

Scattering

(Michel Campillo)

Seismograms and regime of propagation/scaling

From field to intensity and energy

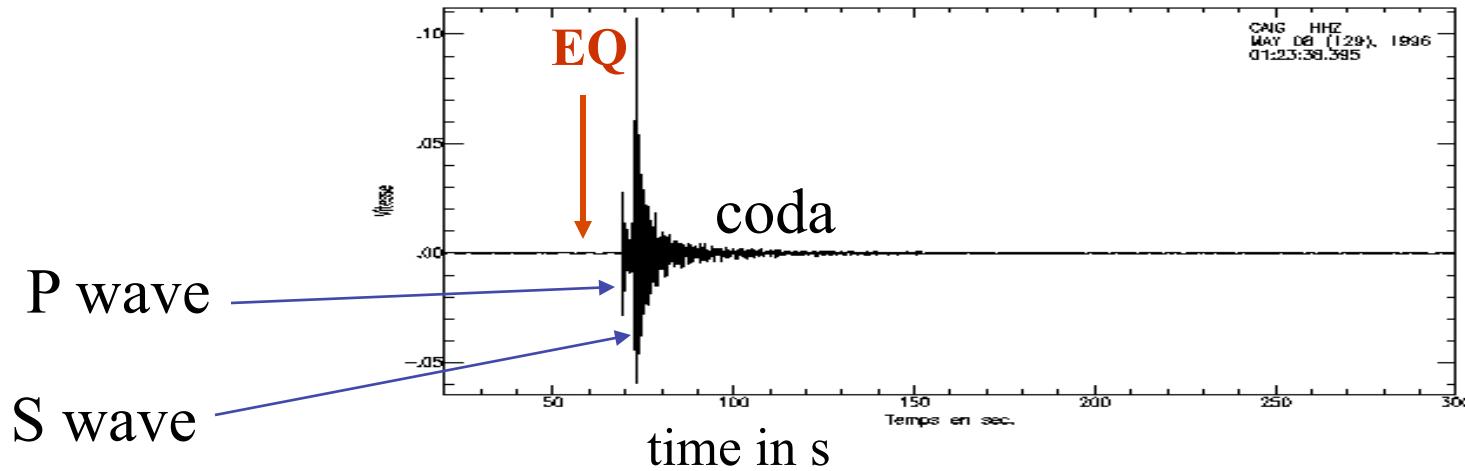
RTE

Diffusion

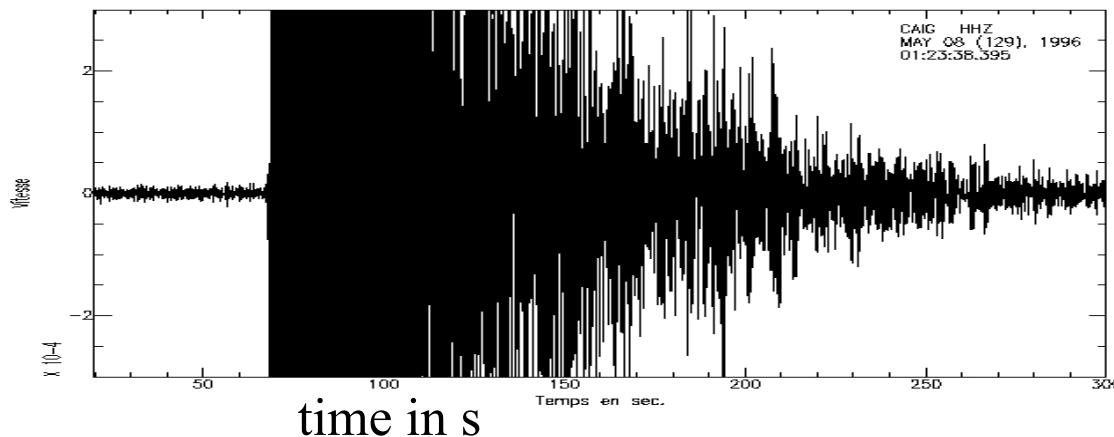
Consequences of scattering: equipartition and isotropization

Scattering in the mantle

Example of a record of a local earthquake in the band .5-20Hz



scale x 1000:

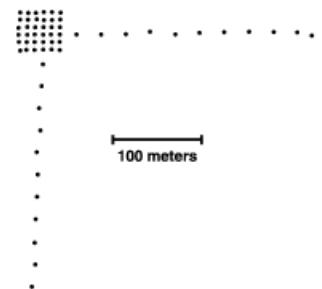
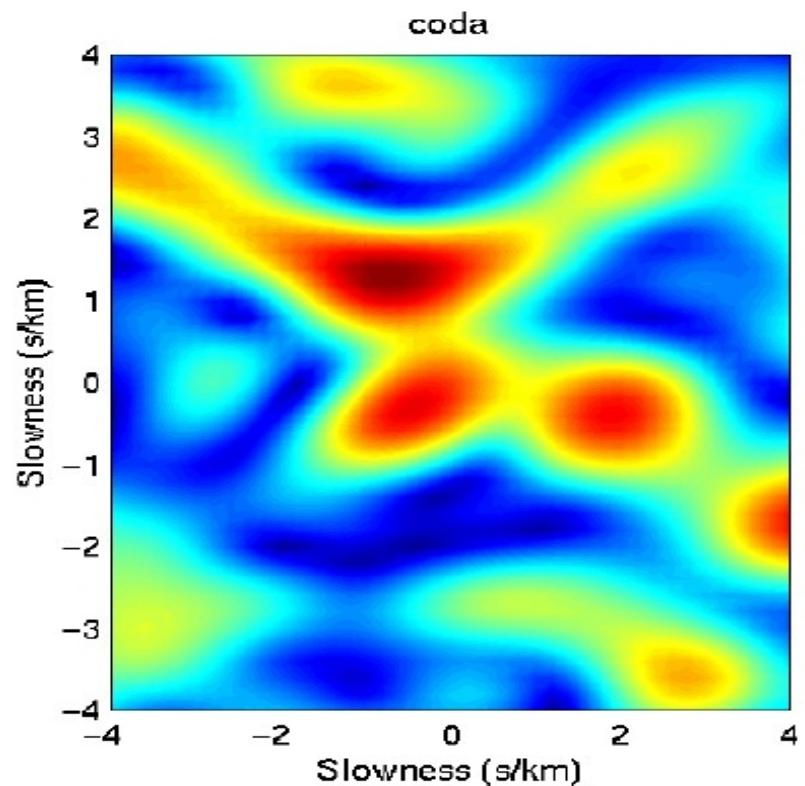
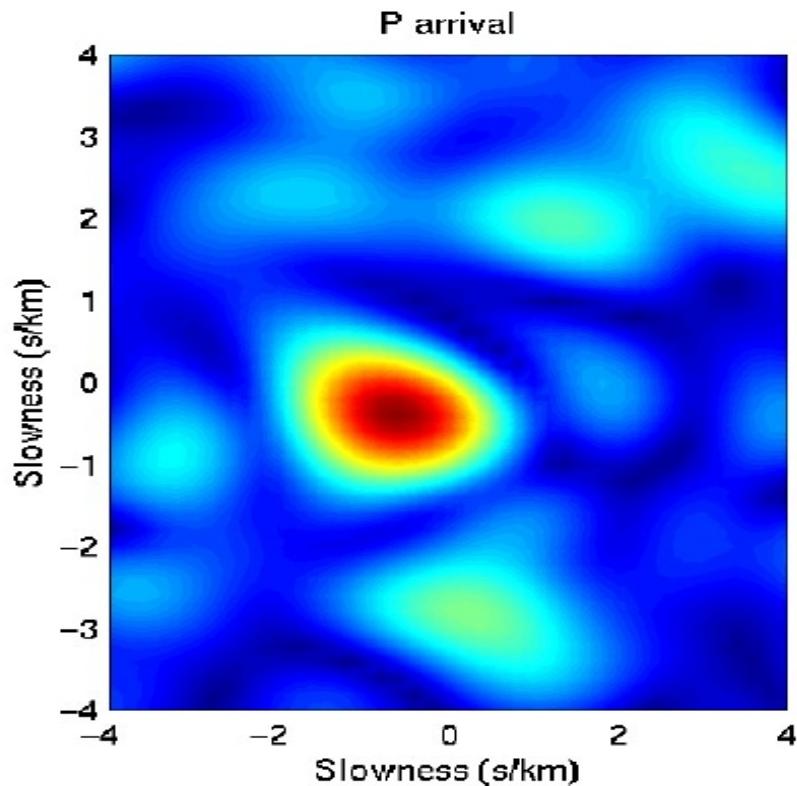


Coda: tail, end of a piece of music....

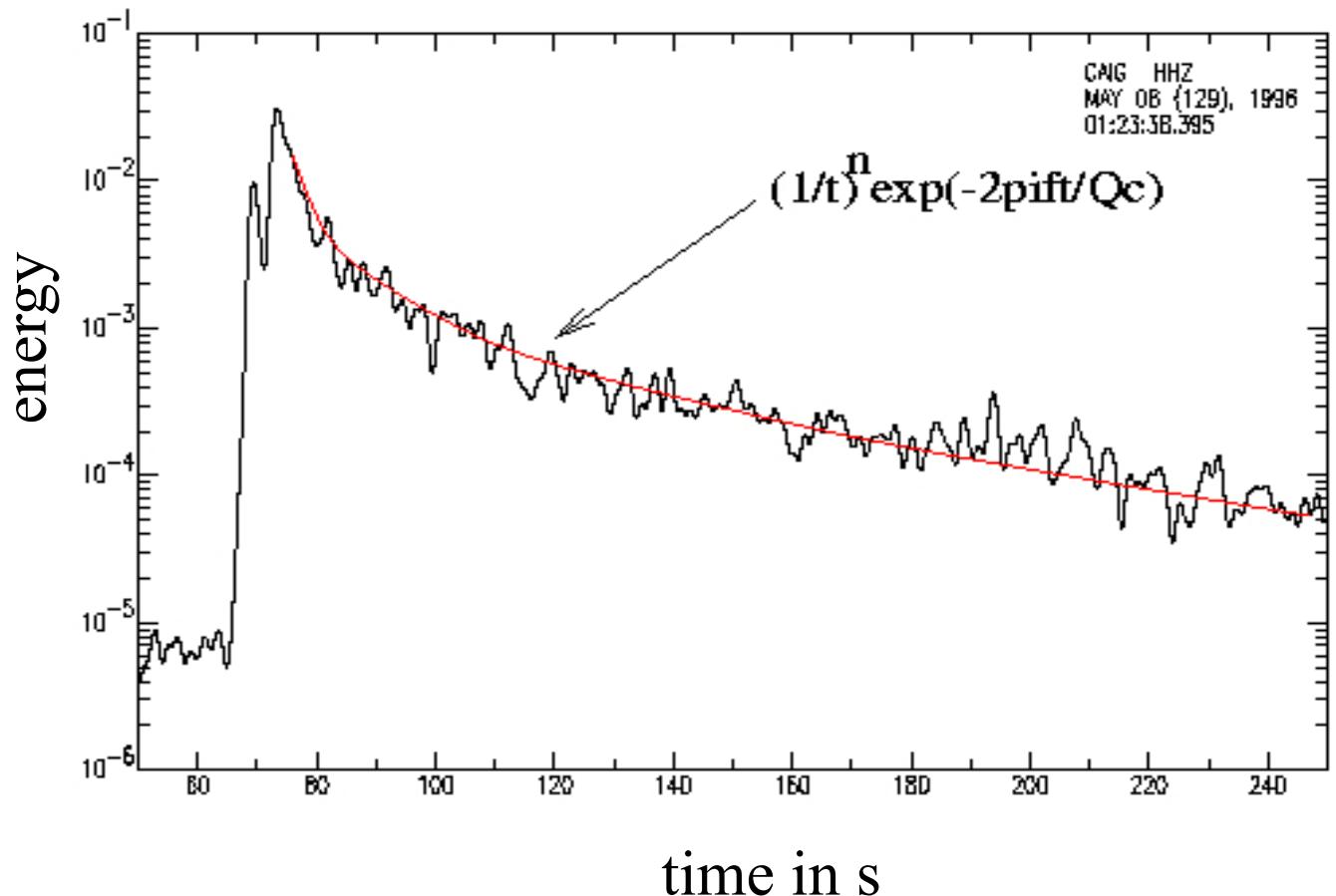
Frequency-wavenumber analysis

(Pinon Flat Seismometer Array)

