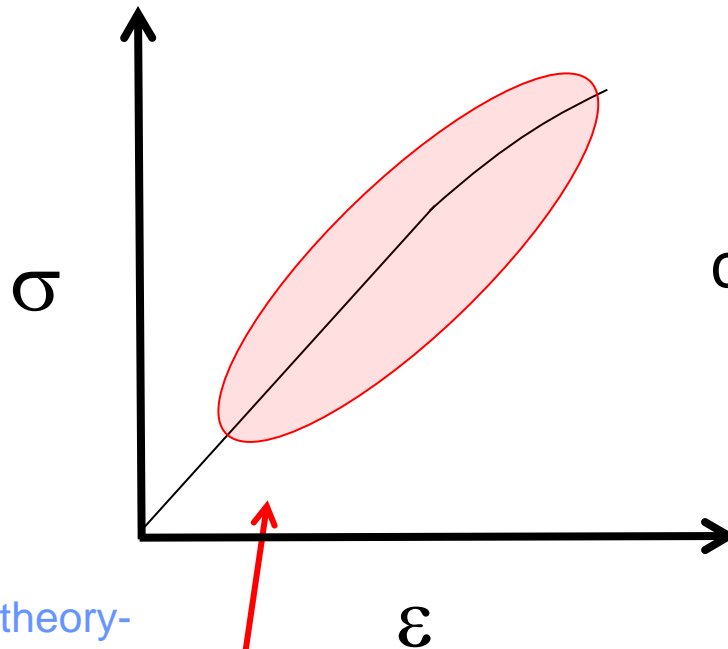
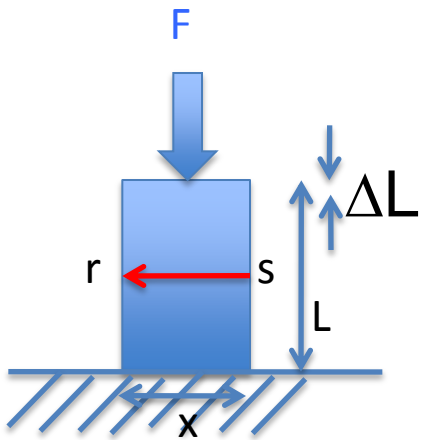
Quasi-static and dynamic nonlinear  
elasticity in consolidated granular materials  
(e.g., rock, concrete...)

# Quasistatic and dynamic elasticity of individual crystals and many metals





# Quasistatic and dynamic elasticity of individual crystals and many metals



Thermodynamics based theory-  
Nonlinearity due to anharmonicity  
in the crystalline lattice.

$$S = Me[(1 + be)]$$

isotropic

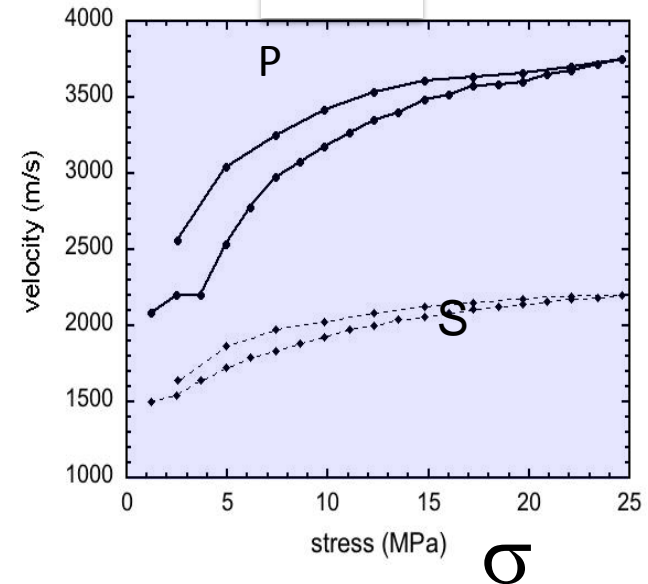
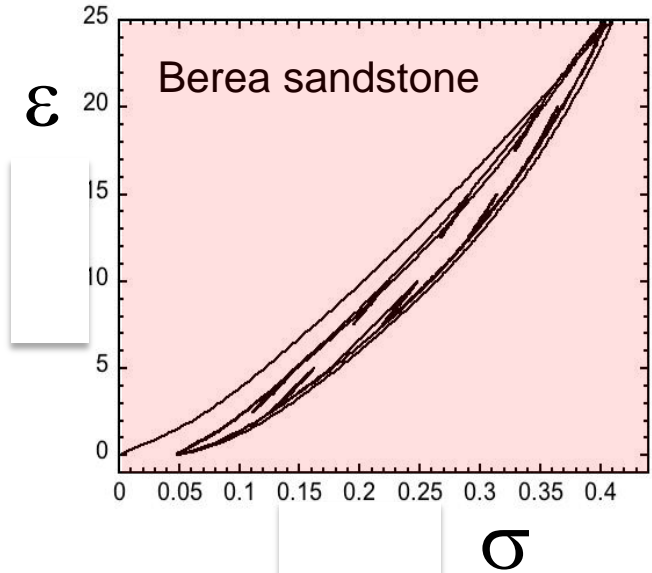
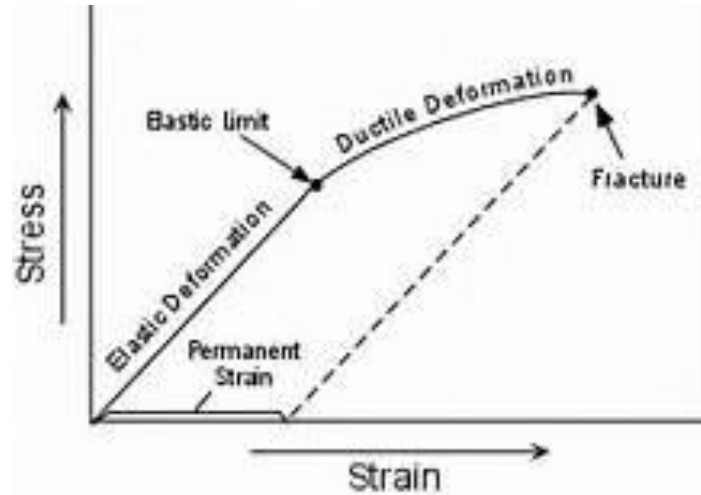
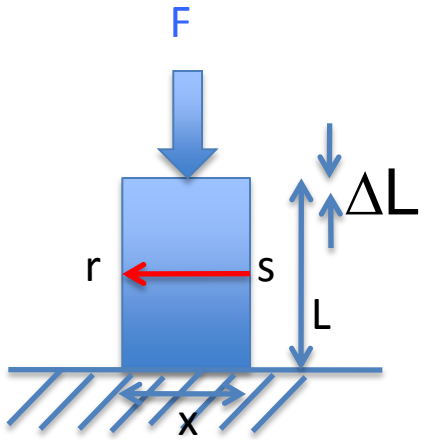
$$b = \frac{3}{4} + \frac{A+B+C}{2rc^2} = \frac{3}{2} + \frac{l+2m}{l+2u}$$

Goldberg

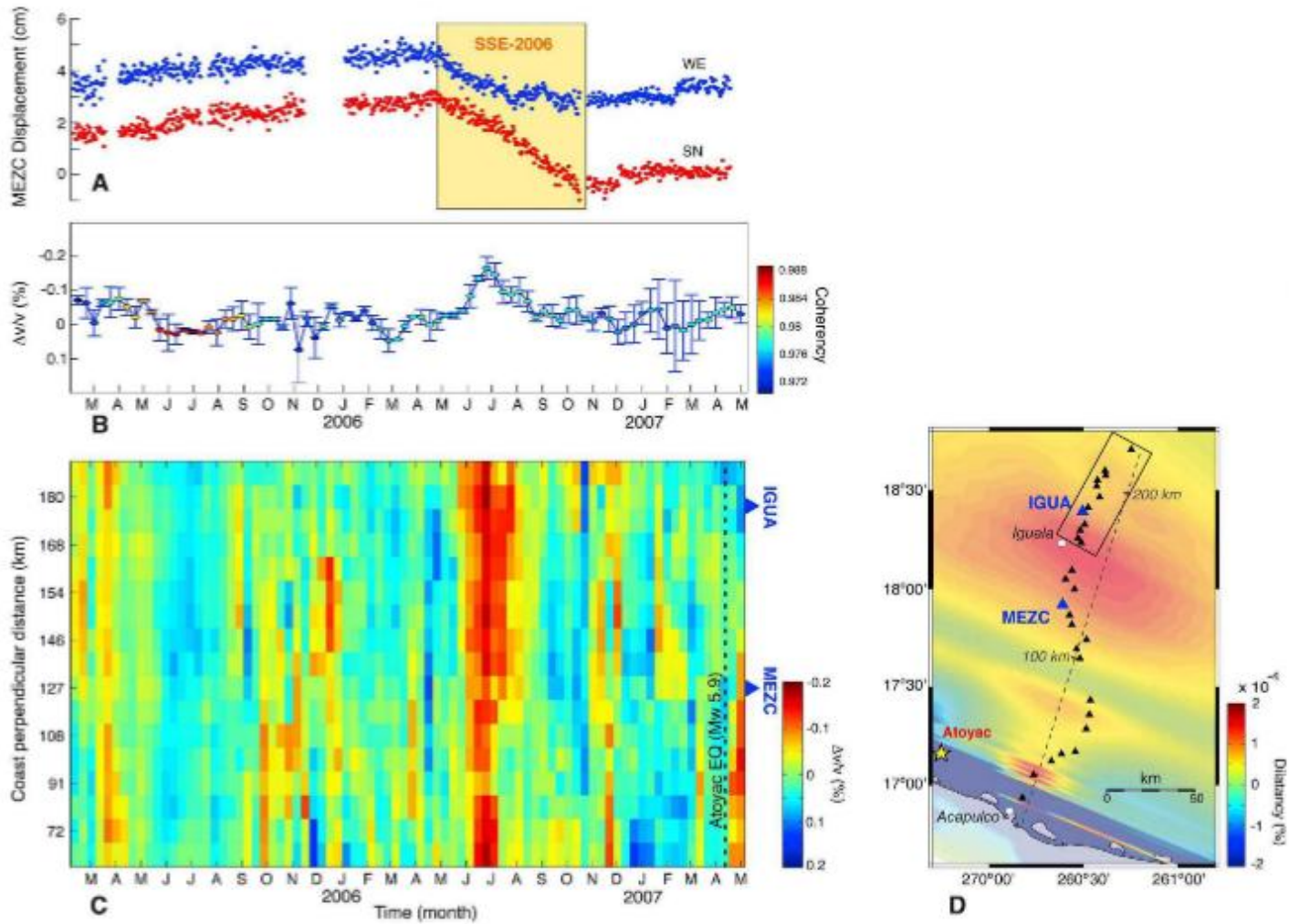
Murnaghan

# Quasistatic elasticity of rock and 'damaged' materials

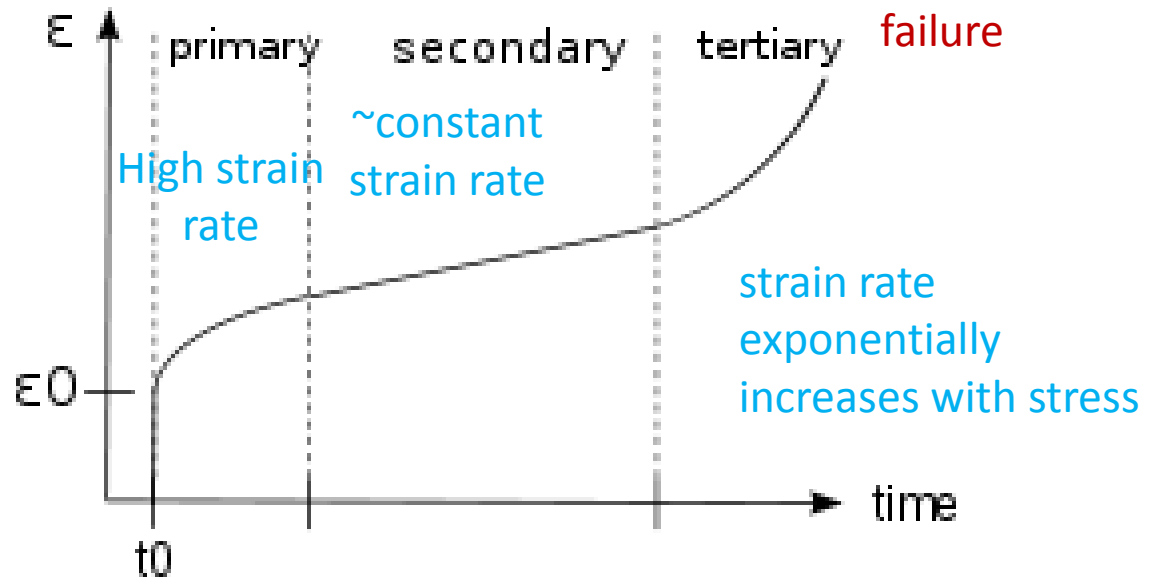
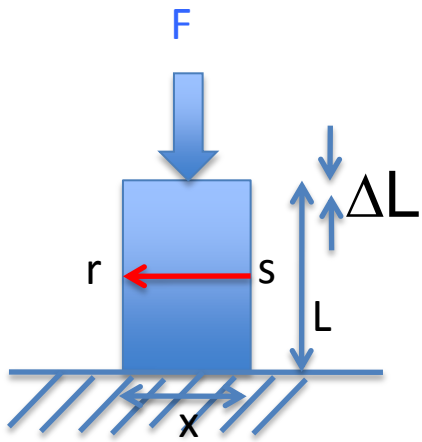
Nonlinearity much larger than metals and single crystals.....