$$u(x,y) \rightarrow u(k_x, k_y)$$

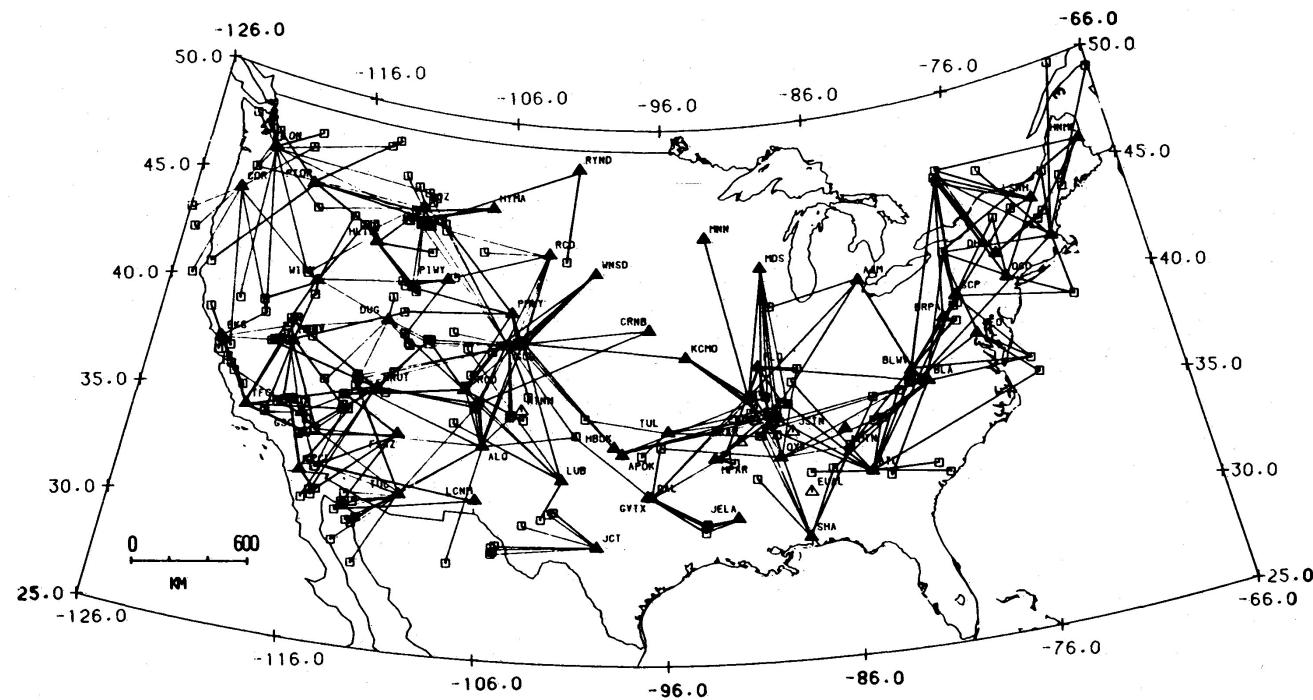


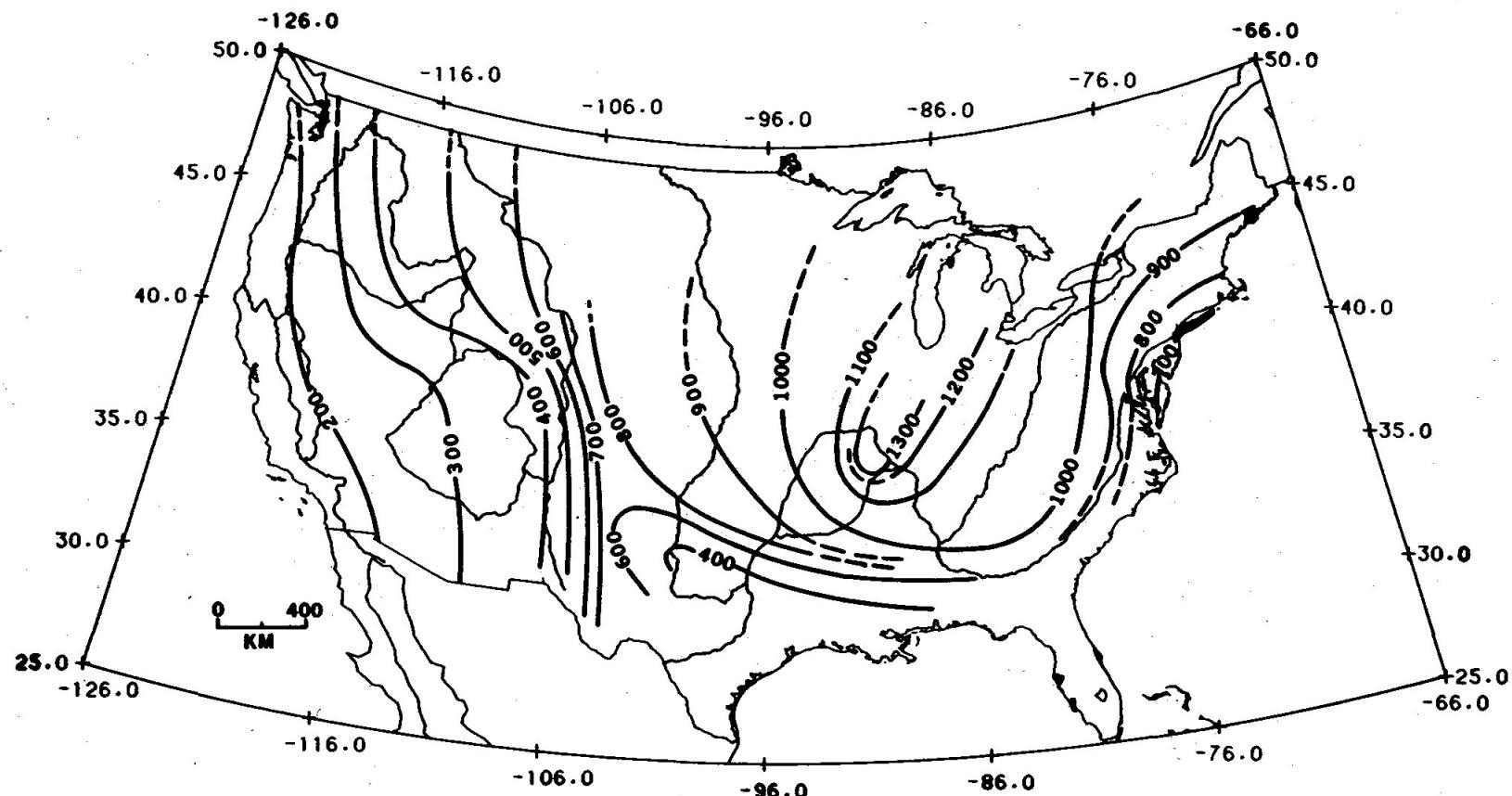
Energy decay in the coda (Aki and Chouet, 1975)



The decay is constant in a region, independently of source and receiver: Q_{coda}

Coda Q in US (Singh and Herrmann, 1983)





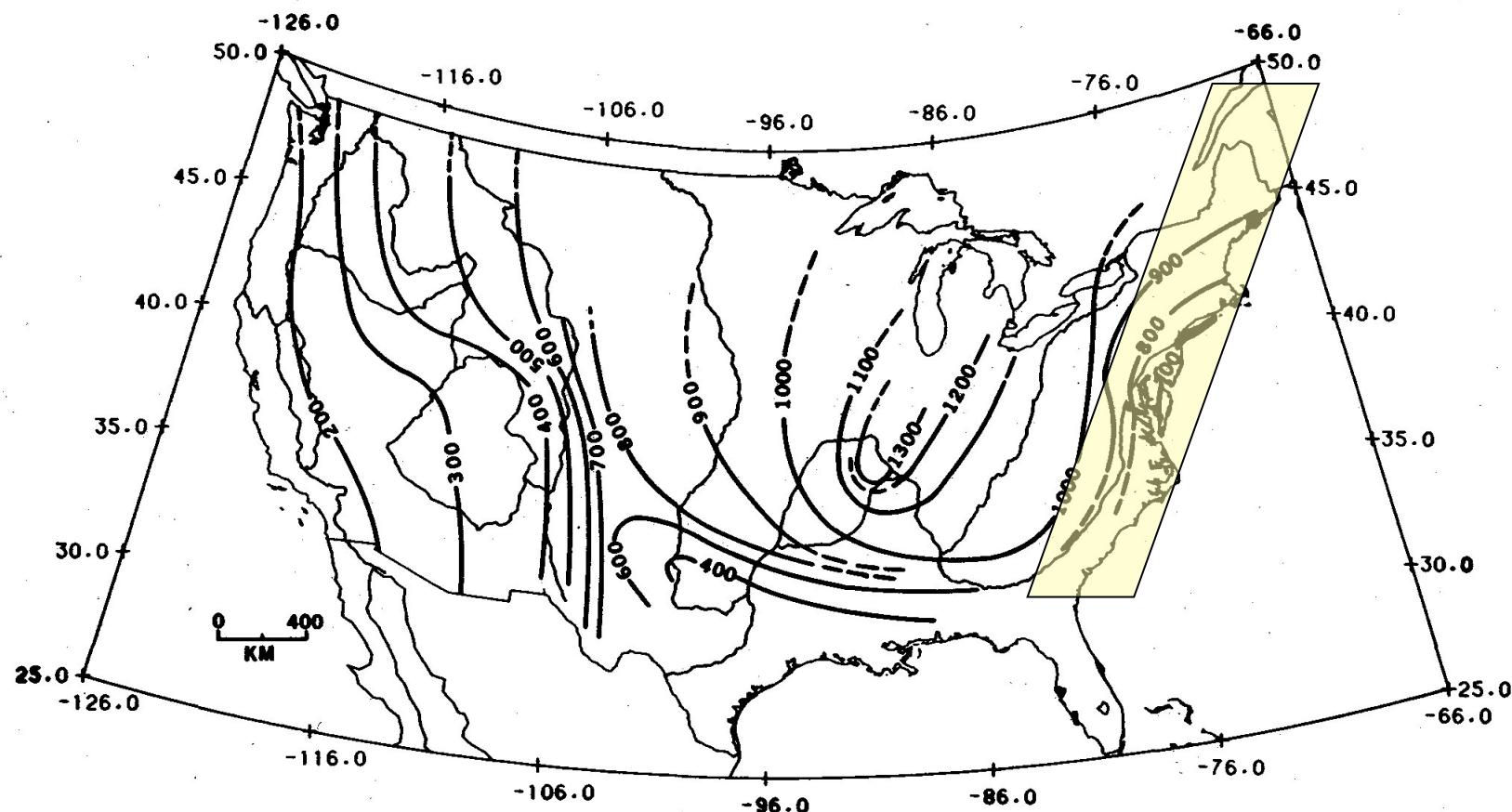


Fig. 15. Contour map of coda Q_0 for the entire continental United States.

Appalachian (Hercynian) belt : $Q_c \sim 600$

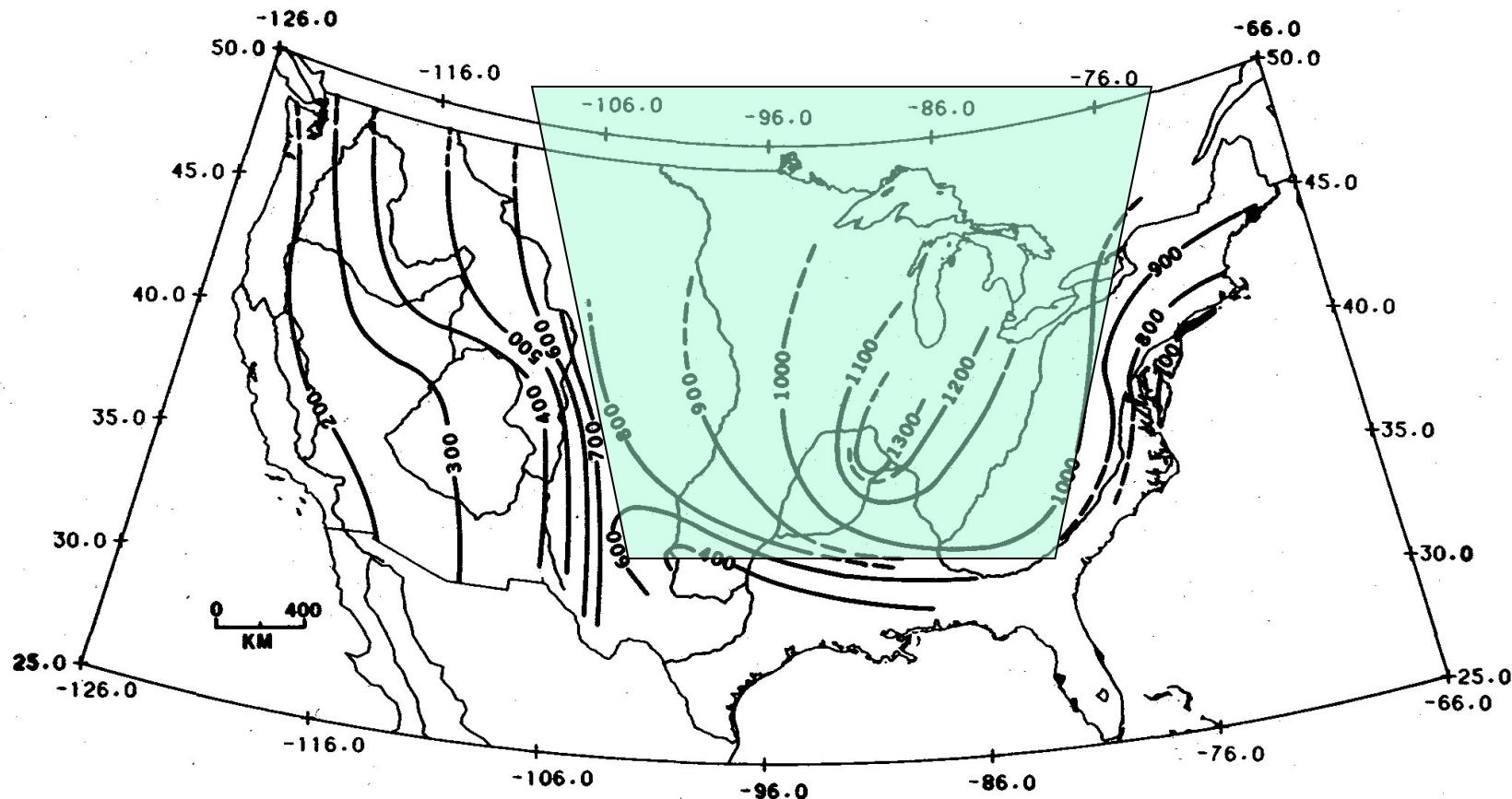


Fig. 15. Contour map of coda Q_0 for the entire continental United States.

Central shield : $Q_c \sim 1000$

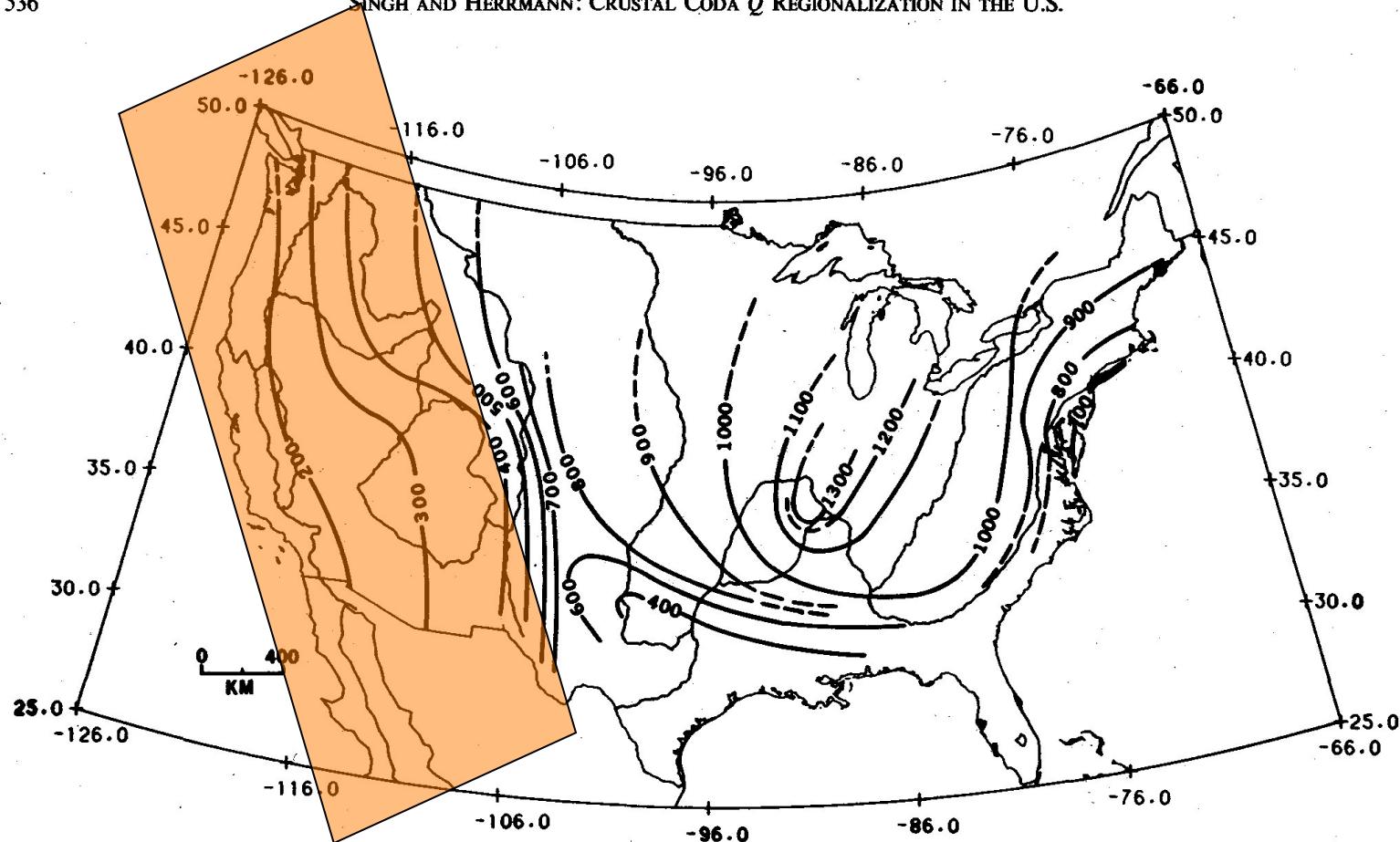
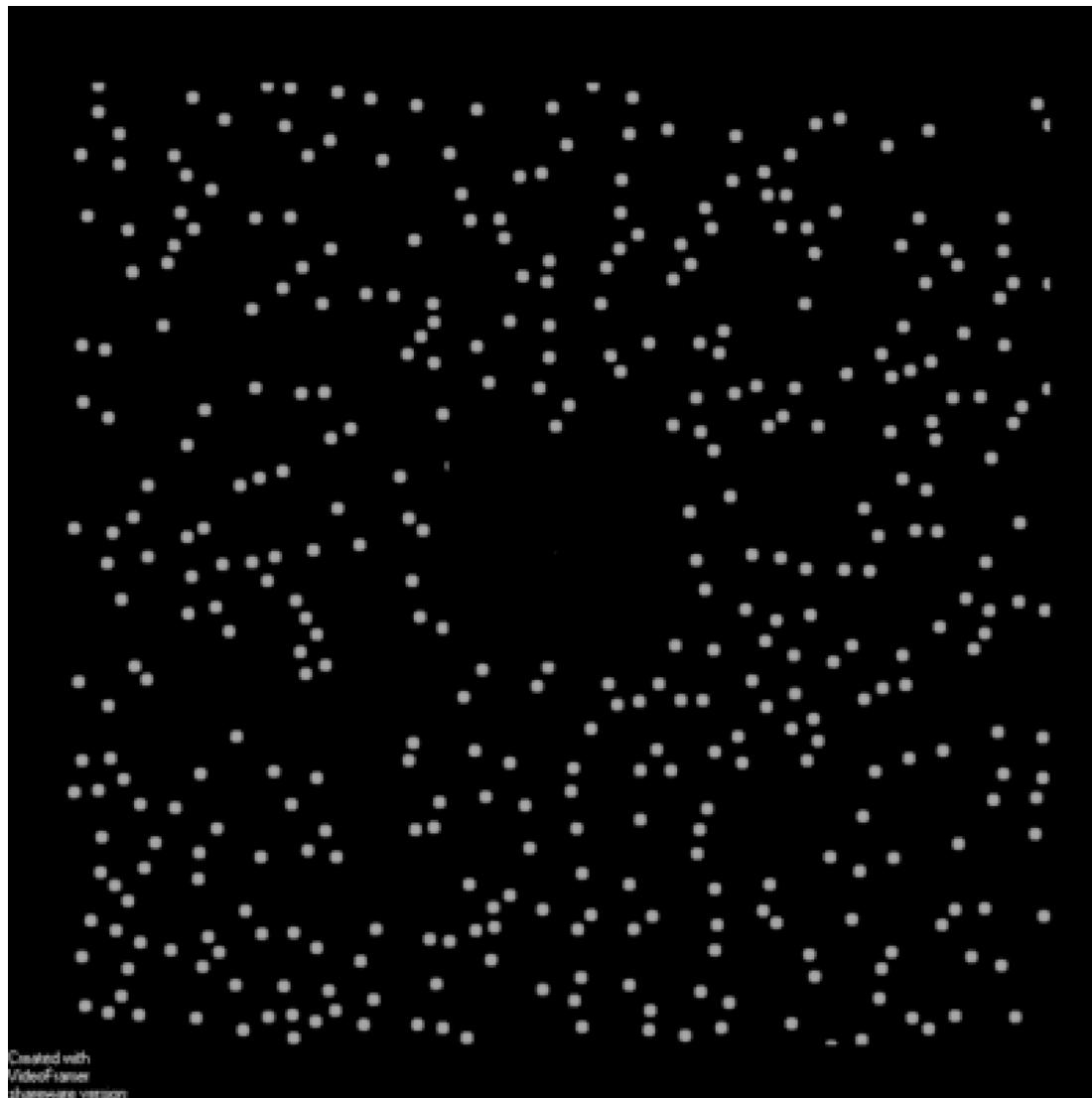


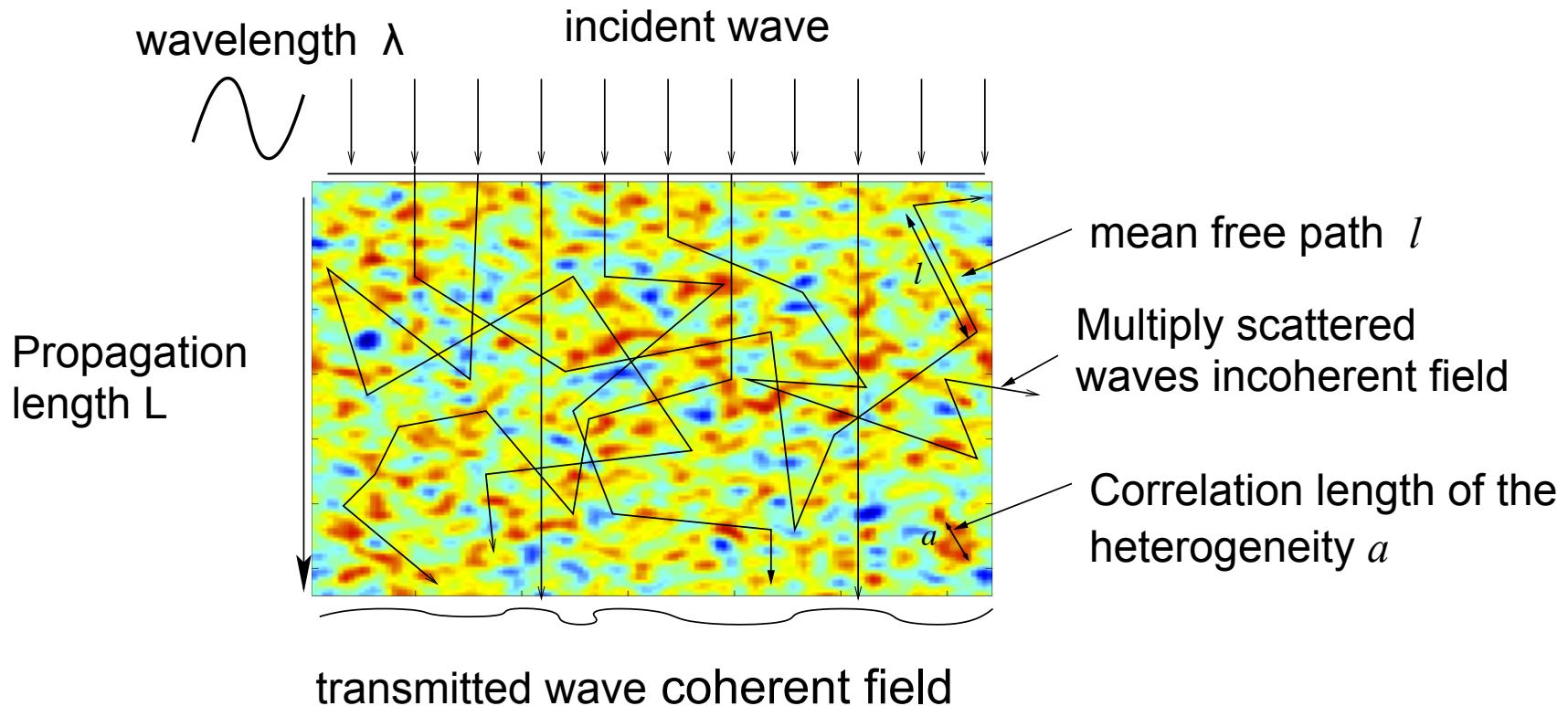
Fig. 15. Contour map of coda Q_0 for the entire continental United States.

Tectonically active western US: $Q_c=100-300$

Transient signals in a complex medium.....



Propagation regimes



$$L < l, l_a$$

$$l_a > L \gg l$$

coherent wave

single to multiple scattering

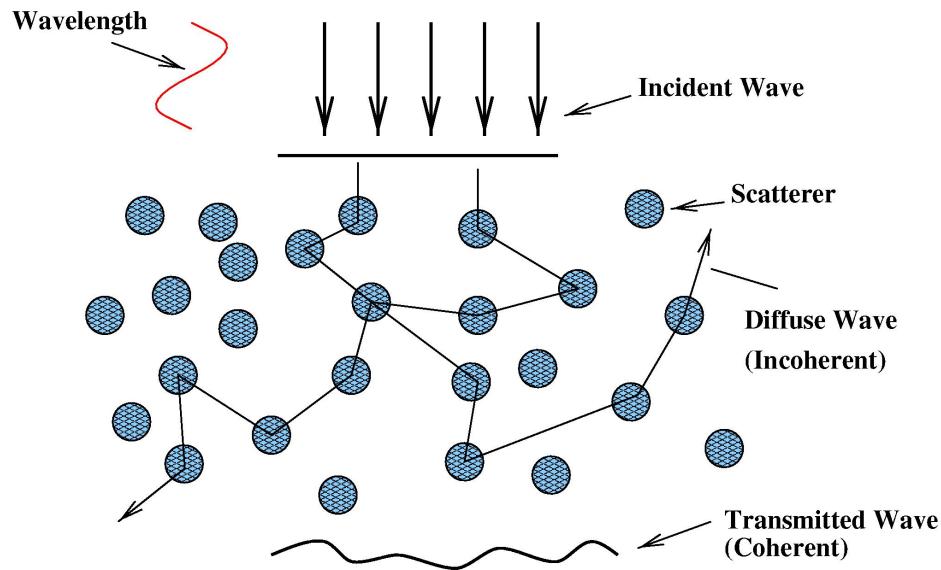
equipartition

mean field

radiative transfer

Diffusion

Wave Propagation through Random Media



Length Scales: λ , Correlation Length, Propagation Distance

Question: Ensemble Average Response?