# Static forcing of slow slip in Mexico

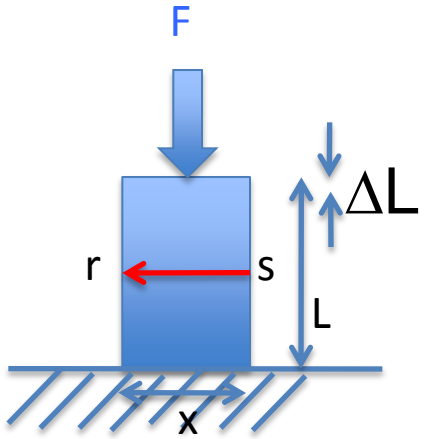


# Creep under static loading

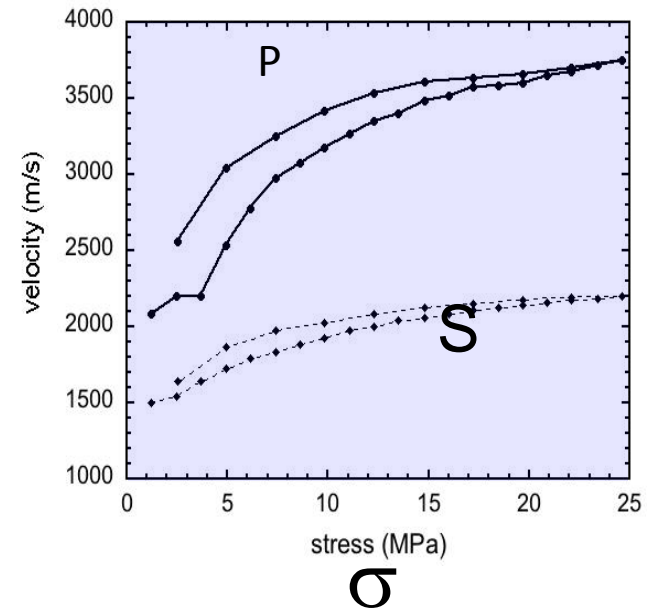
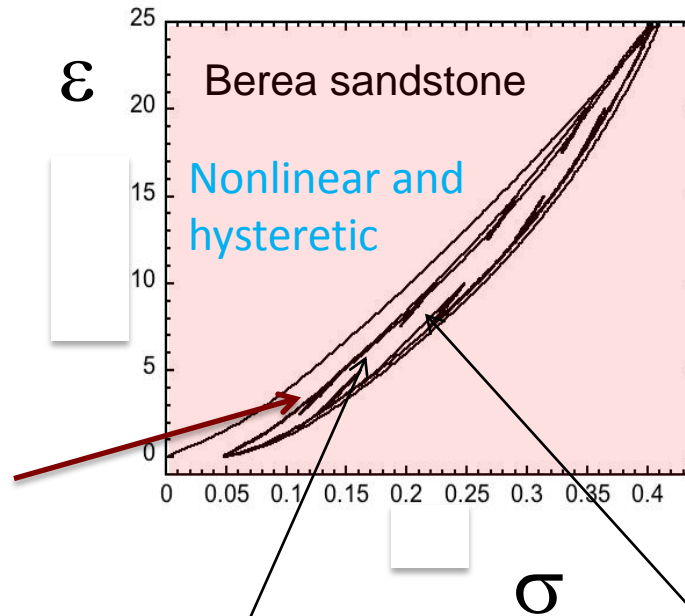


# Elasticity of individual rock and damaged materials

Nonlinearity much larger than metals and single crystals.....



End point memory



Phenomenological model of nonlinearity and hysteresis (Preisach-Mayergoyz)

$$s = K_0 \left( e + be^2 - a \left[ (De)e + \mathbf{sign}(\dot{e}) \frac{e^2 - (De)^2}{2} \right] \right)$$

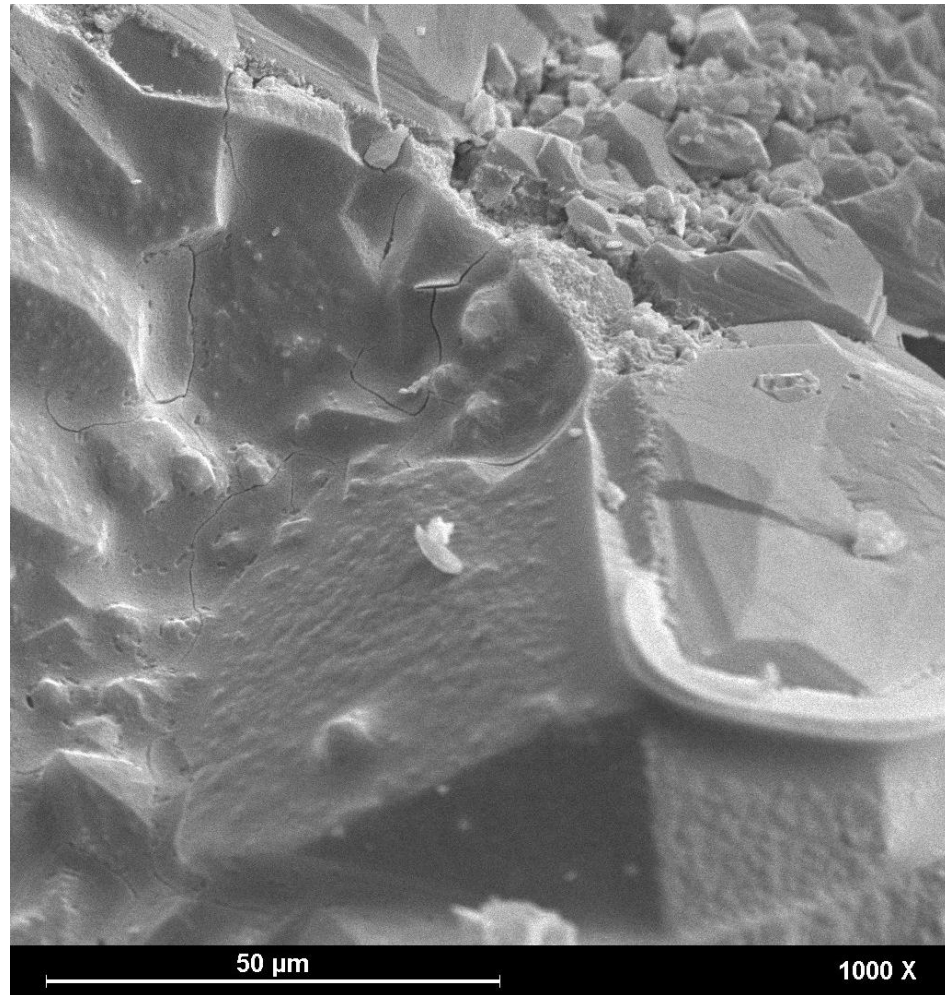
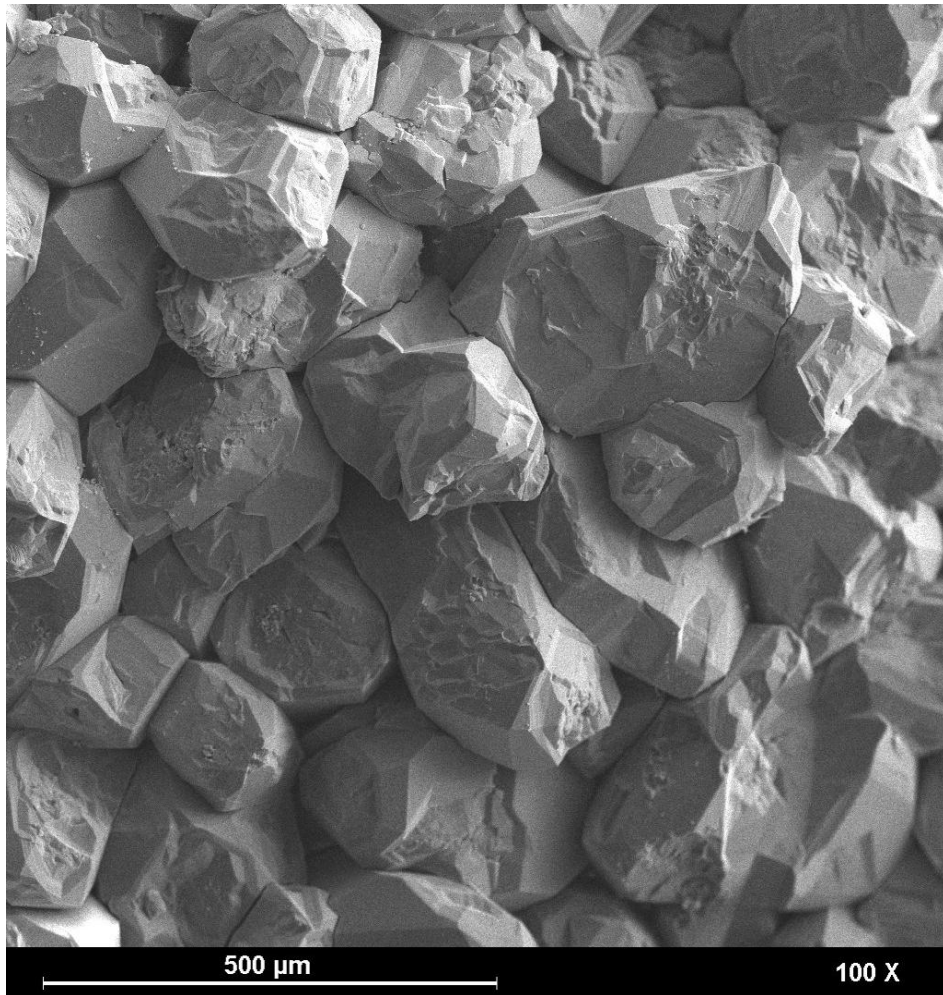
$$e = \frac{du}{dx}, b = \frac{3(l+2m) + 2(l+2m)}{2(l+2m)}$$

M=modulus,  $\epsilon$ = strain,  $\lambda, \mu$  = Lamé coeffs. l, m = 3<sup>rd</sup> order [Murnaghan] coefficients



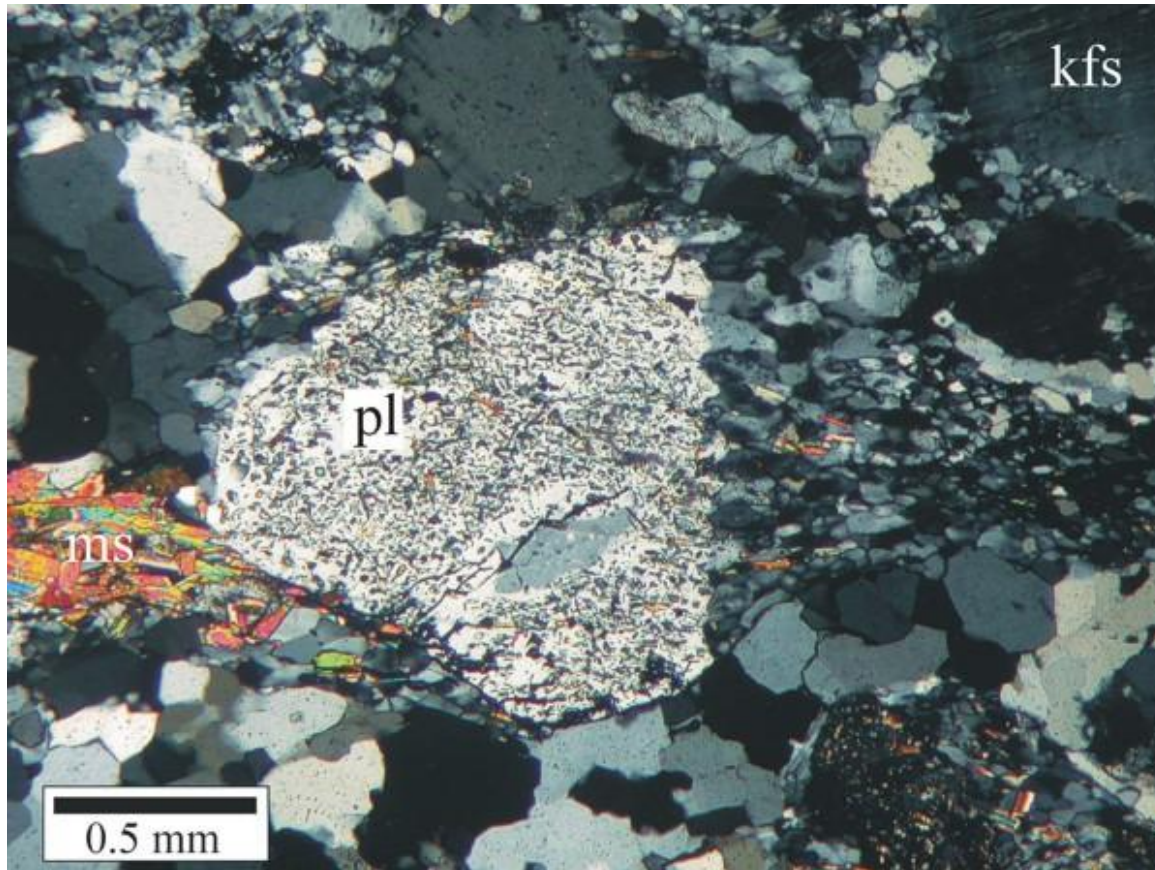
# Why the velocity (and dissipation) change with pressure?