Precise Definition of Coherent and Incoherent Waves

The Concept of Mean Free Path

First Moment of the Green Function:

- Dyson Equation

Ensemble average $\langle G \rangle$: Coherent Field

$$\langle G \rangle = G_0 + G_0 M \langle G \rangle$$

M , Mass Operator Describes all Possible Scattering Situations

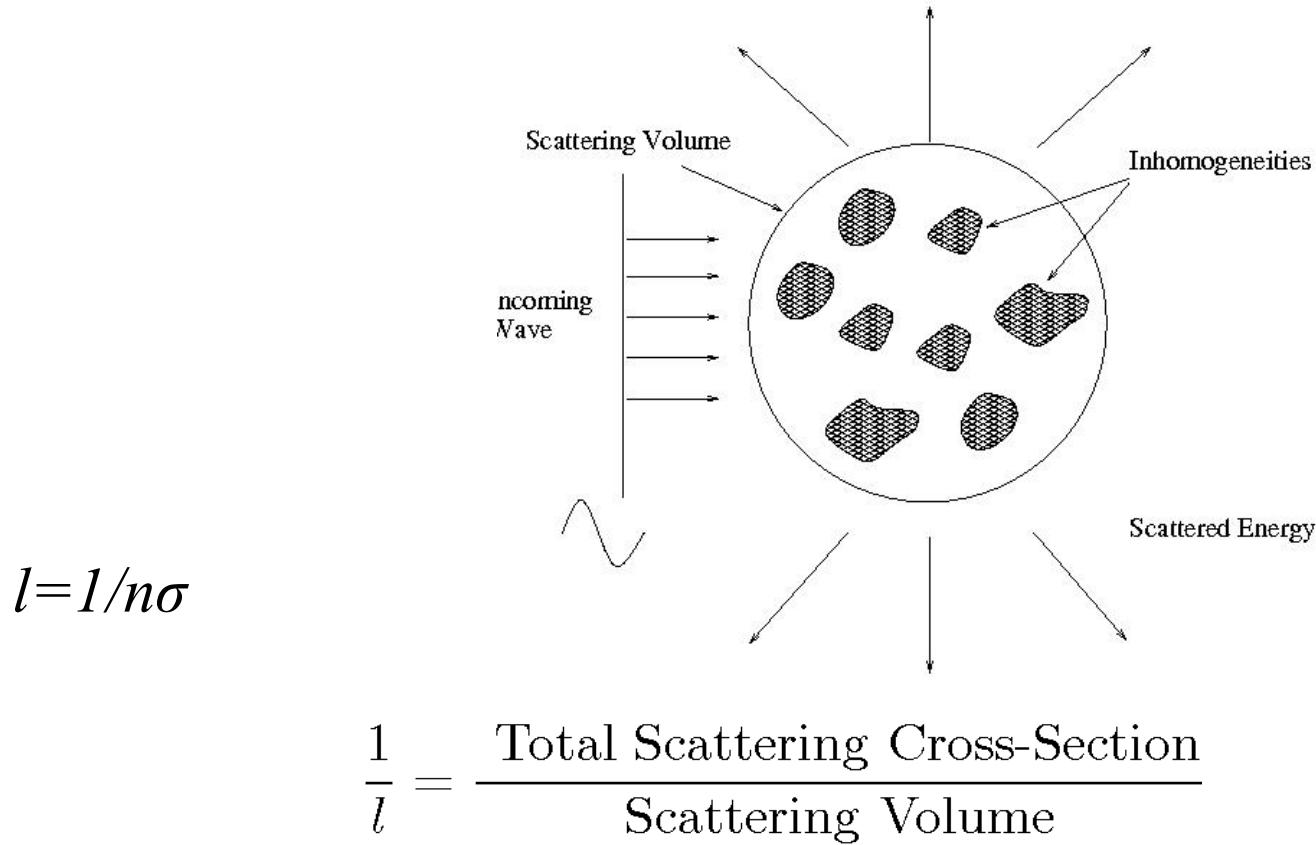
- Approximate Solution

$$\langle G(\vec{r}; \vec{r}_0) \rangle = -\frac{1}{4\pi|\vec{r} - \vec{r}_0|} e^{ik|\vec{r} - \vec{r}_0|}$$

$$k = k_0 + \frac{i}{2l}$$

New Length Scale: Mean Free Path of Waves $l = f(\epsilon, a, \lambda)$

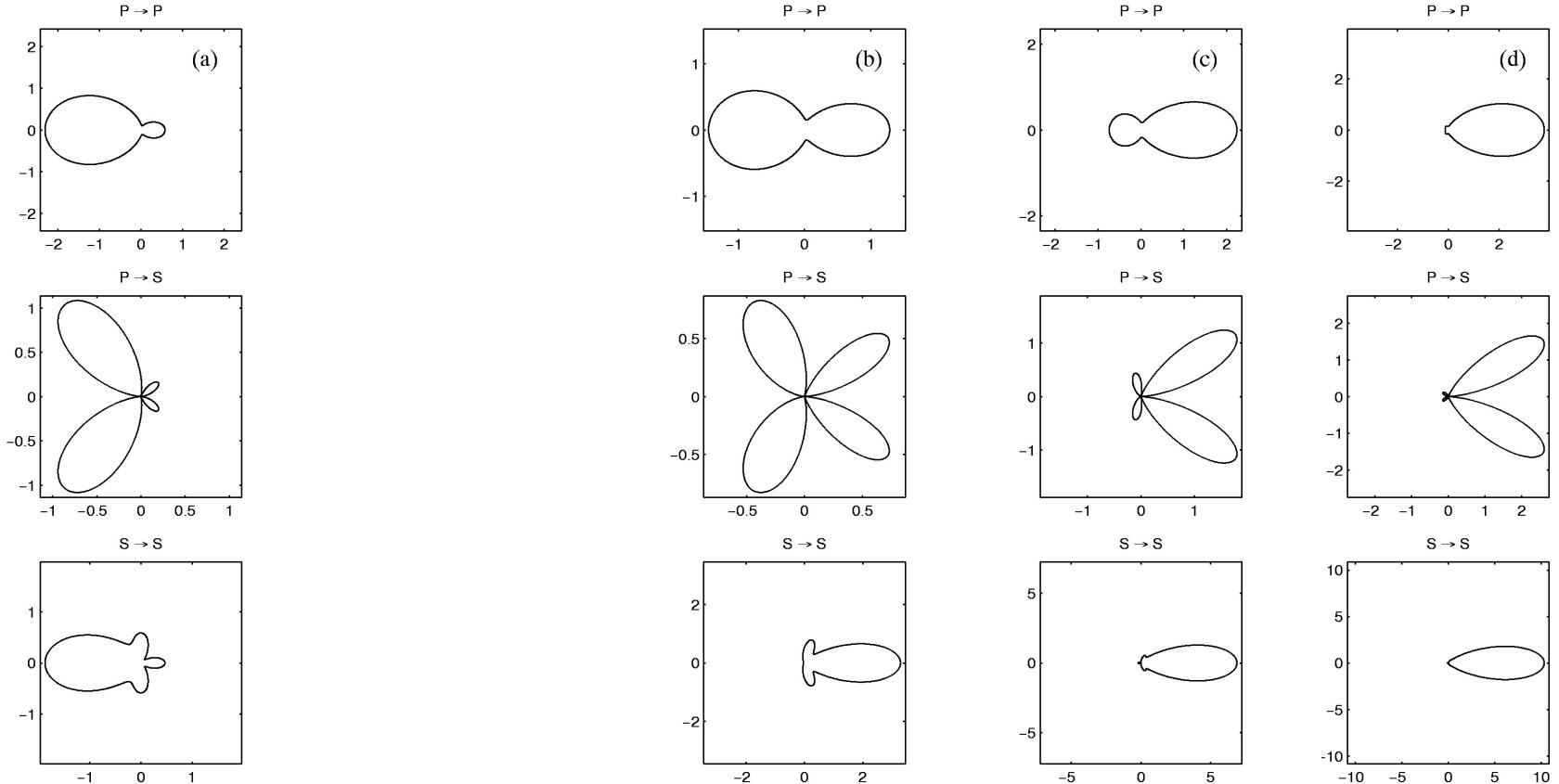
Exponential form of decay: $\exp(-r/2l)$ to be compared with attenuation



Frequency Dependence:

- Rayleigh ($ka \ll 1$): $l \sim \omega^{-4}$
- High-Frequency ($ka > 1$): $l \sim \omega^{-2}$

Differential cross sections of scattering and conversion for a sphere of radius a



$ka \rightarrow 0$: Rayleigh approximation

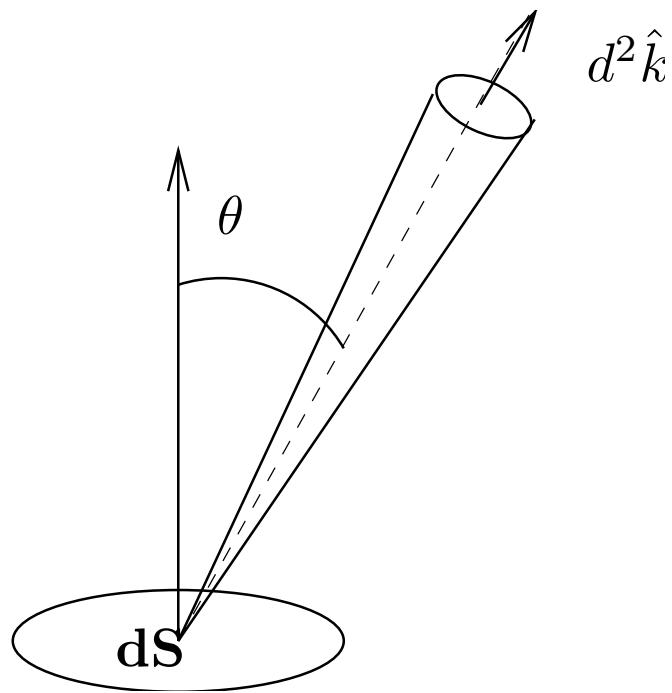
$ka = 1.2$

(averaged in φ for S polarisation)

$ka = 1.6$

$ka = 2.0$

The specific intensity



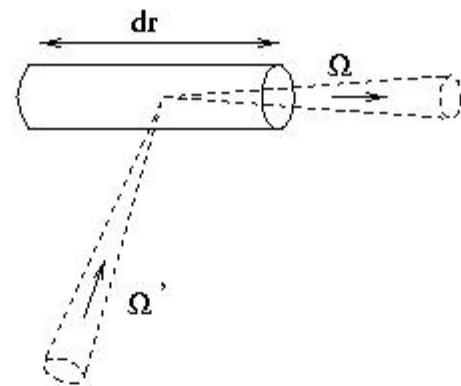
Definition

$\mathcal{I}(\omega, t, \mathbf{r}, \hat{\mathbf{k}}) \times dS \times \cos(\theta) \times dt \times d\omega =$ Amount of energy within
the frequency band $[\omega, \omega + d\omega]$
flowing through dS around
direction $\hat{\mathbf{k}}$ during time dt

Angularly-resolved energy flux
through a surface

Radiative transfer equation

Energy balance of a beam of energy propagating a distance dr in the scattering medium



Variation of Intensity

=

Loss due to scattering into all space directions

+

Gain due to scattering from direction $\vec{\Omega}'$ to direction $\vec{\Omega}$

The Equation of Radiative Transfer

Scalar Case

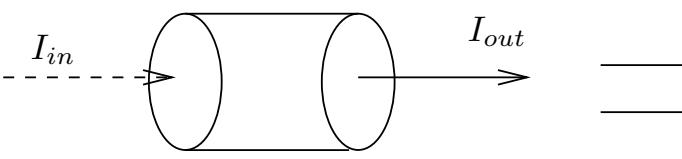
$$\left(\frac{1}{c} \frac{\partial}{\partial t} + \hat{\mathbf{k}} \cdot \nabla_{\mathbf{x}} \right) \mathcal{I}(t, \mathbf{x}, \hat{\mathbf{k}}) = - \left(\frac{1}{l^s} + \frac{1}{l^a} \right) \mathcal{I}(t, \mathbf{x}, \hat{\mathbf{k}}) + \frac{1}{l^s} \oint p(\hat{\mathbf{k}}, \hat{\mathbf{k}'}) \mathcal{I}(t, \mathbf{x}, \hat{\mathbf{k}'}) d^2 \hat{\mathbf{k}'}$$

l^s : scattering mean free path

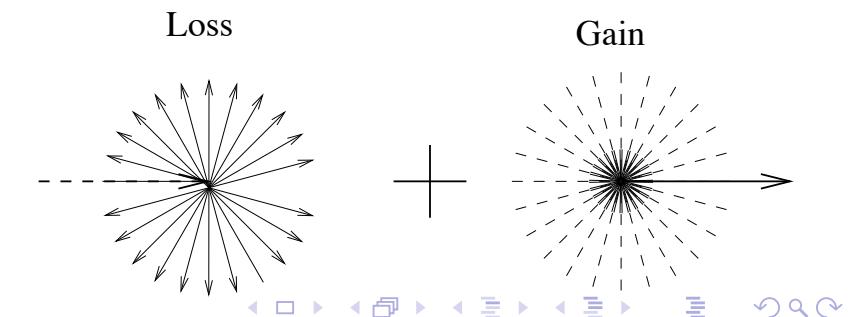
l^a : absorption length

c : wave speed

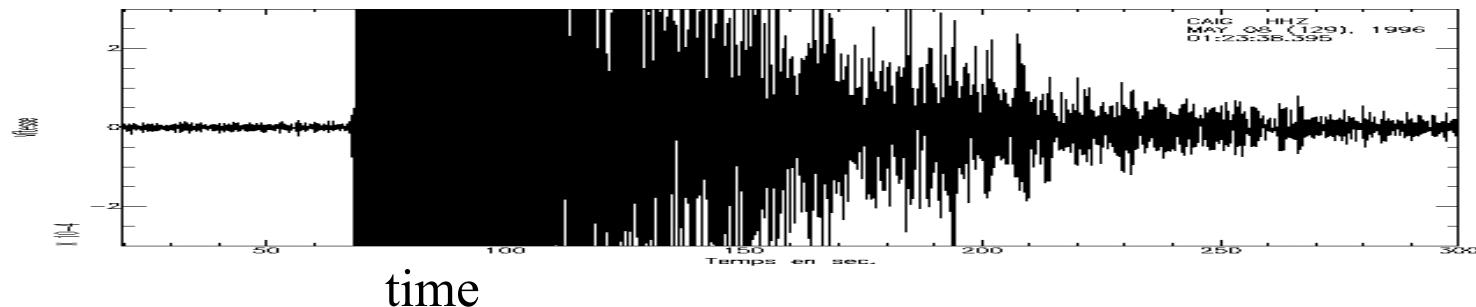
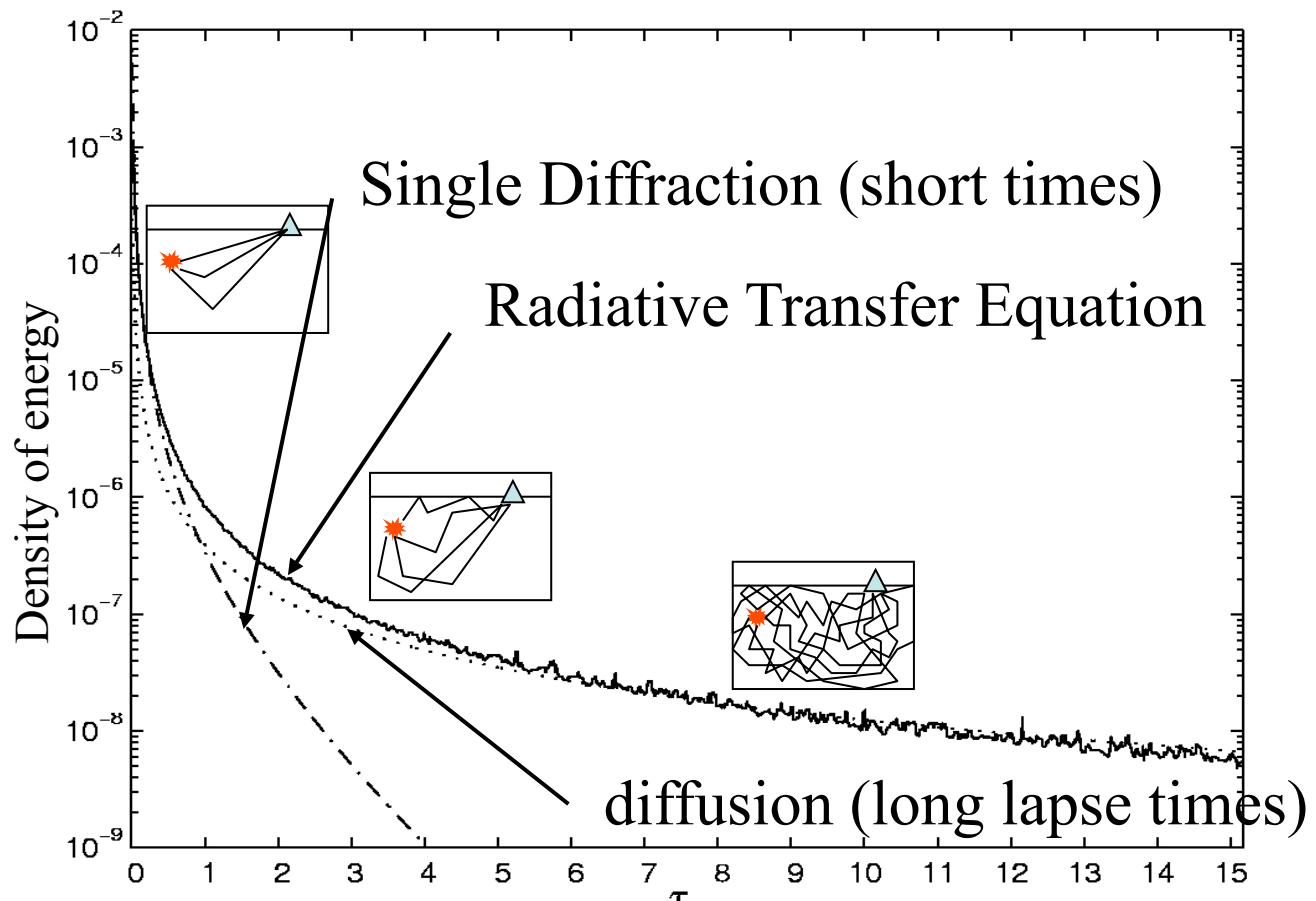
$p(\hat{\mathbf{k}}, \hat{\mathbf{k}'})$: scattering anisotropy