Fontainebleau sandstone (SEM)





# Why the velocity (and dissipation) change with pressure?



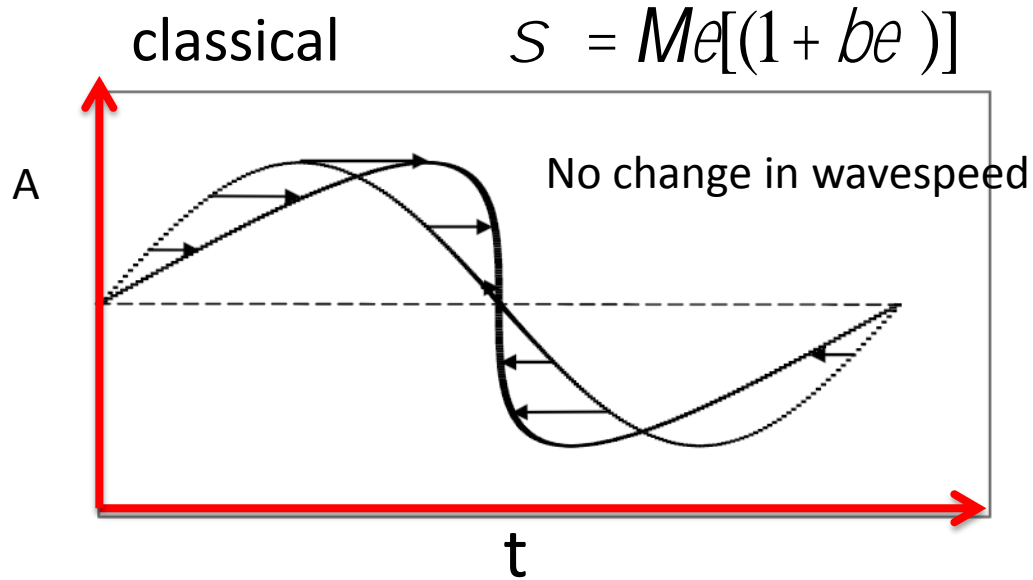
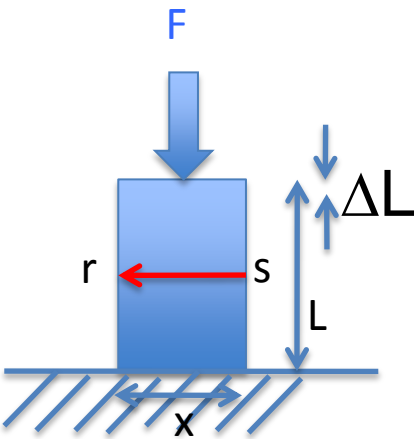
Granite in  
Thin-section

Pressure makes rock  
look like ordinary solids,  
closing cracks

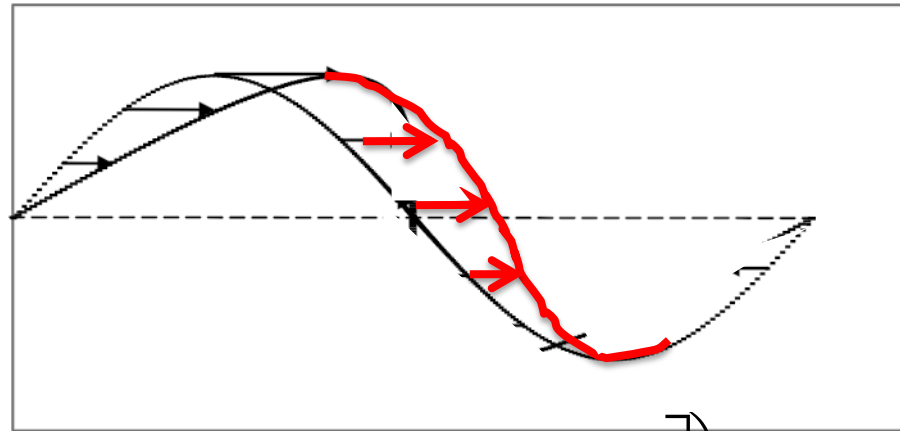
kfs: K-feldspar  
ms: muscovite  
pl: plagioclase

# Dynamic nonlinear elasticity in a traveling wave

Fix the load and measure elastic wave



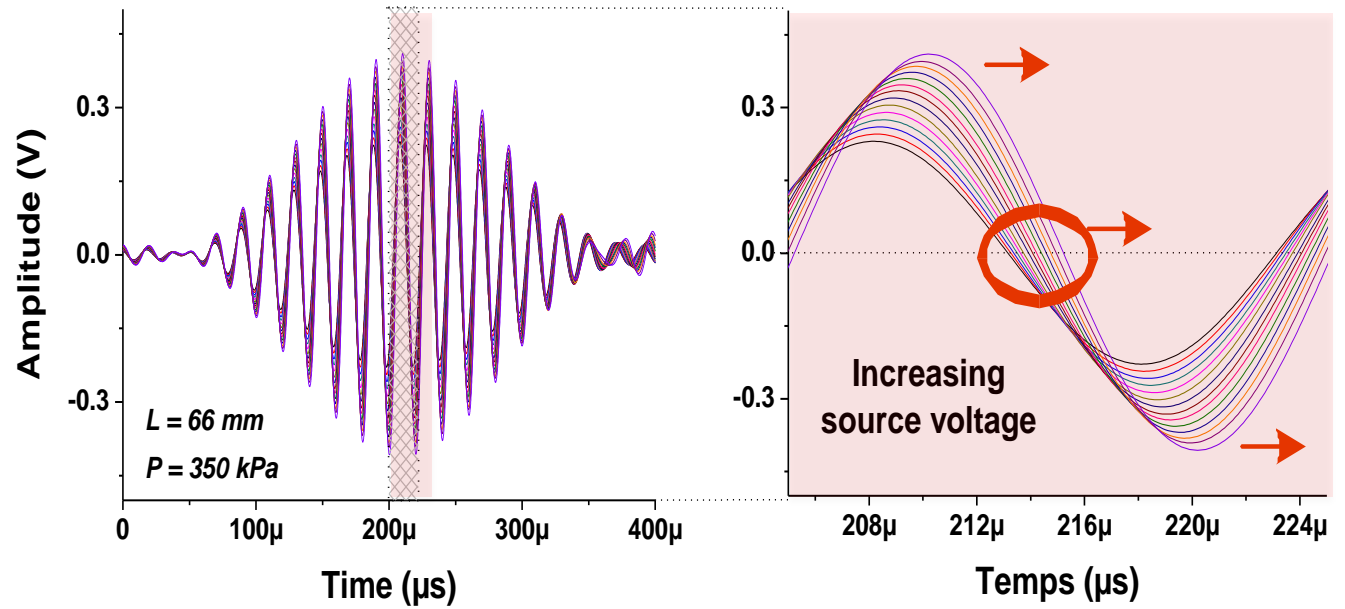
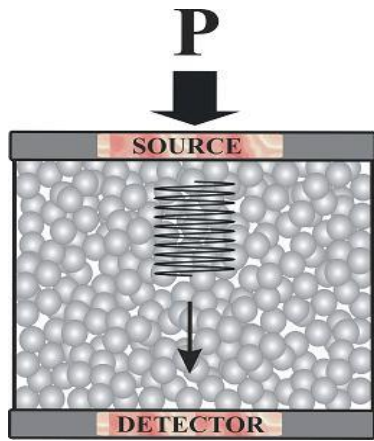
Single crystals,  
many  
metals....



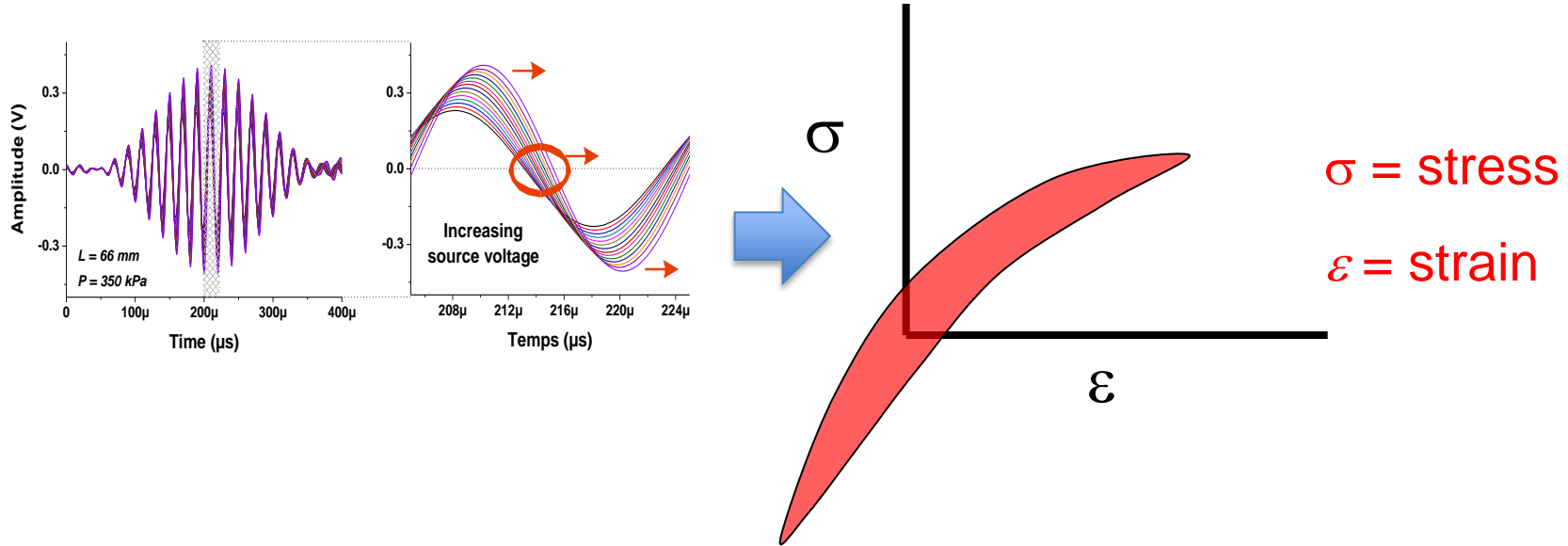
$$s = K_0 \left( e + be^2 - a \left[ (De)e + \text{sign}(\dot{e}) \frac{e^2 - (De)^2}{2} \right] \right)$$

Nonclassical  
(rock, GM....)

# nonlinear acoustics in a traveling wave- example glass bead pack



# Wave dynamics in rock and other earth materials



Like in quasi-static experiments, nonlinearity and hysteresis are present in dynamics when strains exceed  $\sim 5 \times 10^{-7} - 10^{-6}$  [at low effective pressure],  
*Material dependent*

# 1-D phenomenological wave equation

Normal form:

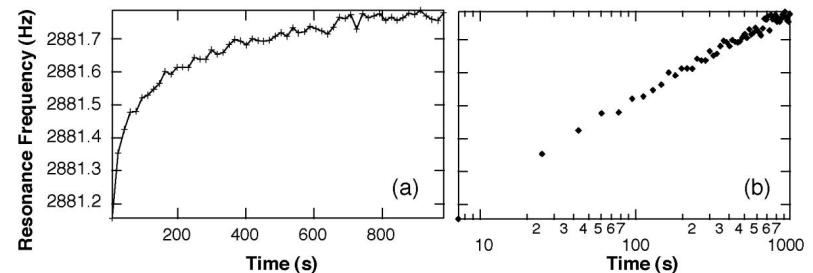
$$\rho \frac{\partial^2 u}{\partial t^2} = \frac{\partial S}{\partial x} + F$$

substituting the nonlinear stress relation

$$\rho \frac{\partial^2 u}{\partial t^2} = \frac{\partial K_0 \left( e + be^2 - a \left[ (De)e + \text{sign}(\dot{e}) \frac{e^2 - (De)^2}{2} \right] \right)}{\partial x} + F$$

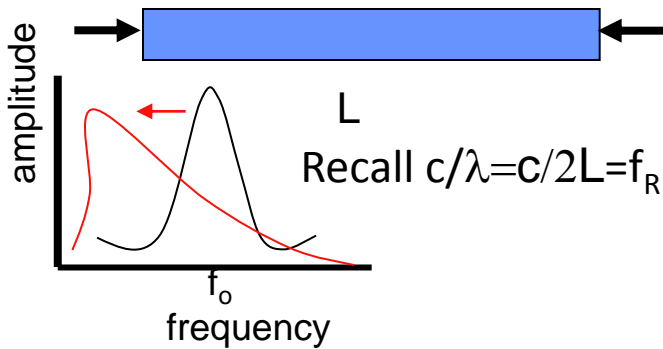
## Problems:

- i. Phenomenologic
- ii. Contains no rate dependence



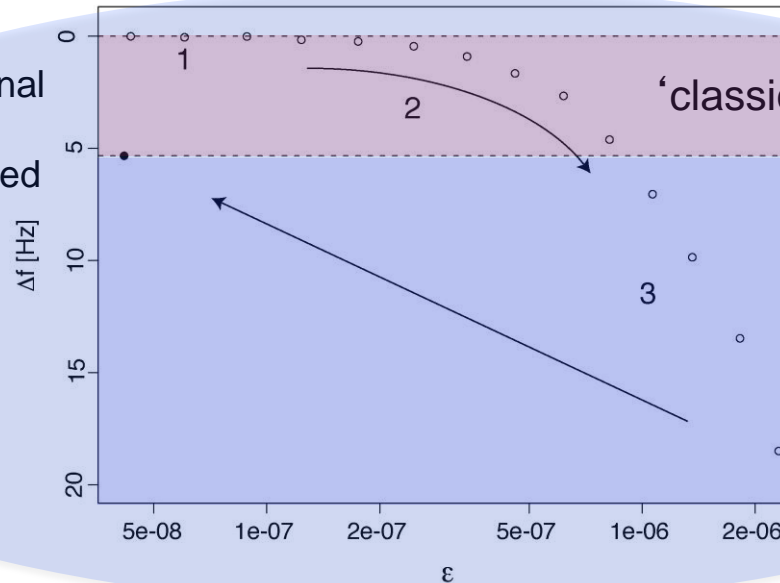
# Exploring regimes of elastic nonlinearity applying wave resonance

Young's mode experiments in a bar



(sandstone)

$\Delta f$   
proportional  
to  
wavespeed

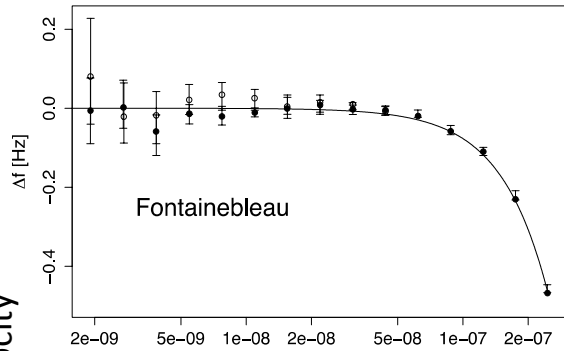


b.  
'classical' behavior

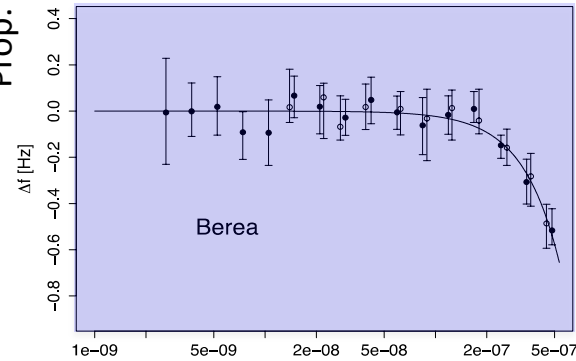
Nonclassical,  
nonequilibrium



# Recovery– the slow dynamics (*aging*)

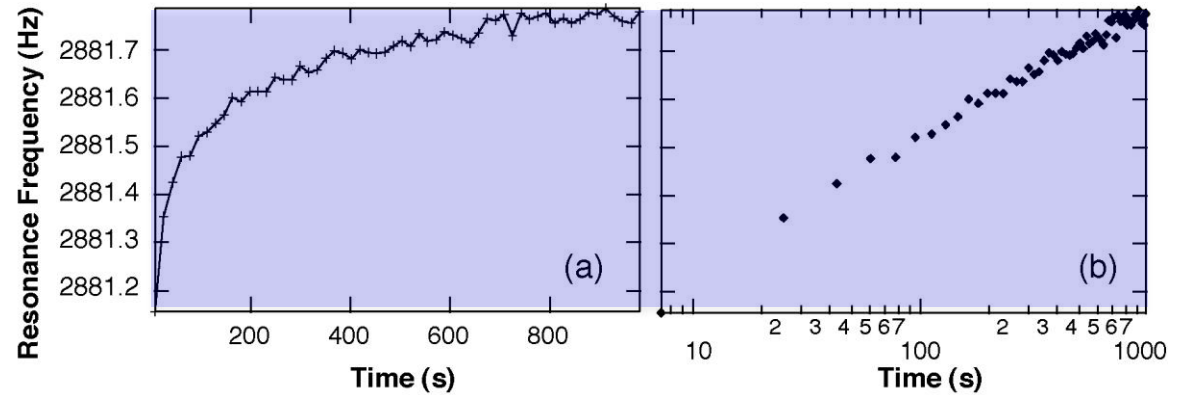


strain



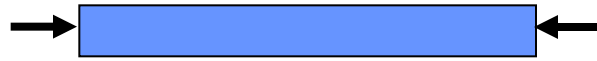
strain

## Recovery in Berea sandstone



*Linear with the log of time--*

# Nonlinear, nonequilibrium behavior in diverse materials

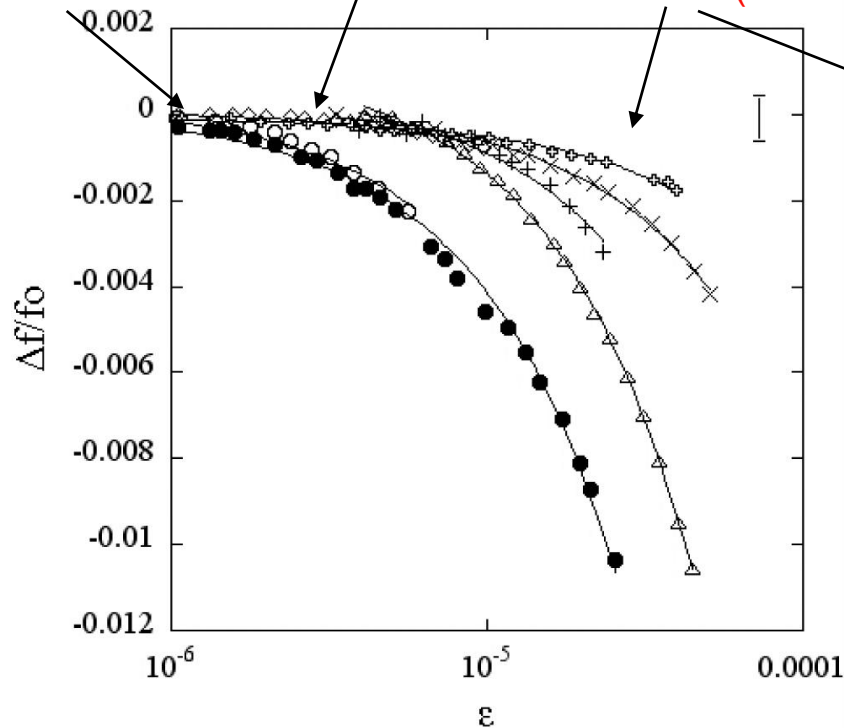


Young's mode experiments

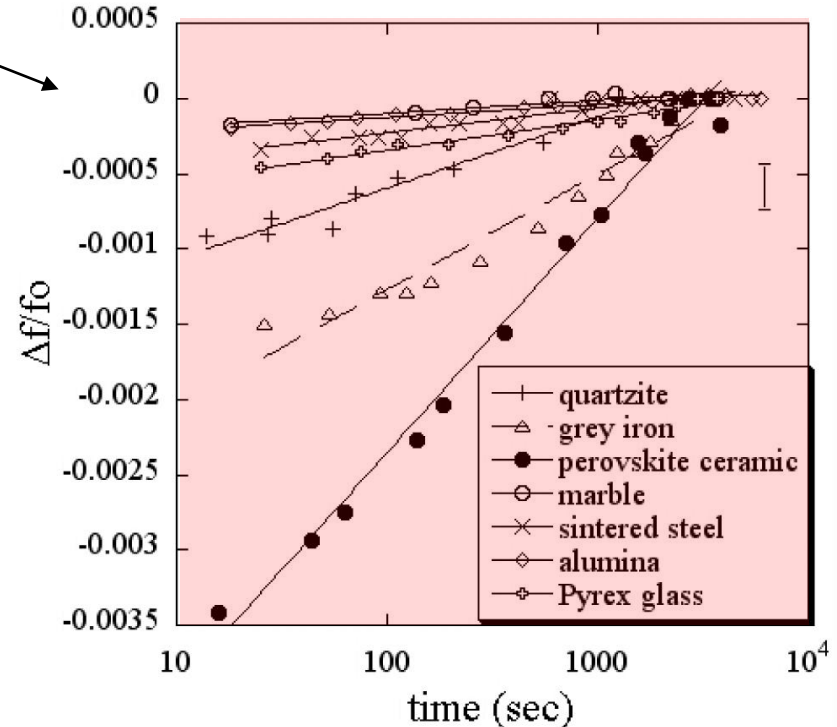
Classical NL behavior

Nonclassical (nonequilibrium) behavior

Linear elastic



Fast dynamics/conditioning

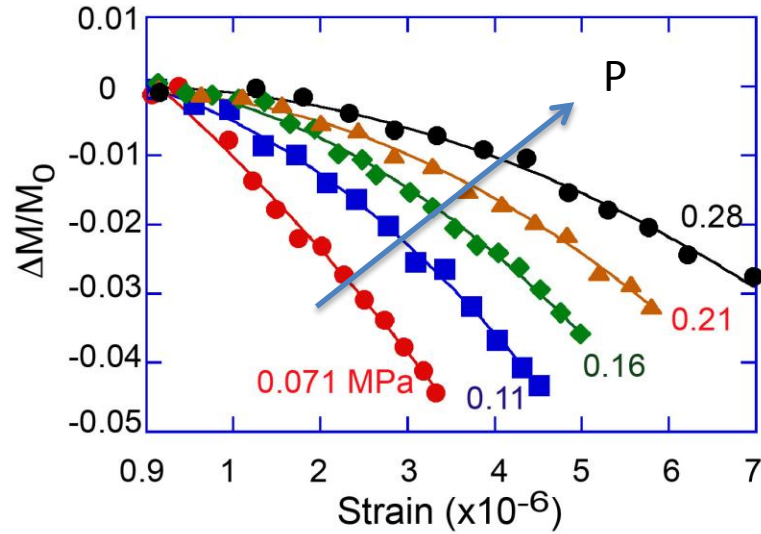
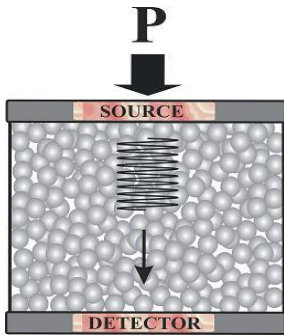


Slow dynamics

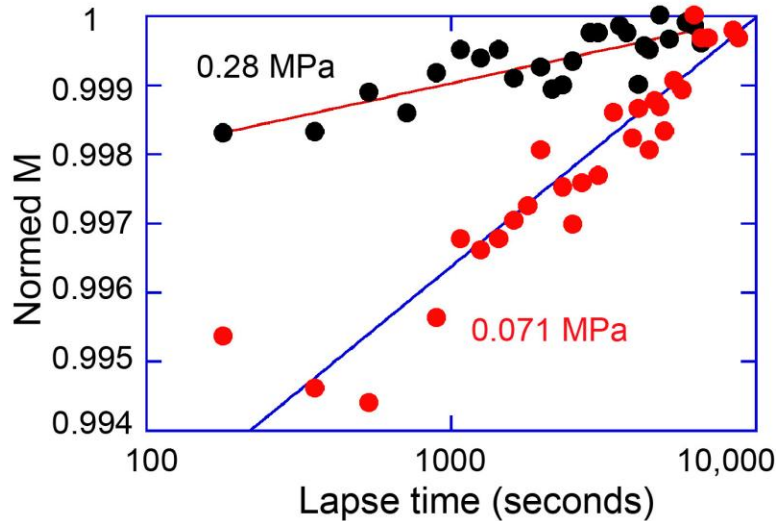
Strong effects on wave attenuation as well.