$$\delta I = I_{out} - I_{in}$$


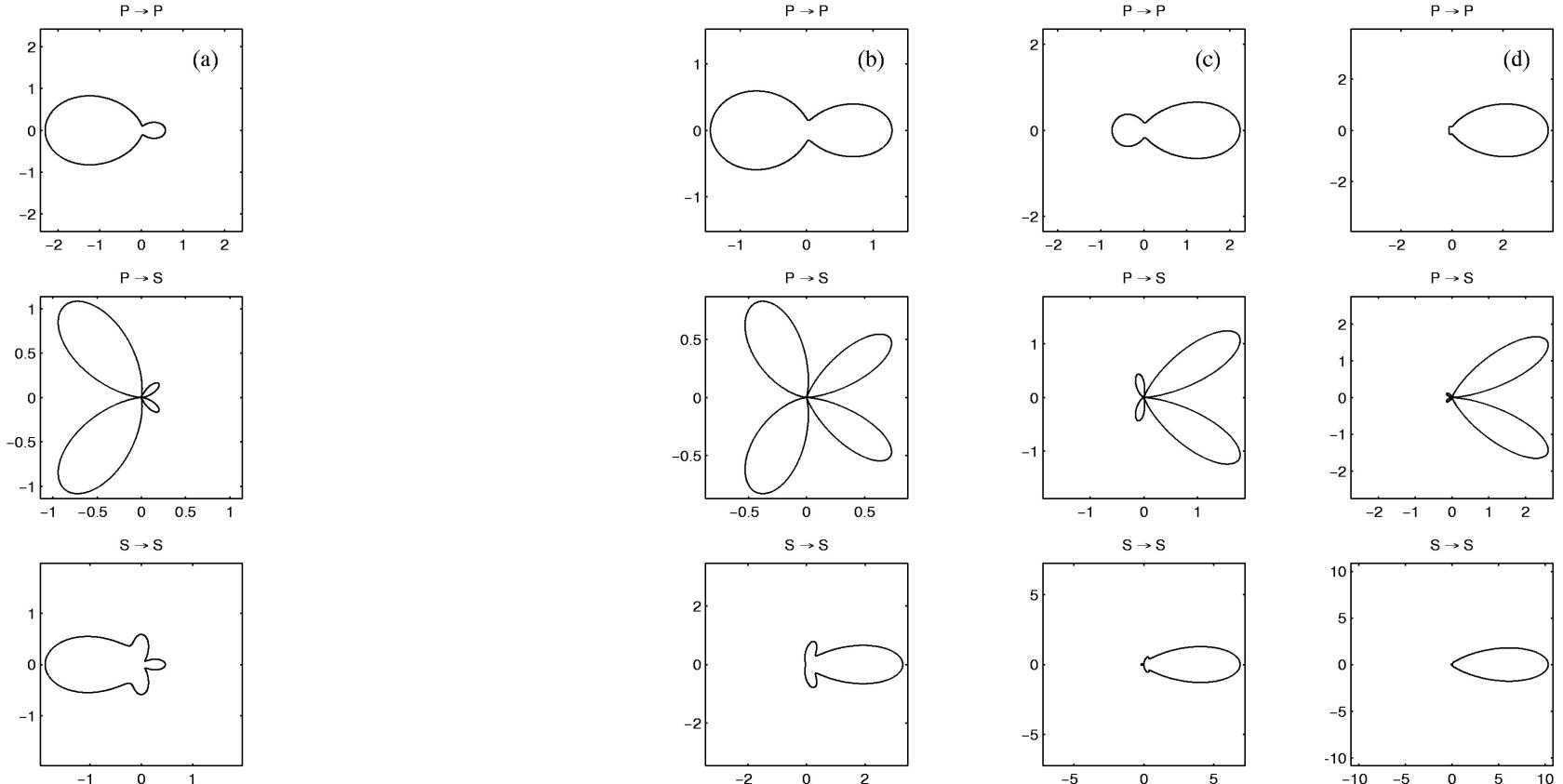
I_{in} I_{out}
 cdt



Propagation regimes and description of energy



Differential cross sections of scattering and conversion for a sphere of radius a



$ka \rightarrow 0$: Rayleigh
approximation

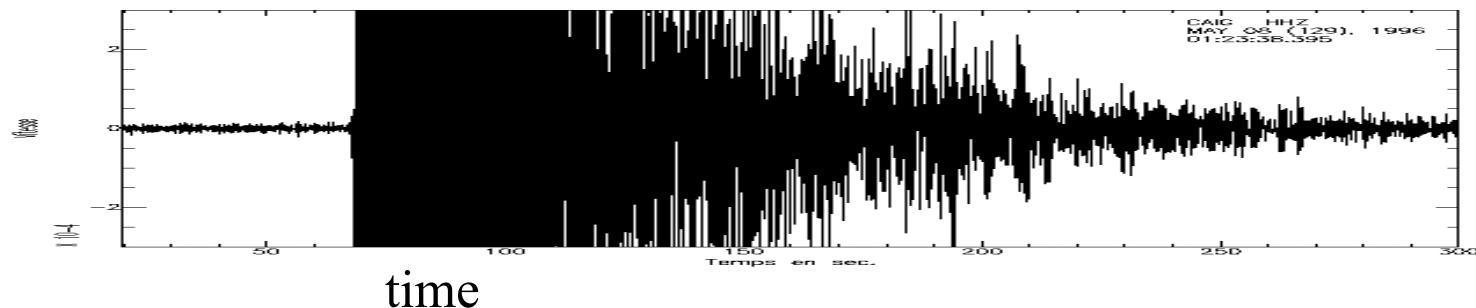
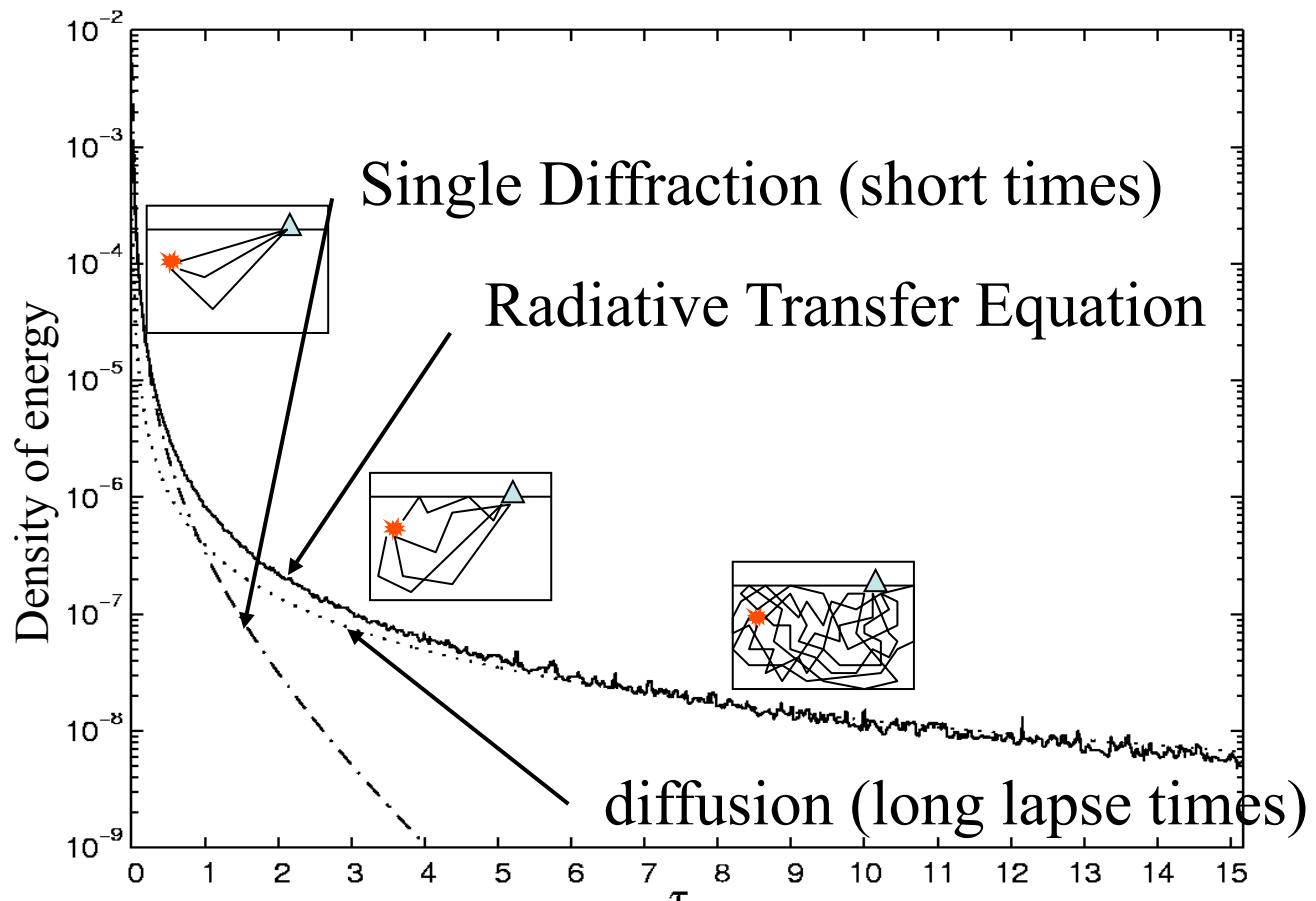
$ka = 1.2$

$ka = 1.6$

$ka = 2.0$

(averaged in φ for S polarisation)

Propagation regimes and description of energy



The Equation of Radiative Transfer

Second moment of the Green's function is governed by the Bethe-Salpeter equation:

$$\langle GG^* \rangle = \langle G \rangle \langle G^* \rangle + \langle G \rangle \langle G^* \rangle K \langle GG^* \rangle$$

K , Intensity Operator describes all scattering situations.

Neglecting recurrent scattering leads to:

$$\partial_t I(t, \vec{\Omega}, \vec{r}) + \vec{\Omega} \cdot \vec{\nabla}_r I(t, \vec{r}, \vec{\Omega}) = -\frac{1}{l} + \frac{1}{4\pi l} \int d\vec{\Omega}' I(t, \vec{\Omega}', \vec{r}) P(\vec{\Omega}, \vec{\Omega}')$$

Describes the transport of the incoherent part of the intensity.