# with applied pressure?

Experiments with rock and canisters of glass beads



Decrease in modulus and increase in dissipation with wave amplitude.

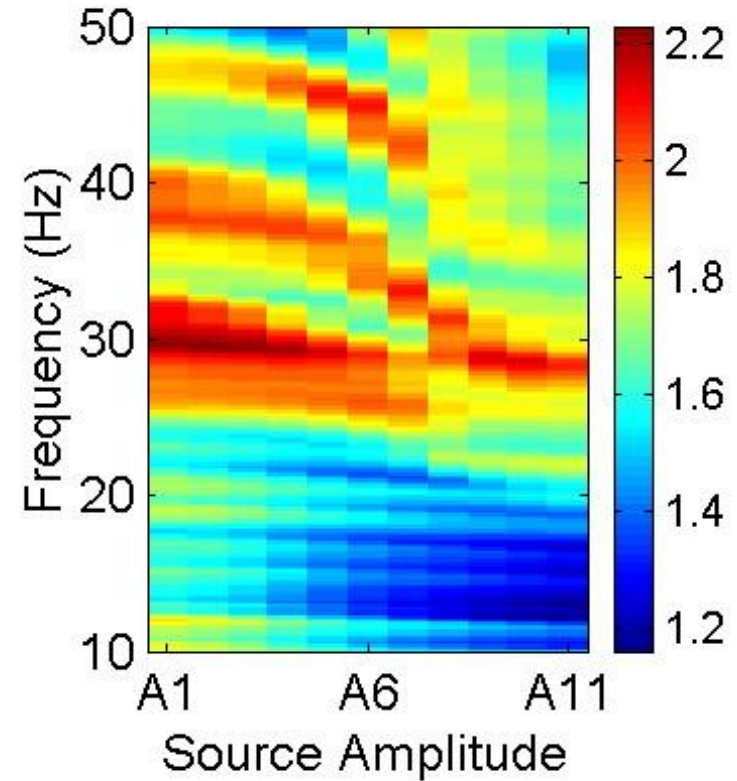


Long recovery to equilibrium (the slow dynamics)

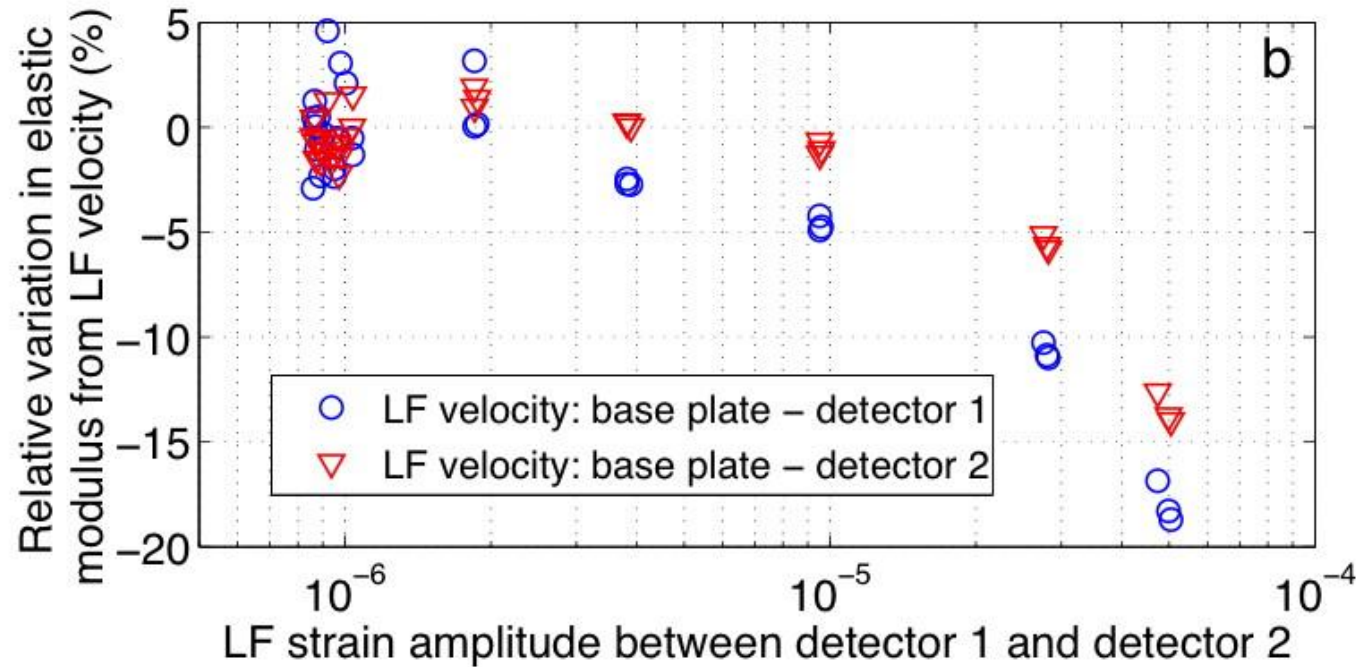
# Induced nonlinear behavior with an active source



T-Rex, University of Texas at Austin

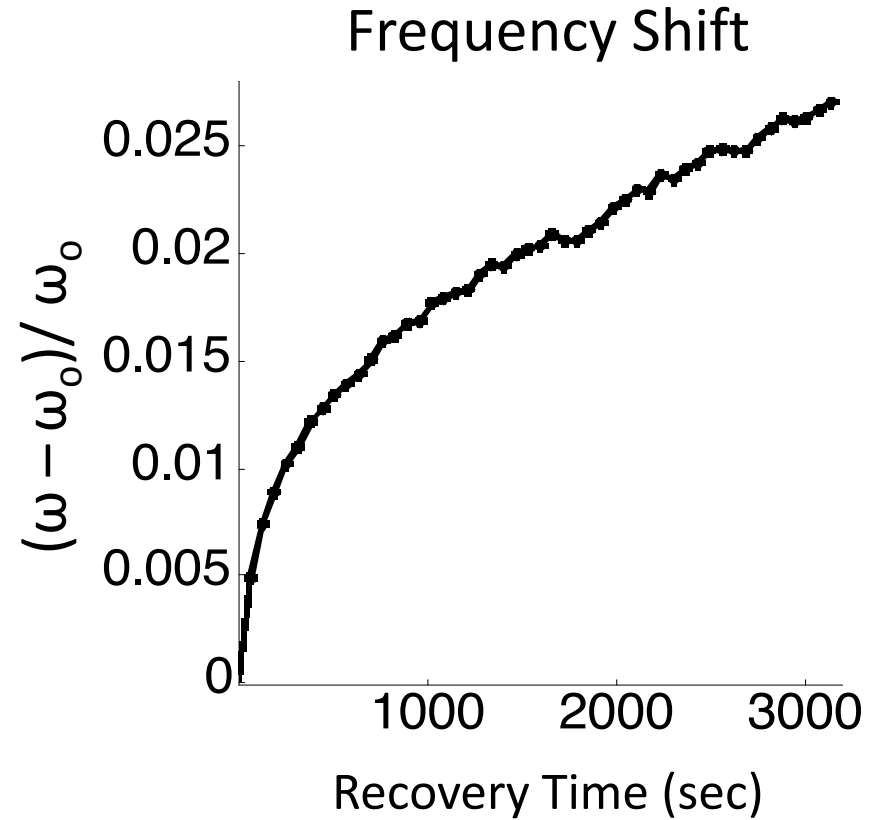
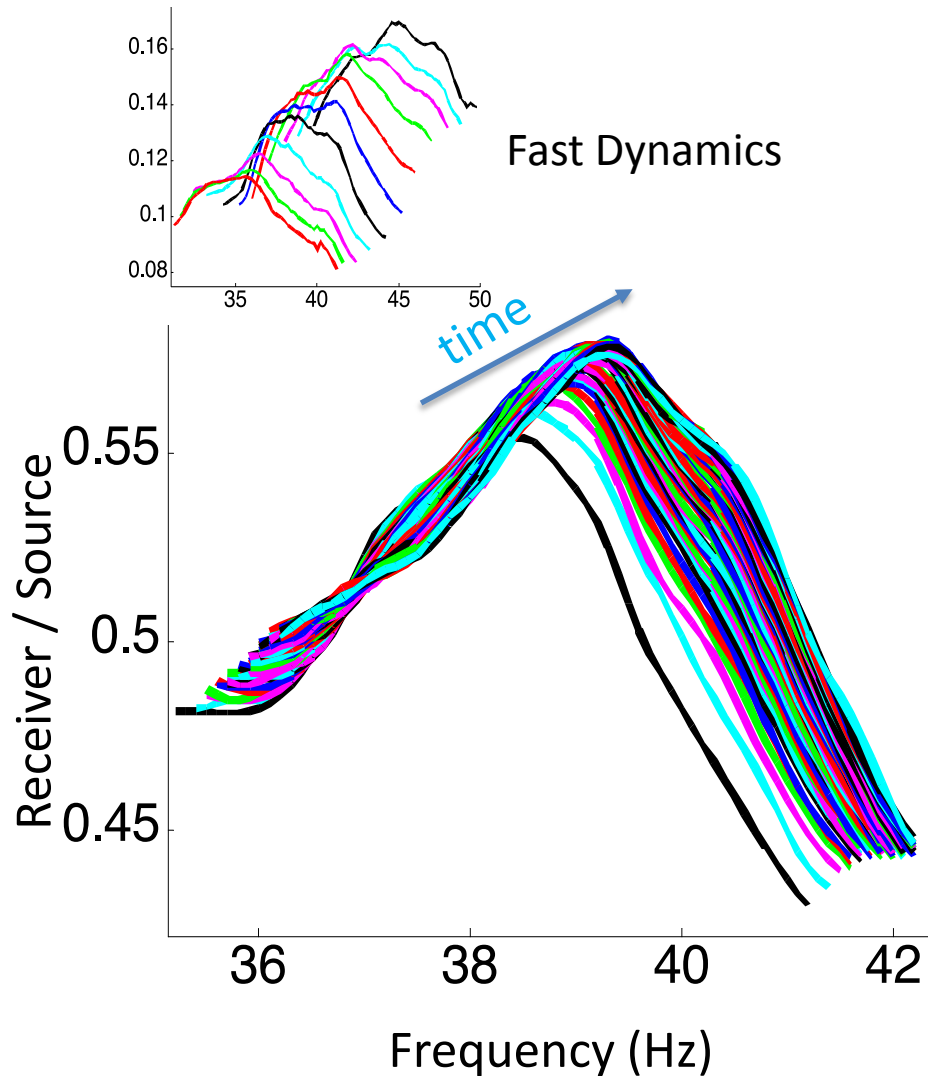