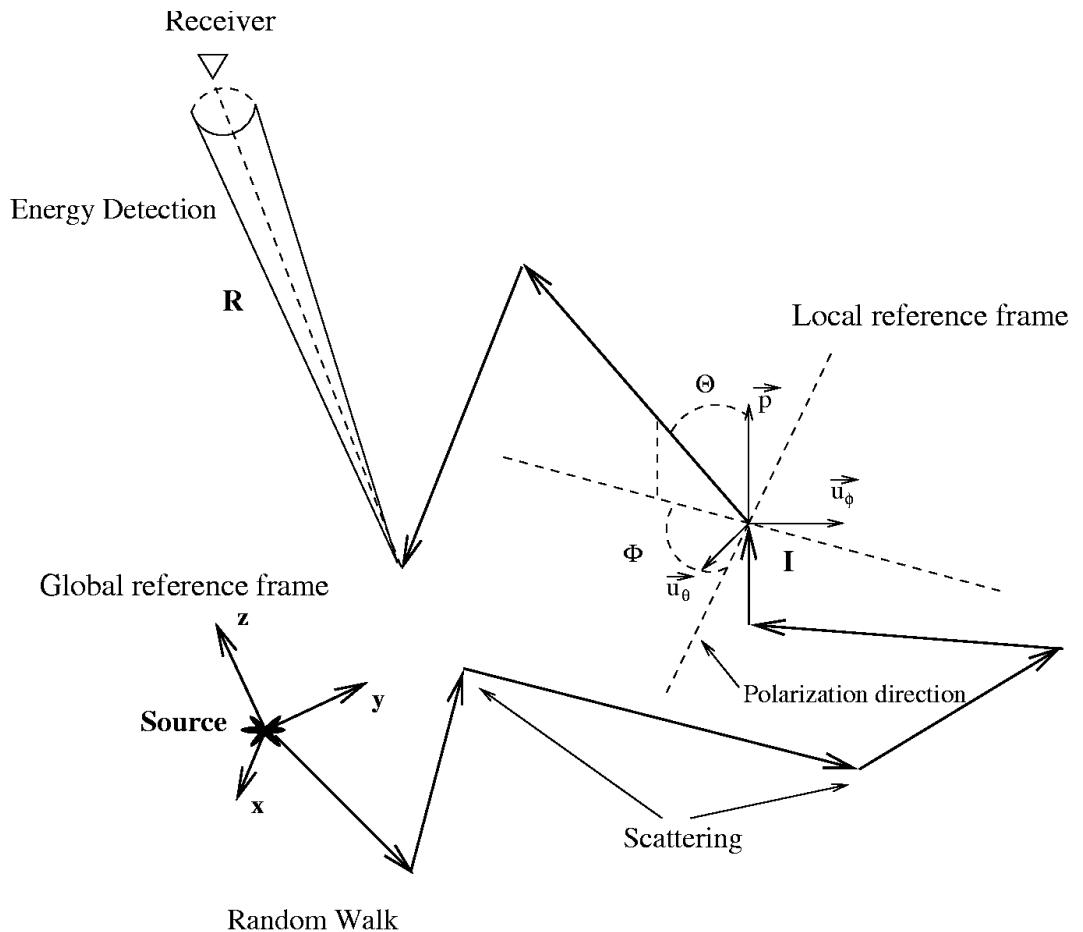
I , Specific Intensity function of space direction, time and position

$P(\vec{\Omega}, \vec{\Omega}')$, phase function (matrix) related to the power spectrum of the inhomogeneities

The radiative transfer equation

« particle analogy »

propagation under the
ray theory assumptions



parameters: $l_p, l_s \dots$, differential cross-sections

Single Scattering Approximation

The waves interact only once with the medium inhomogeneities

First term of an expansion of the intensity in a multiple scattering series:

$$I = I^0 + I^1 + \cdots + I^n + \dots$$

I^0 : Coherent Intensity

I^n : Mean intensity of waves that have been scattered n times

$$I^1 \sim \frac{l}{t^2} e^{-vt/l}$$

When $vt \ll l$ reduces to the Born Approximation

The Diffusion Approximation

General Idea:

- Each scattering distributes energy over all space directions
- After several scatterings the intensity becomes almost isotropic

$$I(t, \vec{r}, \vec{\Omega}) = \text{Angularly Averaged Intensity} + \\ \text{constant} \times \vec{J}(t, \vec{r}) \cdot \vec{\Omega}$$

The current density $\vec{J}(\vec{r}, t)$, points in the direction of maximum energy flow.
Integrating the RT Eq over all space directions leads to:

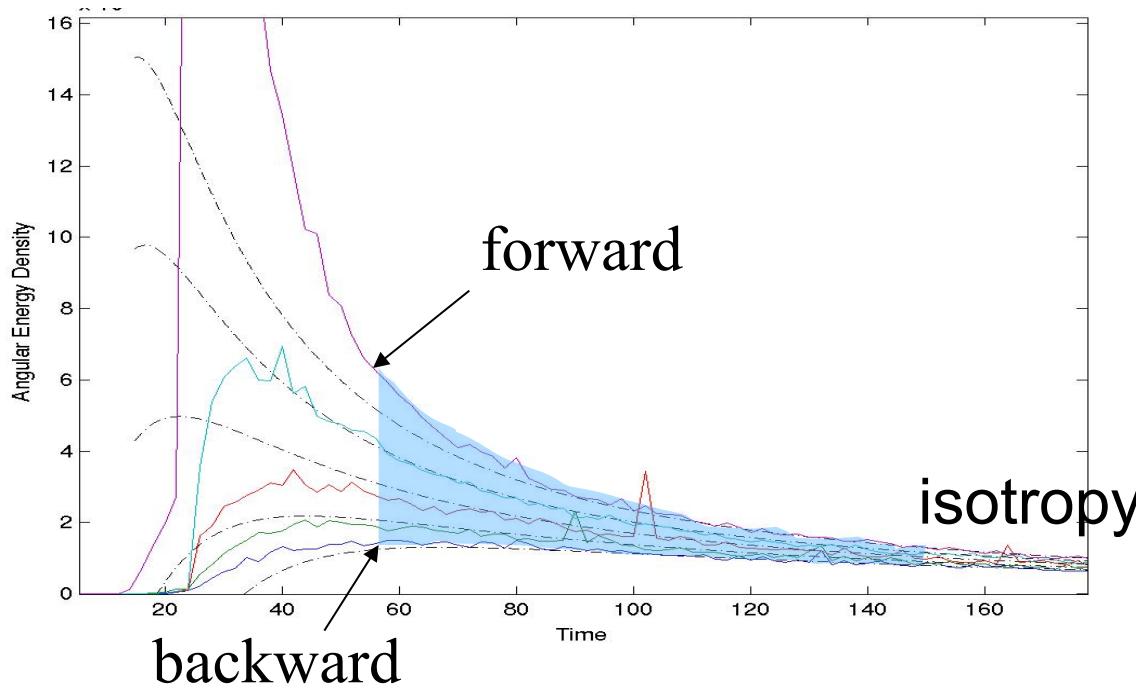
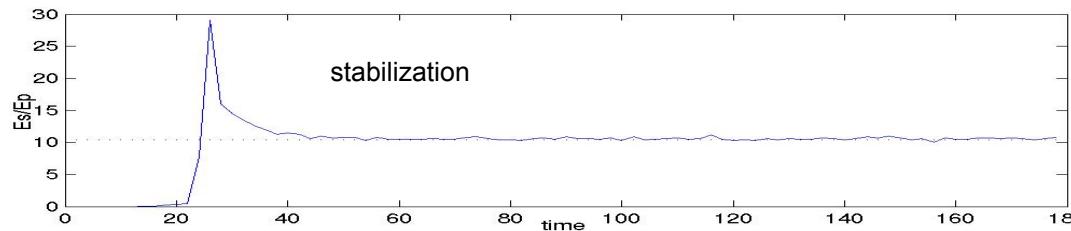
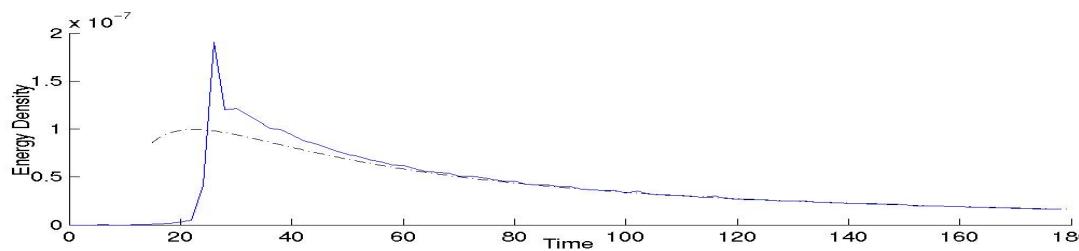
$$\partial_t \rho(t, \vec{r}) - D \nabla^2 \rho(t, \vec{r}) = \delta(t, \vec{r})$$

where ρ is the local energy density.

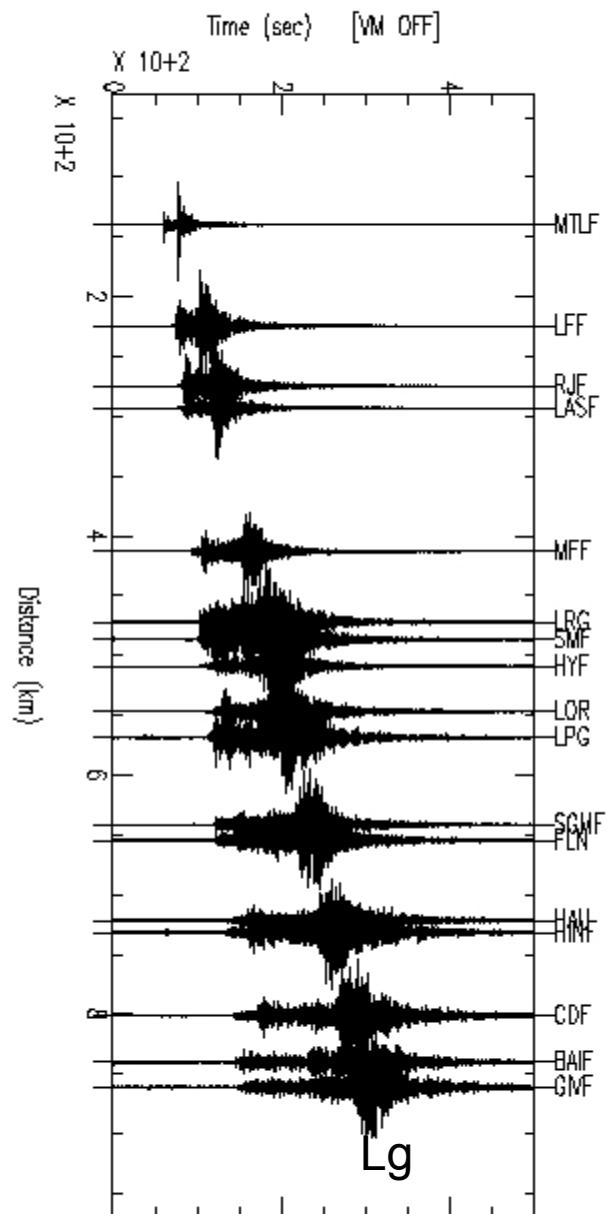
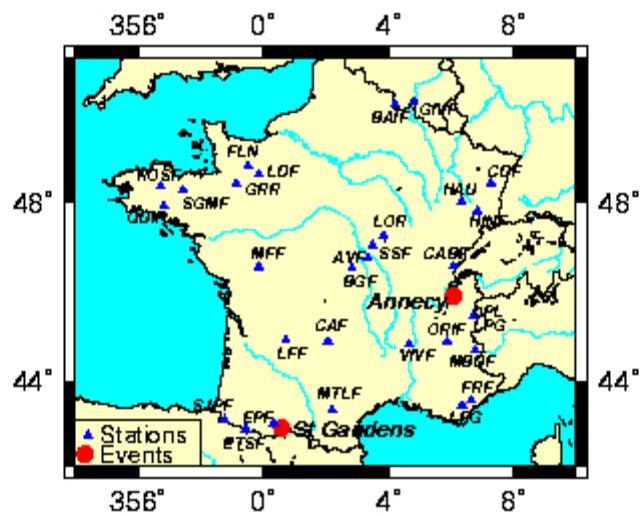
$$\rho(\vec{r}, \vec{r}', t) = \frac{1}{(4\pi Dt)^{d/2}} e^{-|\vec{r}-\vec{r}'|^2/4Dt}$$

$$\rho(t, \vec{r}) \sim \frac{1}{(Dt)^{3/2}} \text{ for large } t.$$

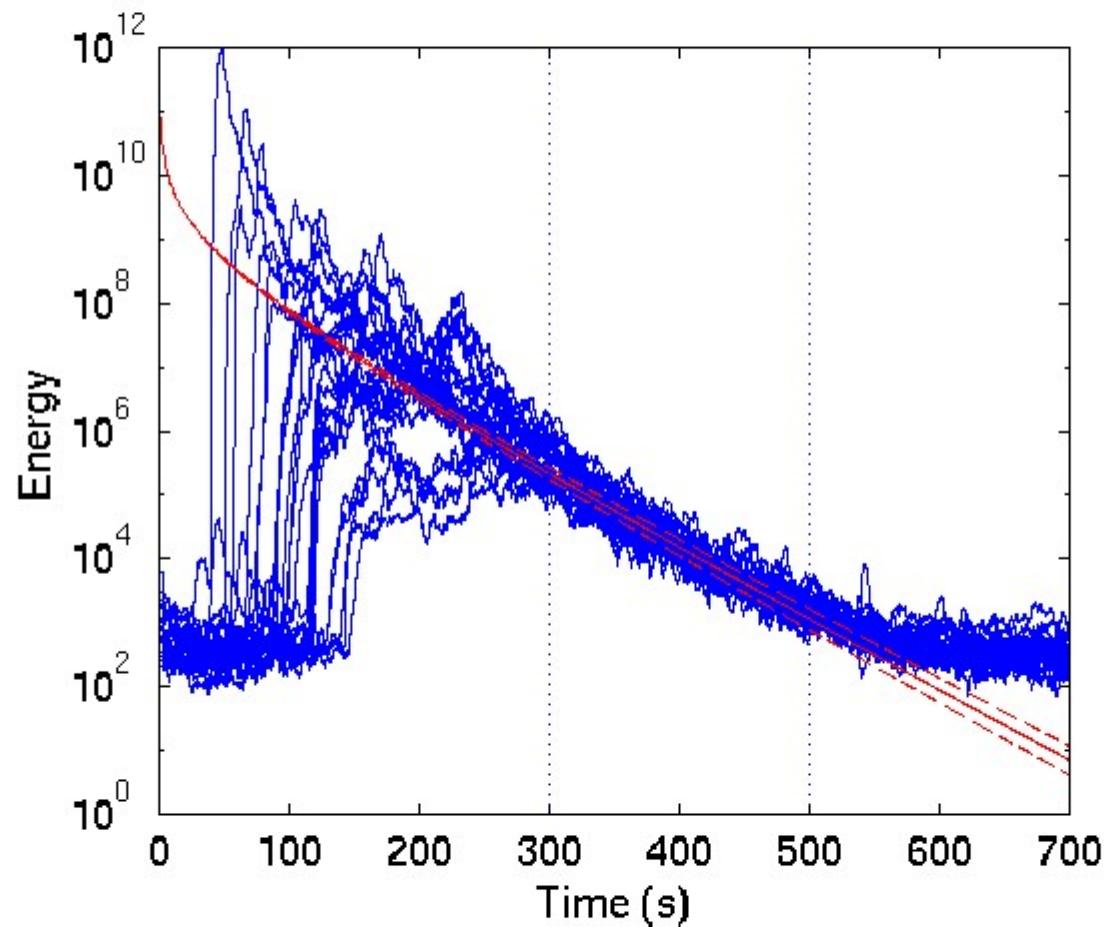
$D = vl/3$ is the diffusion constant of the waves.