# Near surface sediment response from a large vibrator source: site characterization



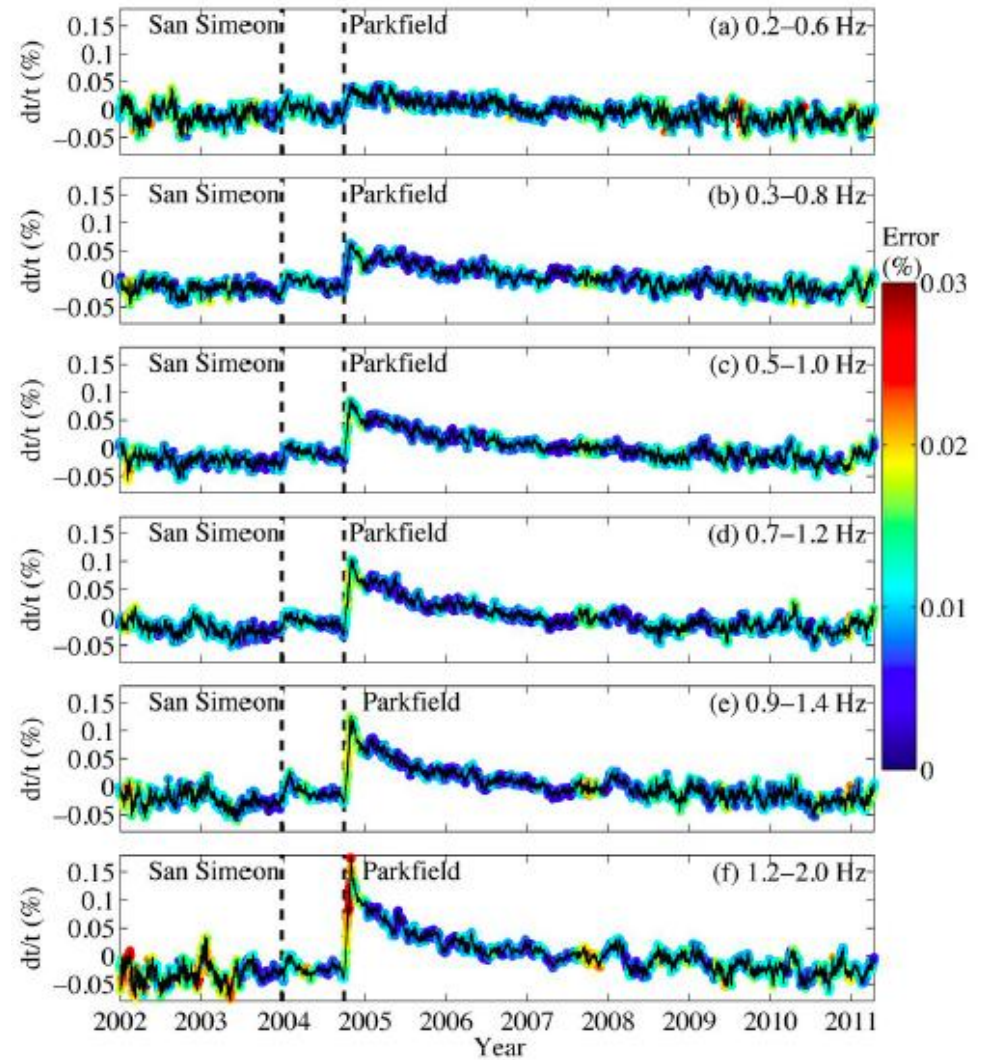
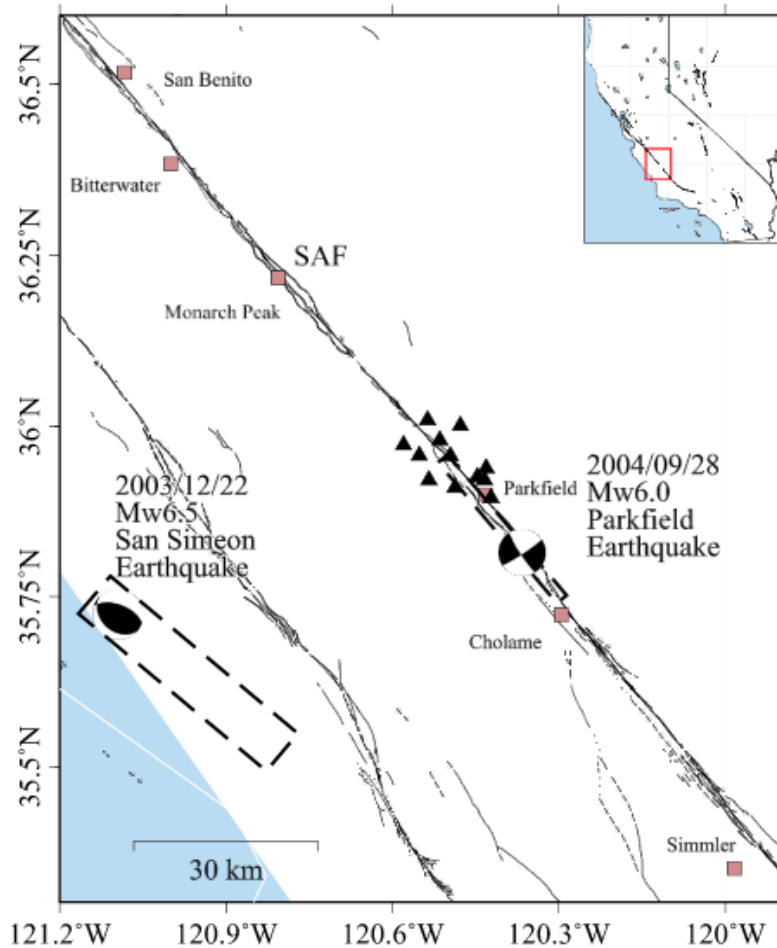


# Slow dynamics

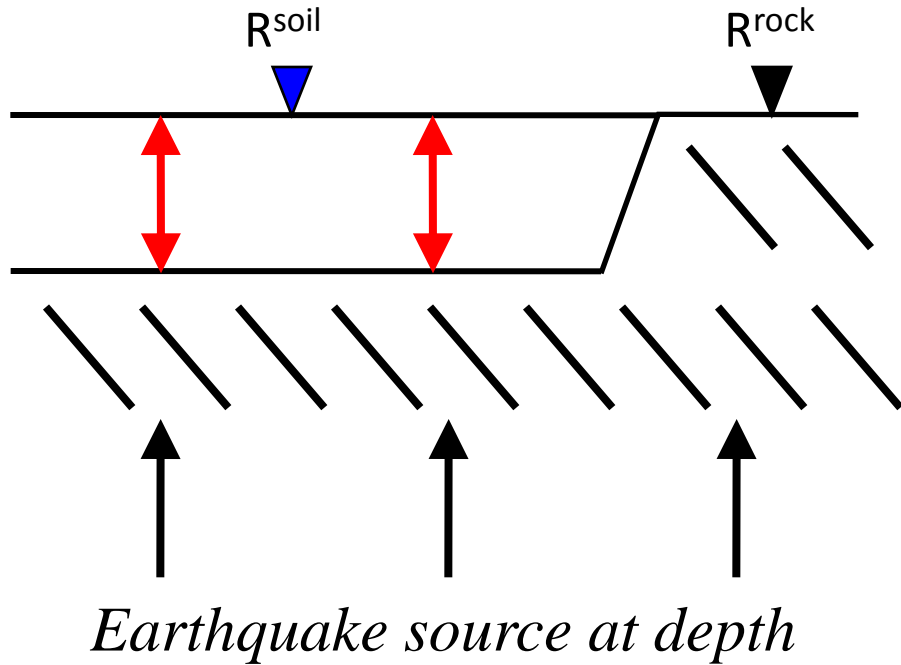




# parkfield

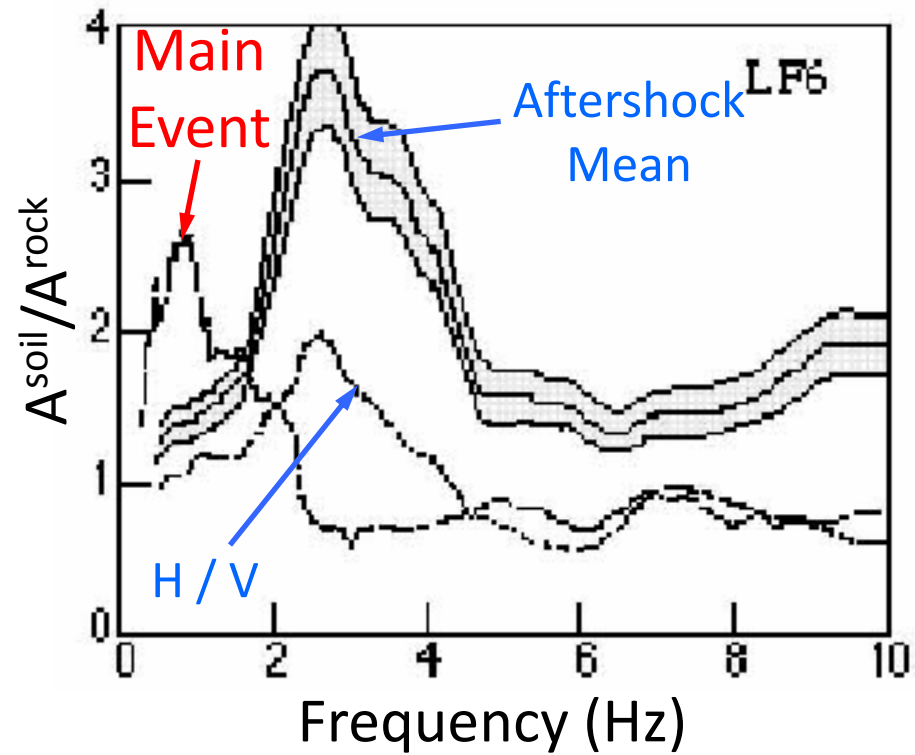


# Strong ground motion



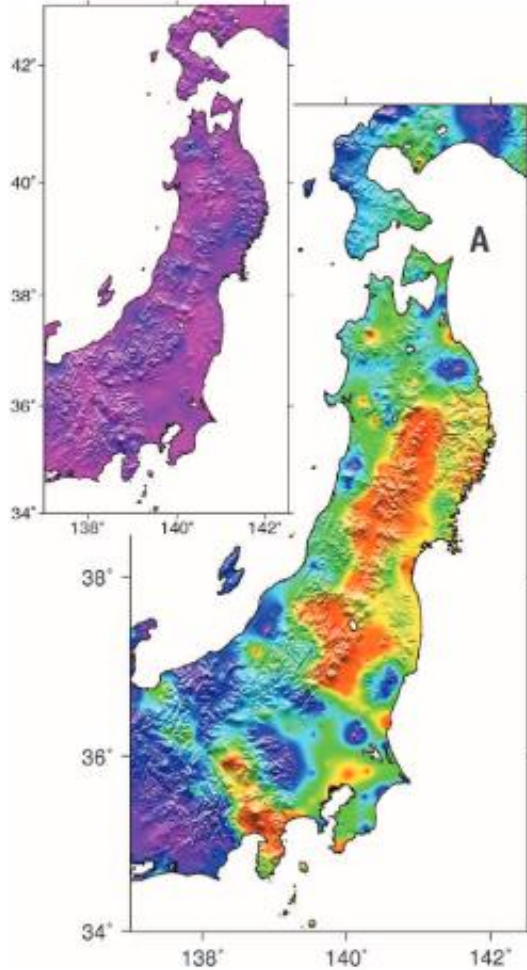
Site Response  $\approx$  
$$\frac{A^{soil}(\omega)}{A^{rock}(\omega)}$$

## 1994 Northridge Earthquake

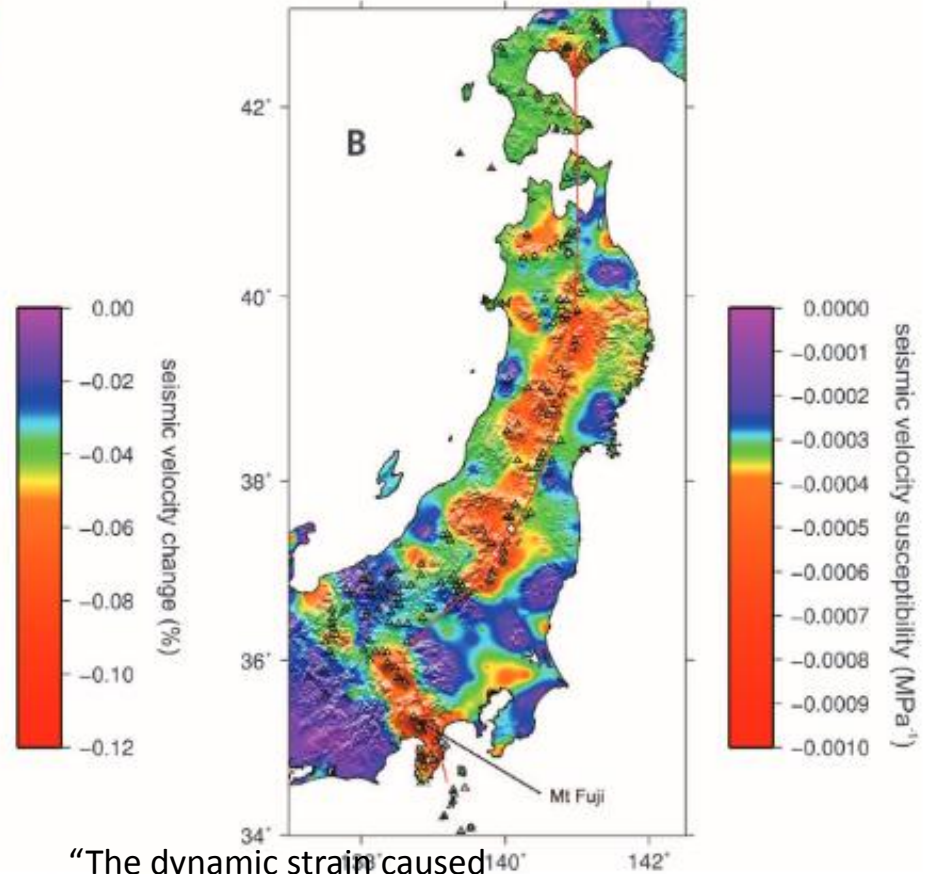


# Quasi-static+ dynamic nonlinear elasticity

Co-seismic velocity change



Velocity susceptibility

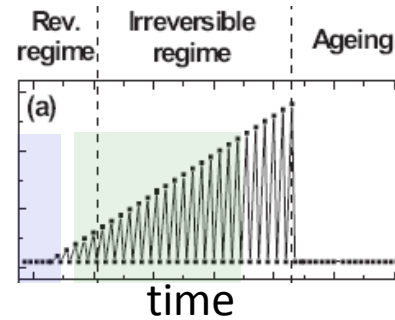
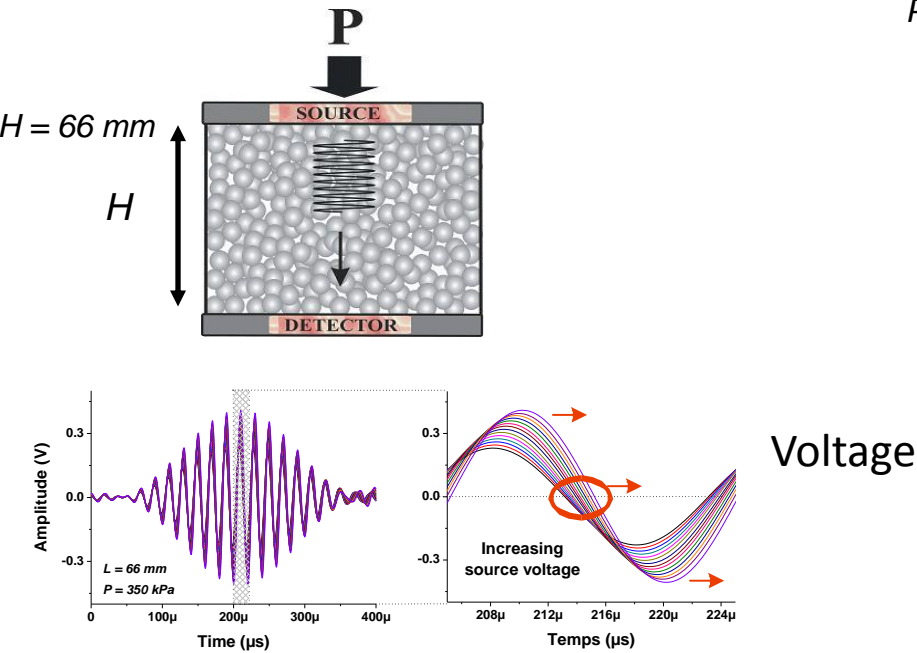


Brenguier et al., Science (2014)

“The dynamic strain caused by the passing of the seismic waves was one to two orders of magnitude higher than the static coseismic strain for Honshu Island. “

# Nonlinear elasticity and plasticity in glass beads

Jia, Brunet, Laurent, *PRE*; 84 (2011); Olson Reichhardt et al., *Physical Review E* (2015); Limrich et al., in review *PRE* 2017



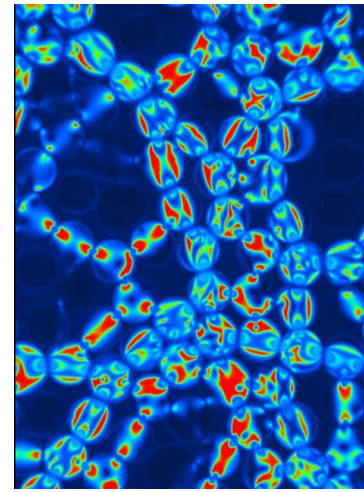
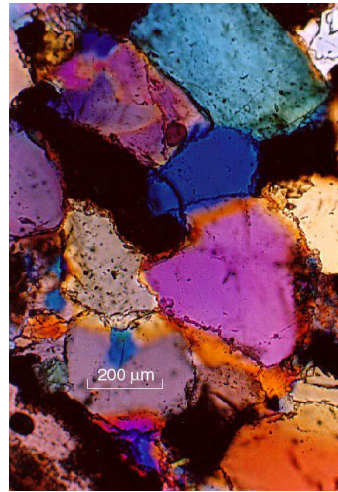
Plasticity is not observed in intact rock at typical dynamic strains, but in highly fractured media, it could be (damage zone in a fault)

# Class of materials exhibiting rock-like “nonclassical” (hysteretic) nonlinear behavior

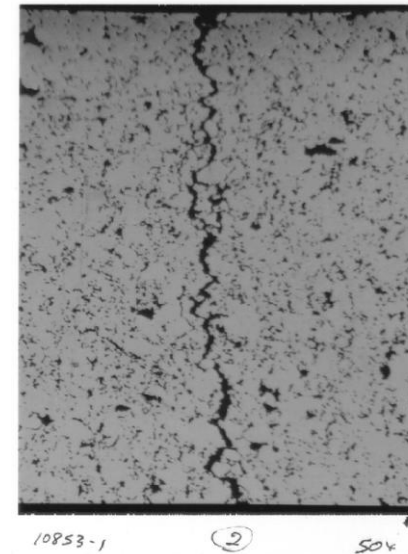
Granular media, some  
ceramics, all(?) rocks,  
all damaged  
materials,  
sintered metal....

A primary  
contribution to  
elastic  
nonlinearity is  
mechanical  
damage.

“Distributed”



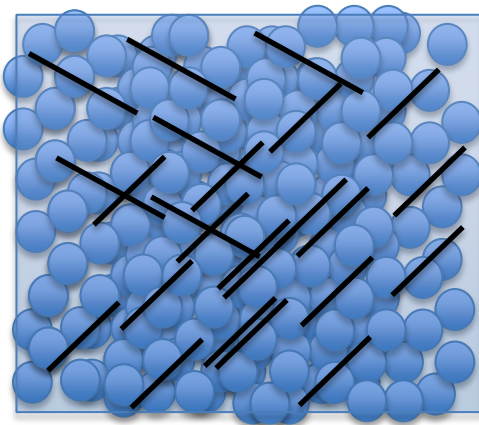
“Localized”





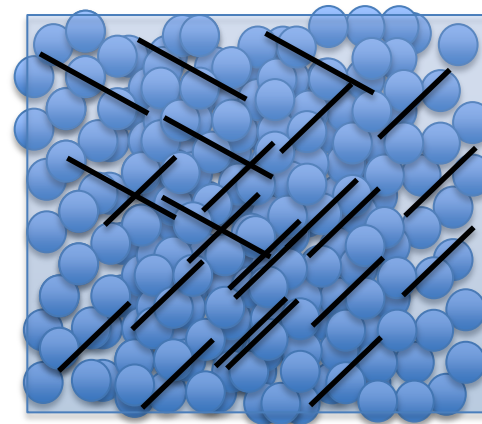
# Mechanism?

$$F \cos(\omega t - kx), \omega < 2\pi$$



Stress transferred by skeleton. Cracks/bonds/dislocations... are periodically closed by cyclic forcing. Little mobilization of damage features.

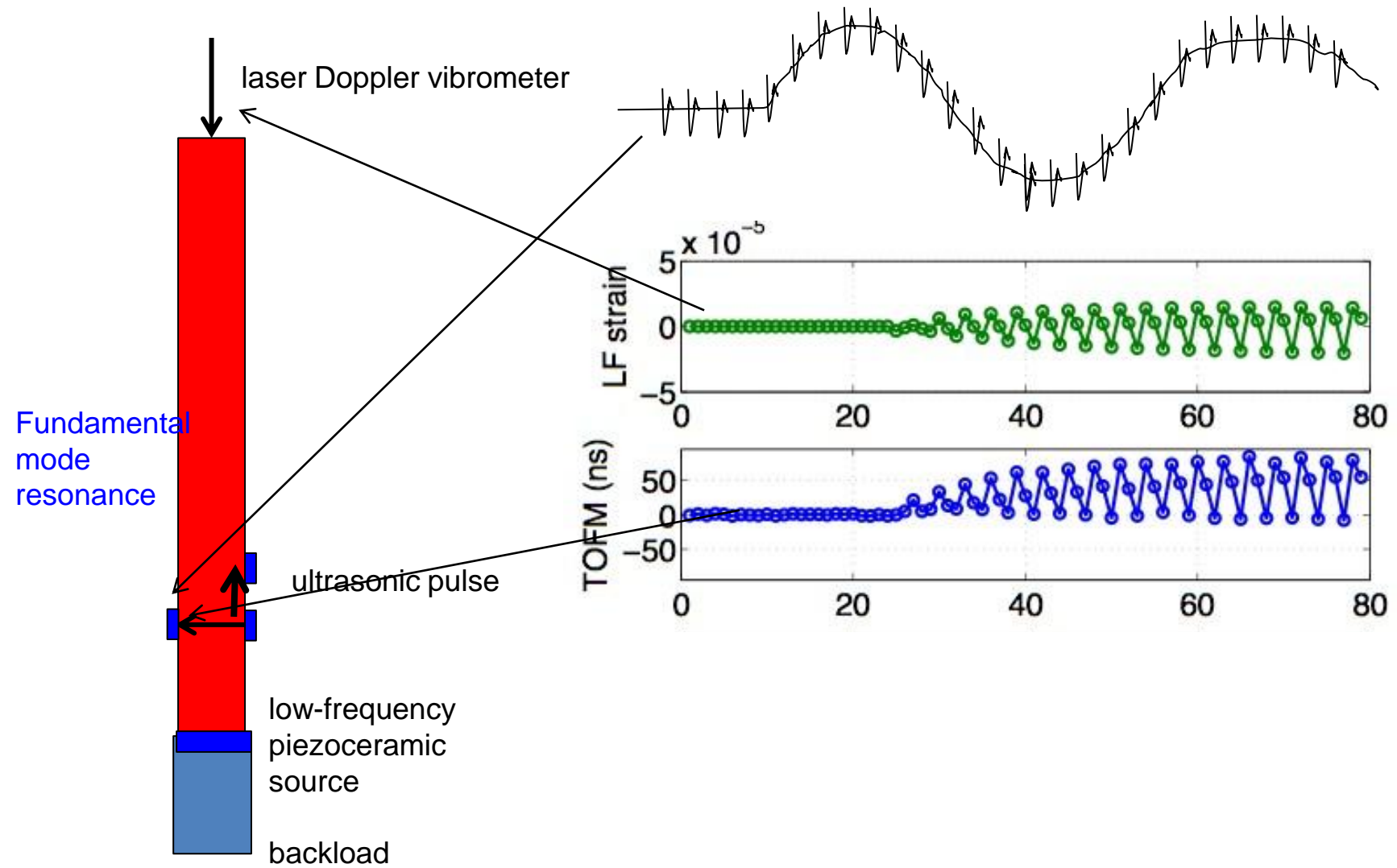
$$F \cos(\omega t - kx), \omega \gg 2\pi$$



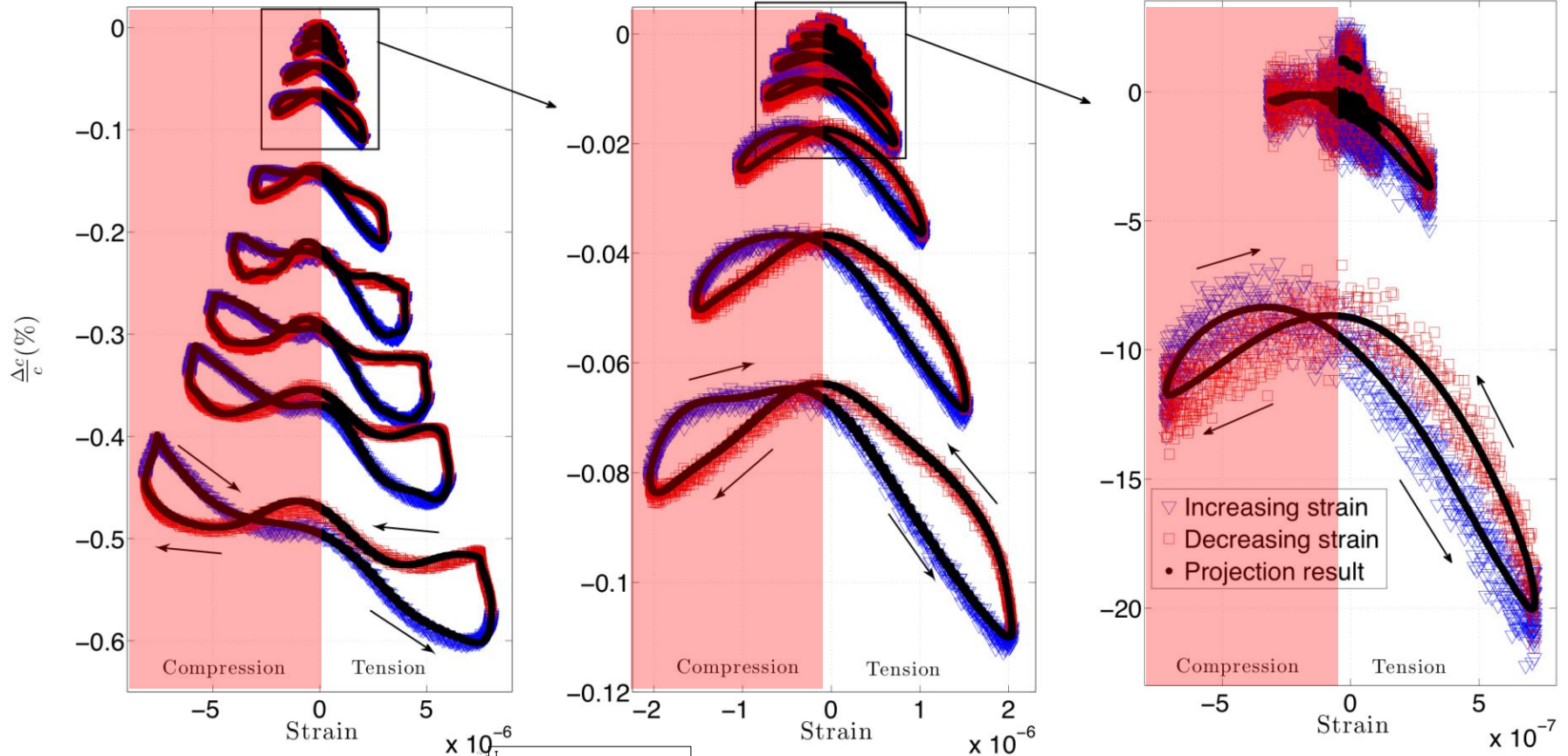
Progressively stronger stress mobilization of damage features, primarily by shearing (?)