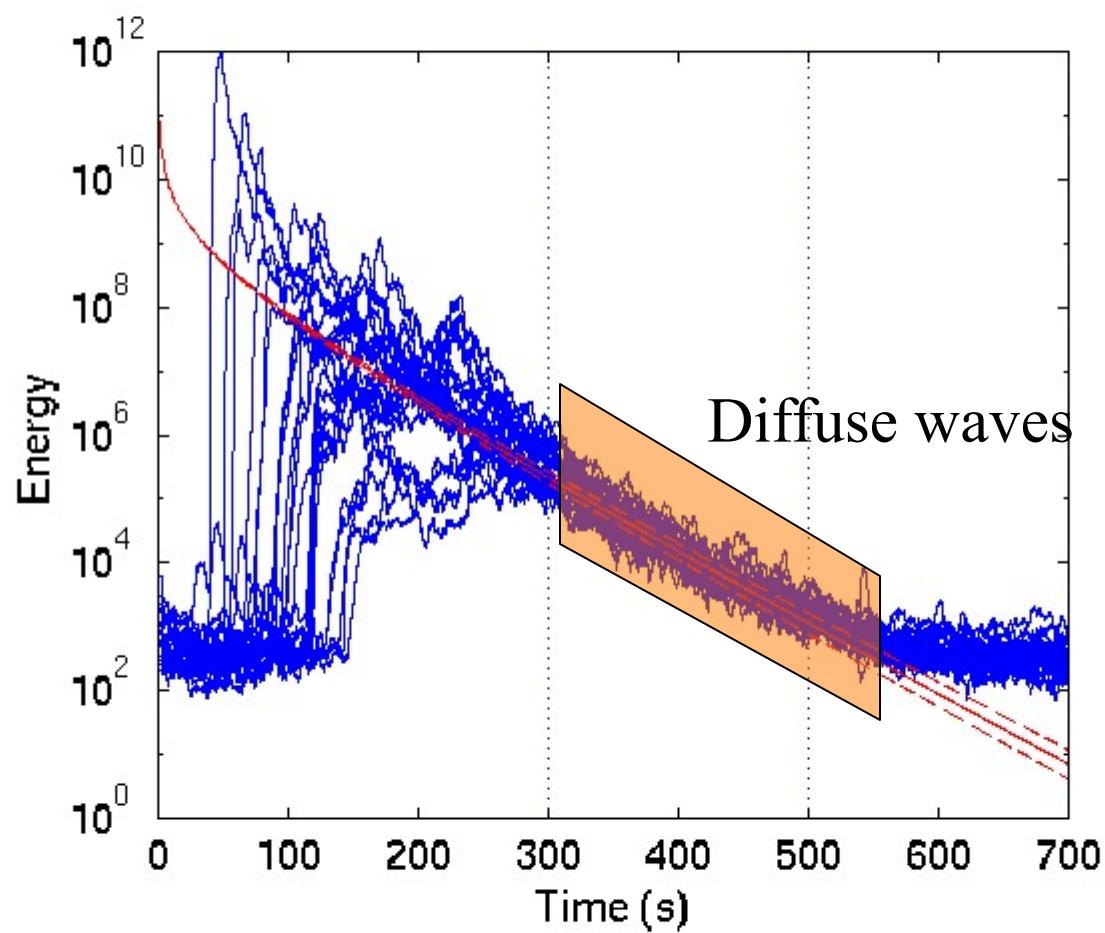


Regional seismograms



Observations at distances between 150 and 800 km!!



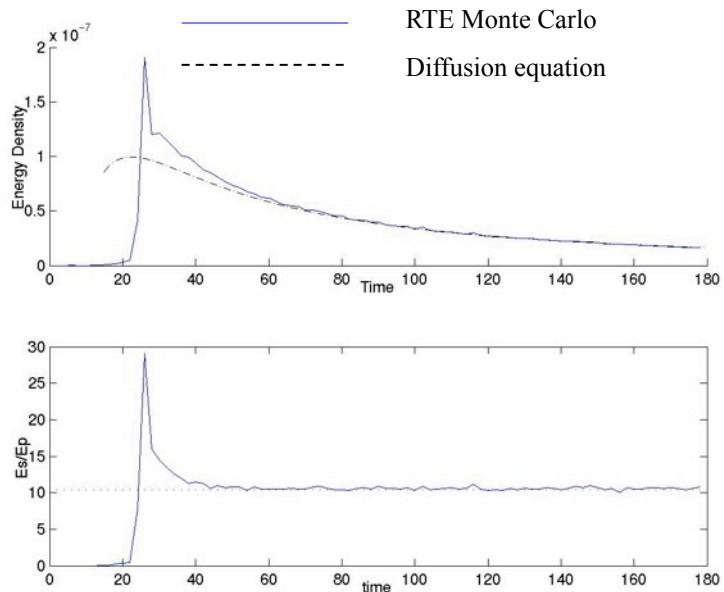


Searching for a marker of the regime of scattering...

Equipartition principle for a completely randomized (diffuse) wave-field: in average, all the modes of propagation are excited to equal energy.

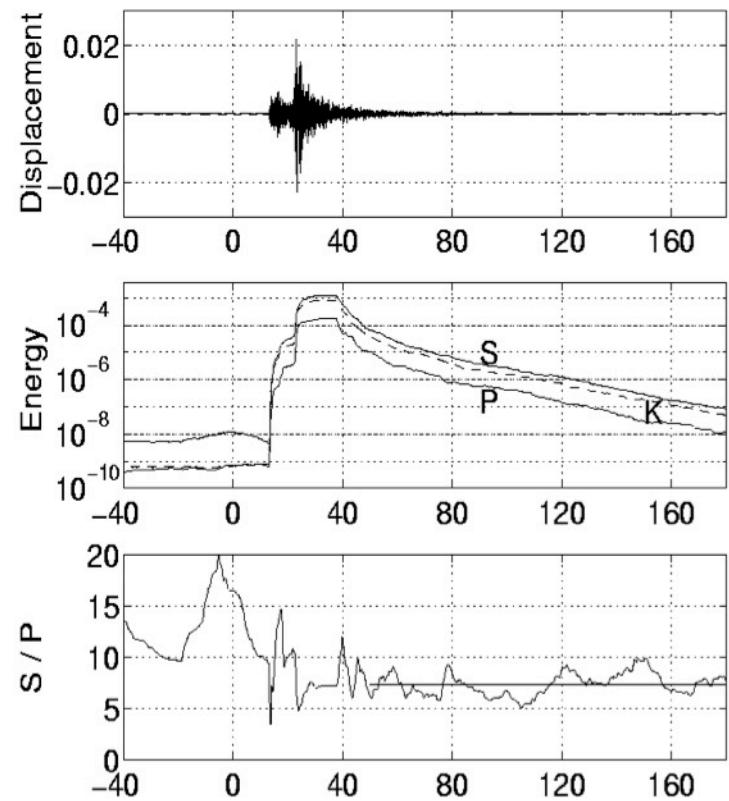
Implication for elastic waves (Weaver, 1982, Ryzhik et al., 1996): P to S energy ratio stabilizes at a value independant of the details of scattering!

Numerical simulation

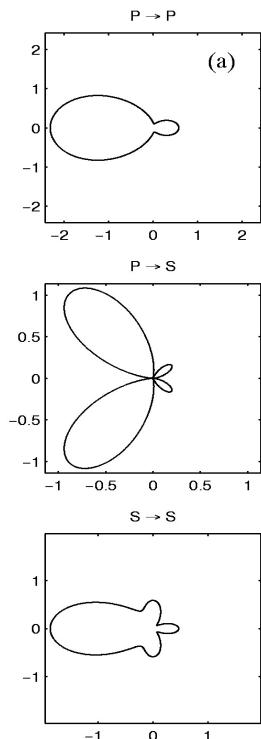


Observations

Event 11

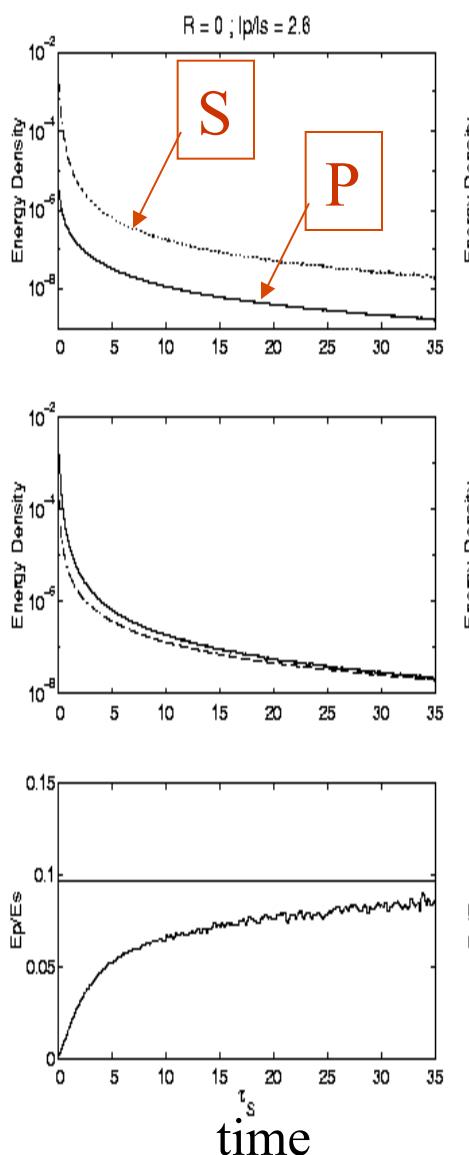


Cross sections



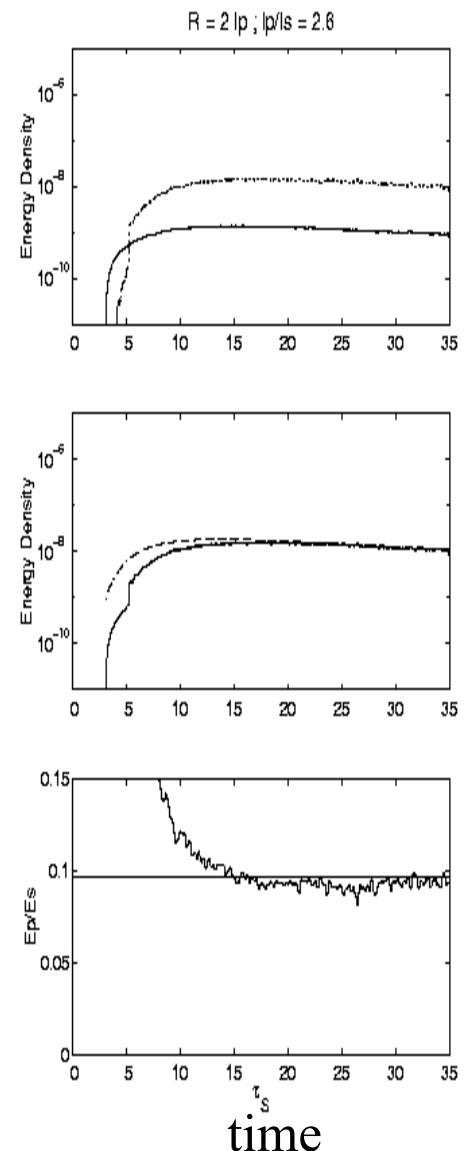
$ka \rightarrow 0$: Rayleigh approximation

Numerical solutions of the RTE



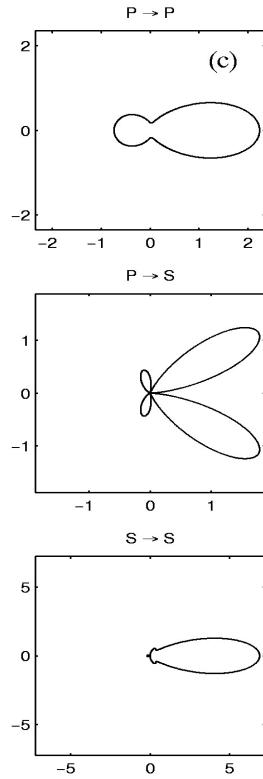
total energy
vs
diffusion app.