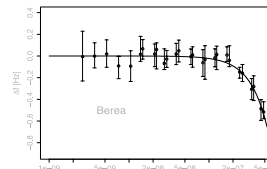
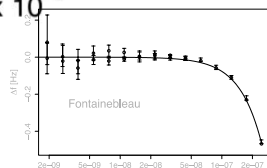
# Probing elastic nonlinearity with dynamic acousto-elasticity



# Dynamic acousto-elasticity reveals exotic behavior

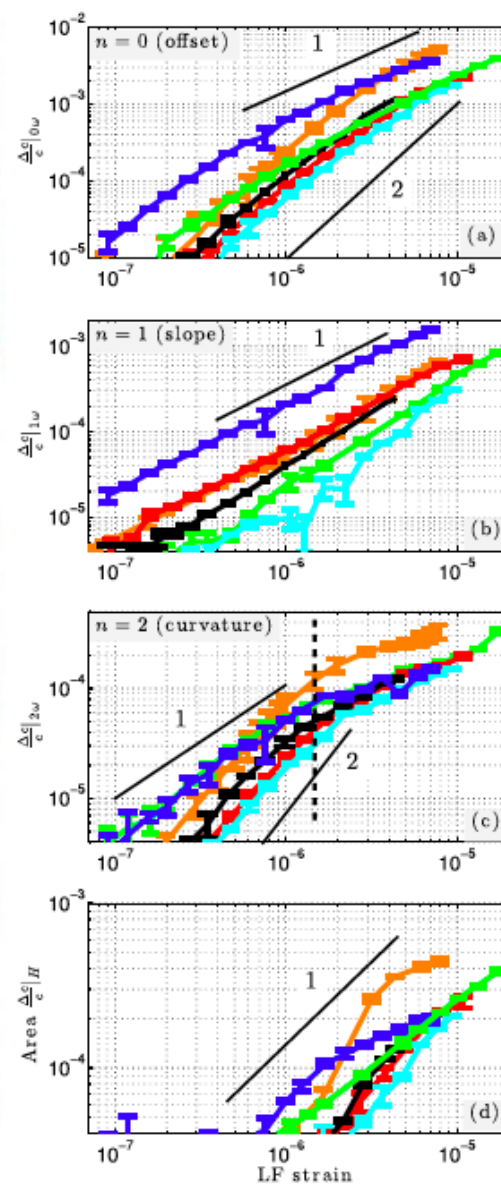
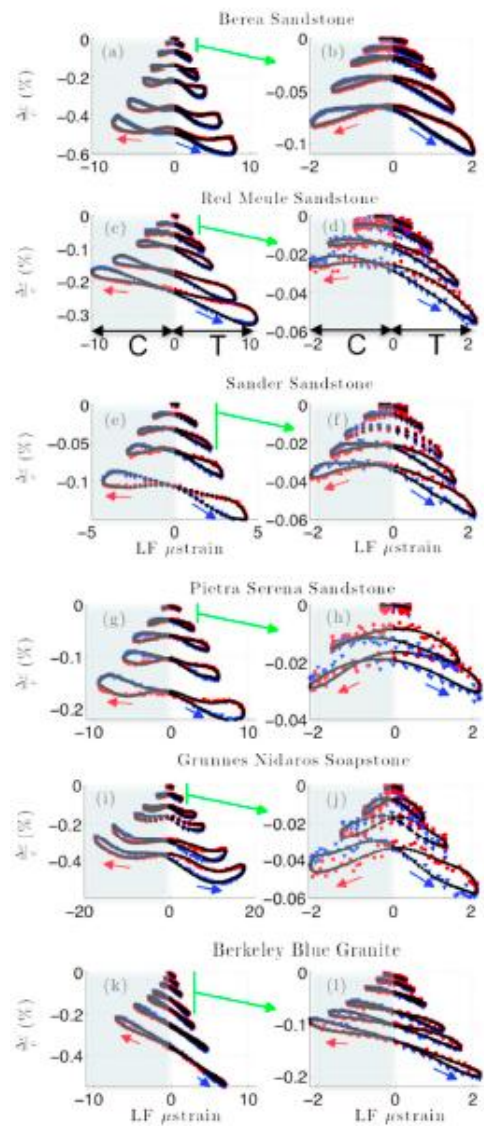


Note progressive offset with amplitude—'conditioning'= resonance peak shift under resonance conditions



Highly complex behavior that can be described with additional harmonics

# Nonlinear parameters from DAE



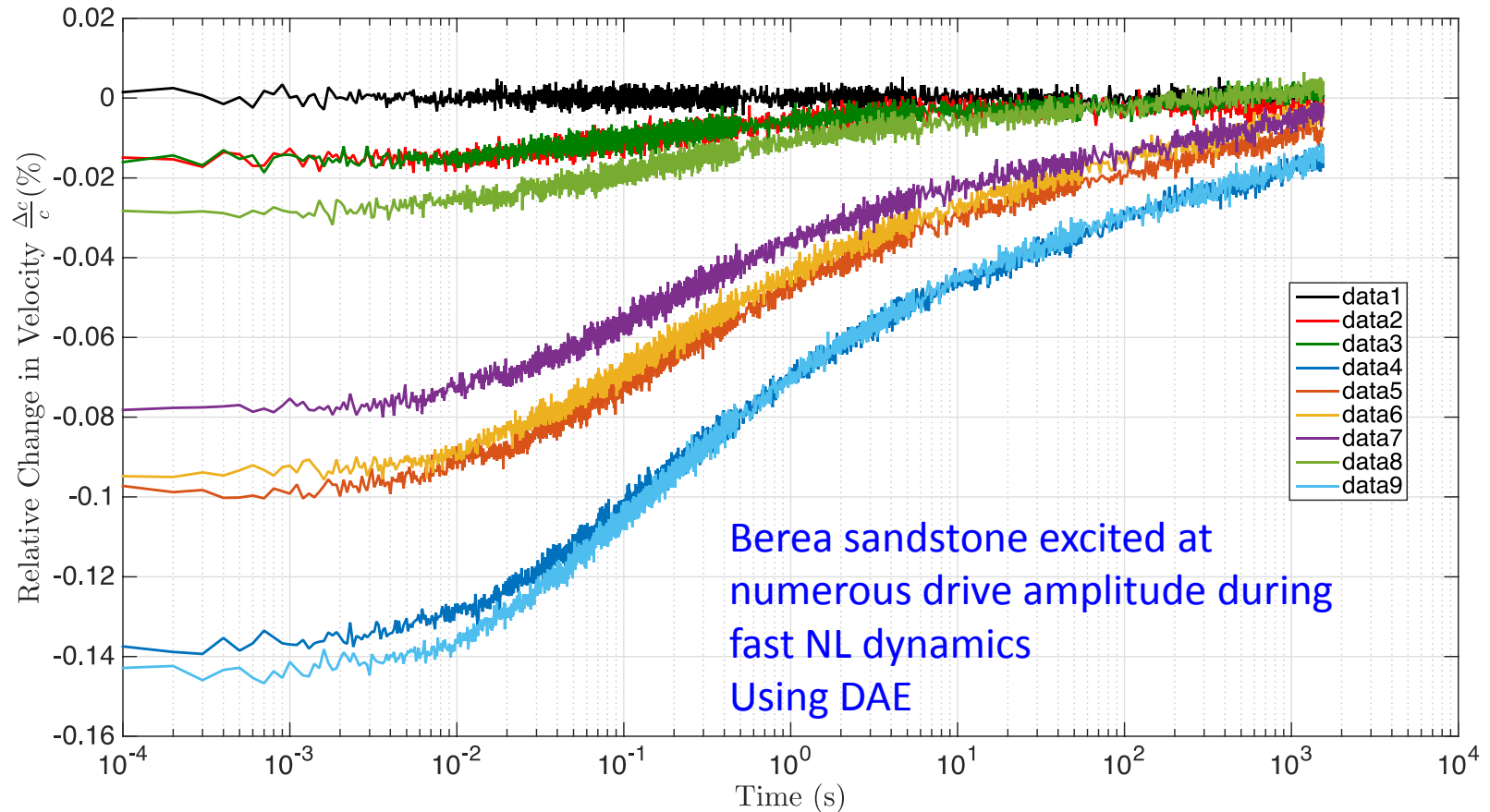
offset

slope

curvature

hysteresis

# Recovery– the slow dynamics: early time



Spectrum of barriers.....

It doesn't matter how beautiful your theory is, it doesn't matter how smart you are. If it doesn't agree with experiment, it is wrong.

*Richard P. Feynman*

# Some Models of non-classical nonlinearity

Applicability

## P-M, Ahrennius model

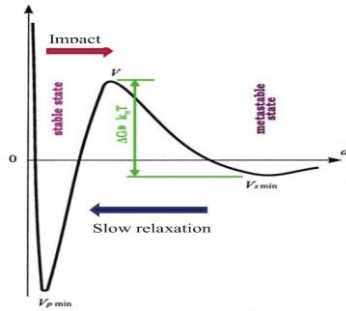
Phenomenologic, basic hysteric units:

$$s = K_0 \left( e + be^2 - a \left[ (De)e + \text{sign}(\dot{e}) \frac{e^2 - (De)^2}{2} \right] \right) + \exp(1/KT)$$

Bentahar et al.,  
Phys Rev. B, 2016

All systems

## Contact physics model



Lebedev and Ostrovsky  
Acoust. Phys.  
2014.

Relaxation of metastable contacts:

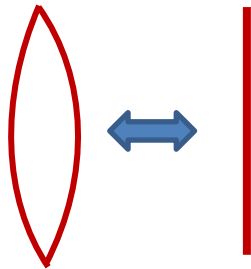
$$\frac{dE_s}{dt} = AE_0 \exp\left(-\frac{E_s}{\Lambda}\right) \quad (\Lambda = kT/V)$$

Sound speed relaxation (analytical result)

$$\frac{\Delta c}{c} = \frac{E_s}{2E_0} = \frac{\Lambda}{2E_0} \ln \left[ \frac{t}{\tau} + e^{-|E_{s,0}/\Lambda|} \left( 1 - \frac{t}{\tau} \right) \right]$$

Consolidated and pre-compressed unconsolidated materials.  
Models hysteresis and slow time relaxation  
Provides a contact scale

## Crack healing model



Snieder et al,  
Geophys. J. Intern.  
Adv. Access, 2016

Superposition of relaxations:

$$R(t) = \int_{\tau_{min}}^{\tau_{max}} \frac{1}{\tau} e^{-t/\tau} d\tau$$

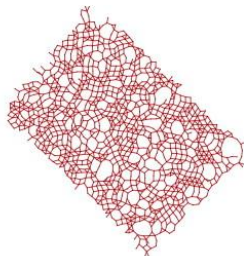
Snieder, next

Strain relaxation:

$$\dot{\epsilon} = -\frac{1}{n\eta} \left( N\sigma - \sum_{i, \text{closed}} \kappa \epsilon_i \right)$$

Consolidated materials with microcracks.  
Models hysteresis and slow time relaxation

## STZ model



Lieu et al,  
JGR Solid Earth  
2017 (submitted)

Effective temperature

$$\Lambda \equiv \frac{N_+ + N_-}{N} = 2e^{-1/\chi}$$

[Daub last Monday]

Linear resonance:

$$\omega_{res}^2 = \left( \frac{\pi}{H} \right)^2 \frac{M}{\rho G}$$

Non-consolidated materials with grain restructuring.  
Models fast time processes  
Including nonlinear resonance and losses



“the story is true. I may have  
skipped over some of the details  
along the way.”

*The Ben Zion, Cargese, 2017*

# Take away: nonlinear fast and slow dynamics