Energy ratio

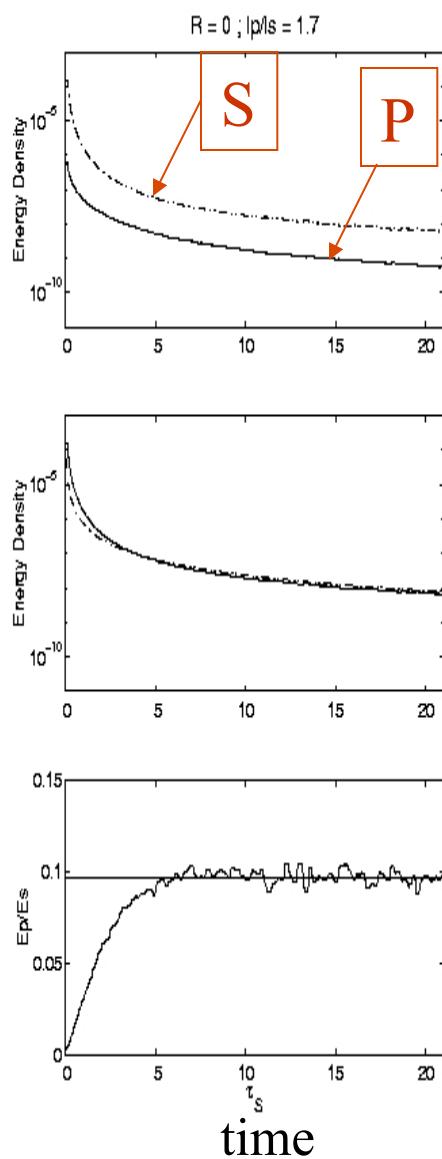


Cross sections

RTE solutions (Monte Carlo)

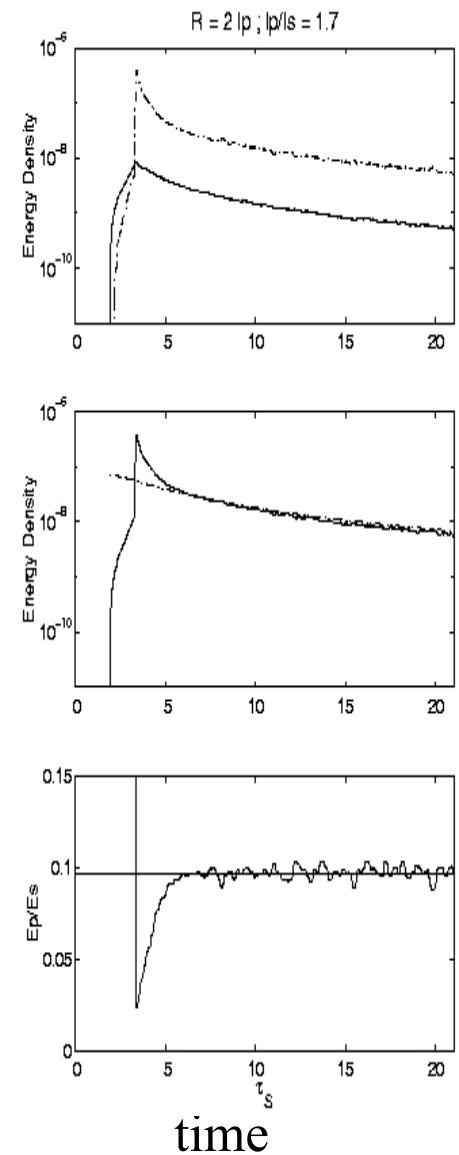


$ka=1.6$

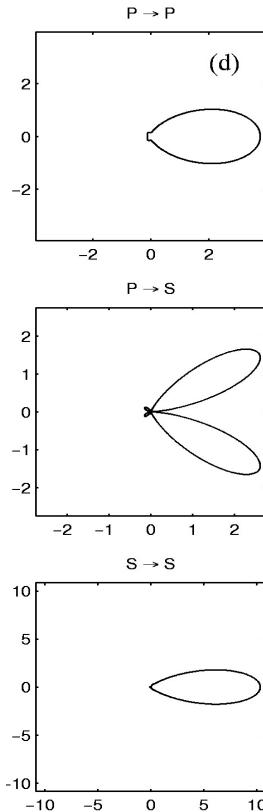


total energy
vs
diffusion app.

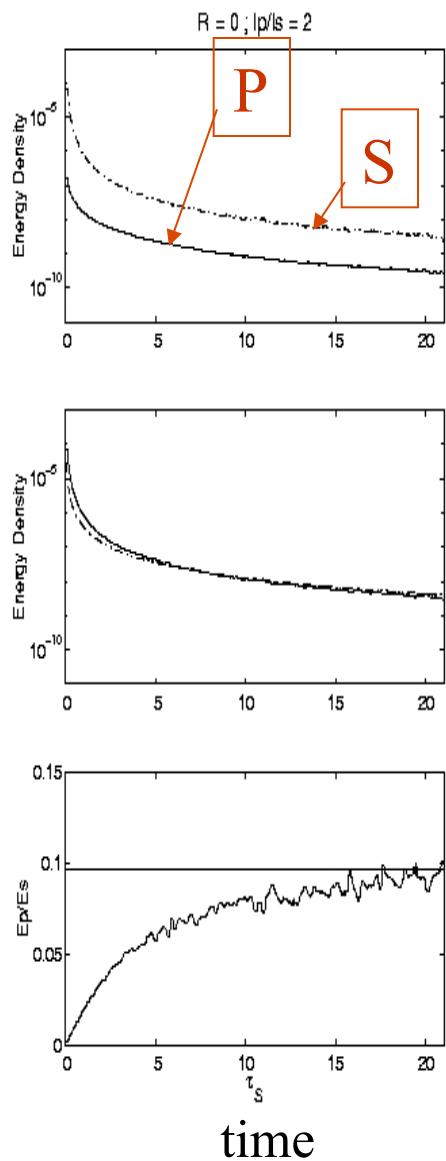
Energy ratio



Cross sections

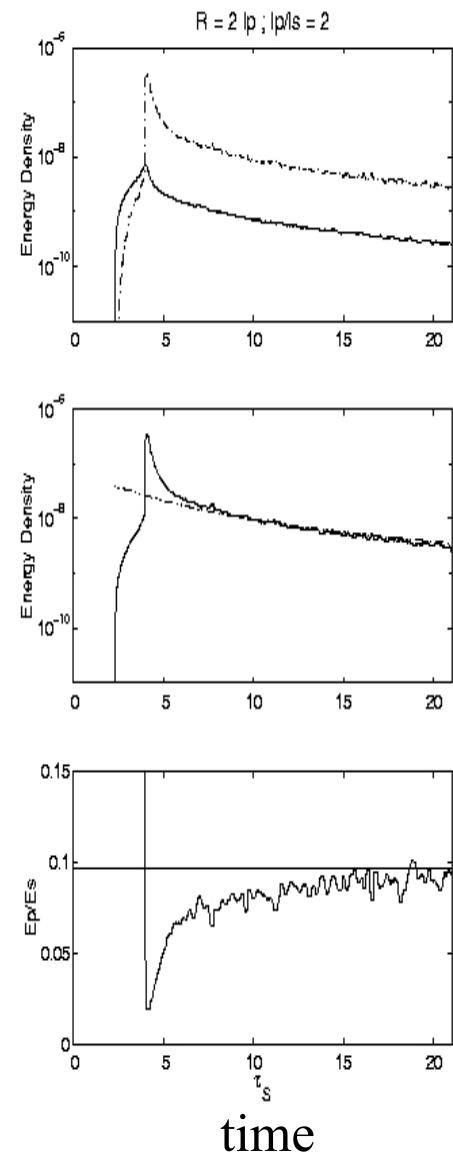


$ka=2.0$



total energy
vs
diffusion app.

Energy ratio
is the same for
different
values of ka



Energy in an Elastic Solid

$$E = K + P + S + I$$

$$E = \frac{1}{2}\rho(\partial_t \mathbf{u})^2 + \left(\frac{\lambda}{2} + \mu\right)(\operatorname{div} \mathbf{u})^2 + \frac{\mu}{2}(\operatorname{curl} \mathbf{u})^2 + I$$

I contains mixed partial derivatives
 $K = H^2 + V^2$

Focus on the ratios:
 $P/S, K/(P + S), I/(S + P), H^2/V^2$

Equipartition predicts: Any Ratio of Energies Becomes Independent of Time

Measurement of the deformation energy requires evaluation of partial derivatives of the wavefield

S to P Energy ratio as a marker of the regime of scattering...

$$G_{i,j}(\vec{R}, \vec{S}, t) = \sum_n \varepsilon_n \Phi^n(\vec{R}) \exp(-i\Omega_n t)$$

where ε_n are random independent variables (finite body)

(ldos)

Consequence for an infinite inhomogeneous solid:

$$\frac{E_s}{E_p} = 2 \left(\frac{v_p}{v_s} \right)^3$$

Independent of the Details of the Scattering !

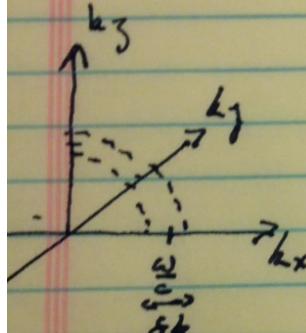
Independent of the position in a full space with homogeneous reference

Partition of energy (Full elastic space)

Multiphon scattering, large t
→ "equipartition"
[reference medium + disorder]

Phase space of the full space elastic problem

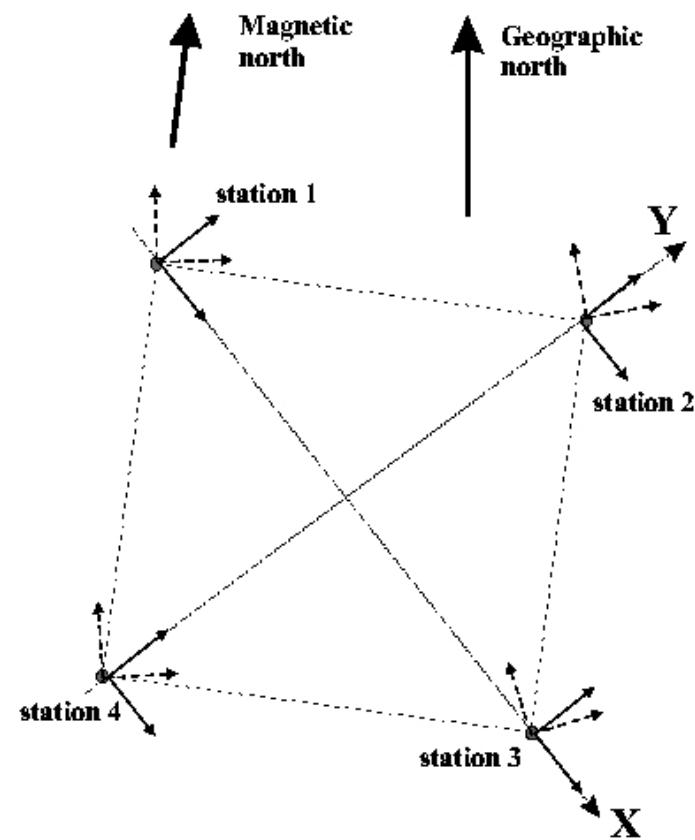
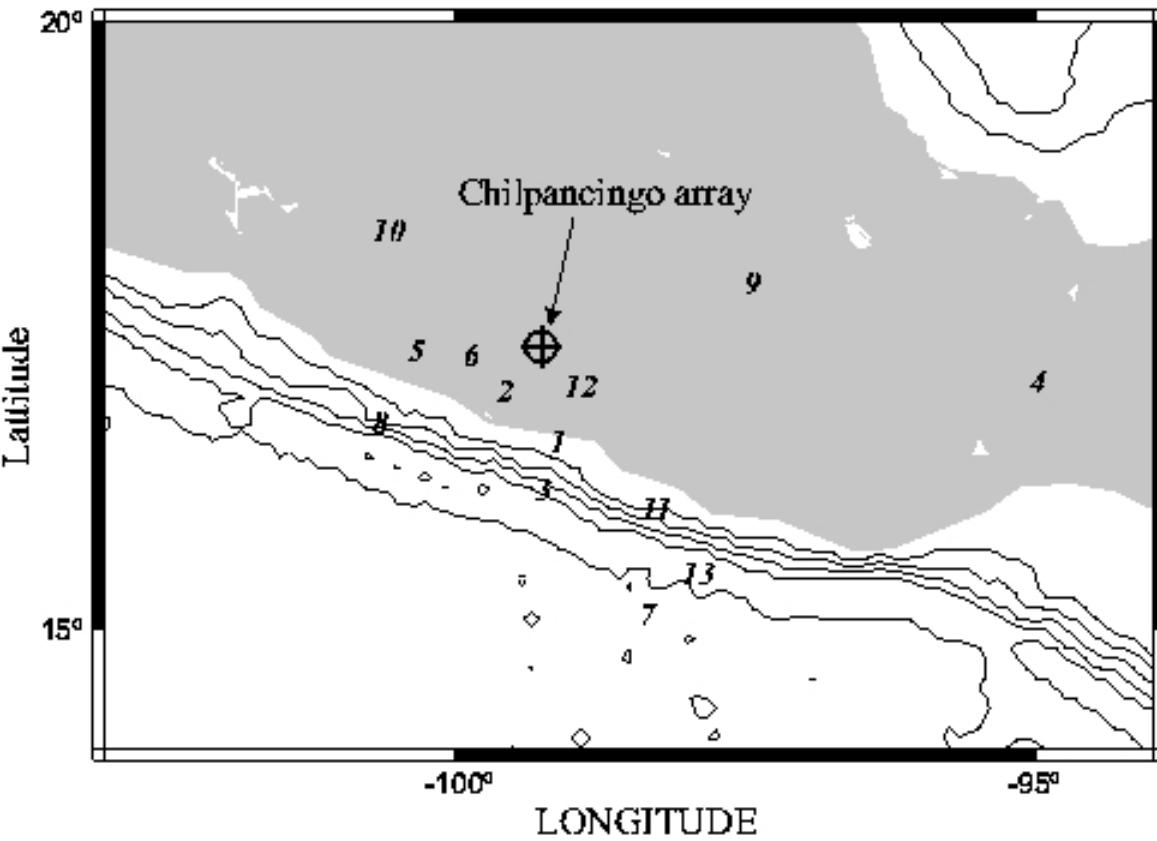
→ all propagating plane waves
excited at same level of energy



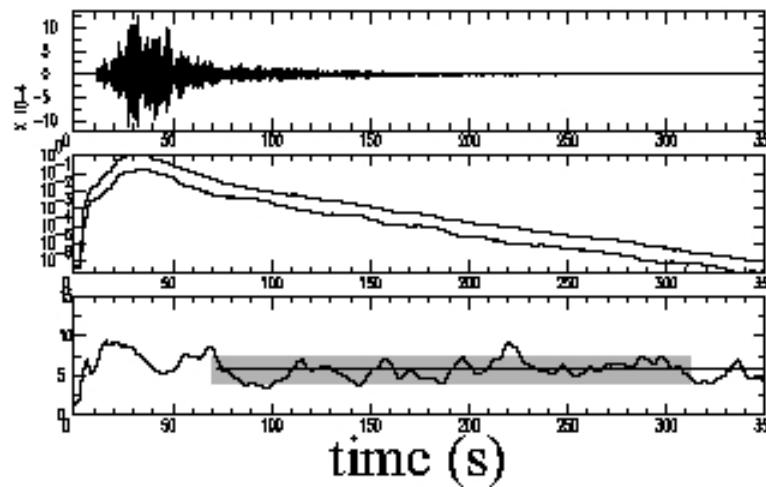
$$\text{Energy in a band } \omega \pm \frac{\delta\omega}{2}$$

→ Volume for P-waves

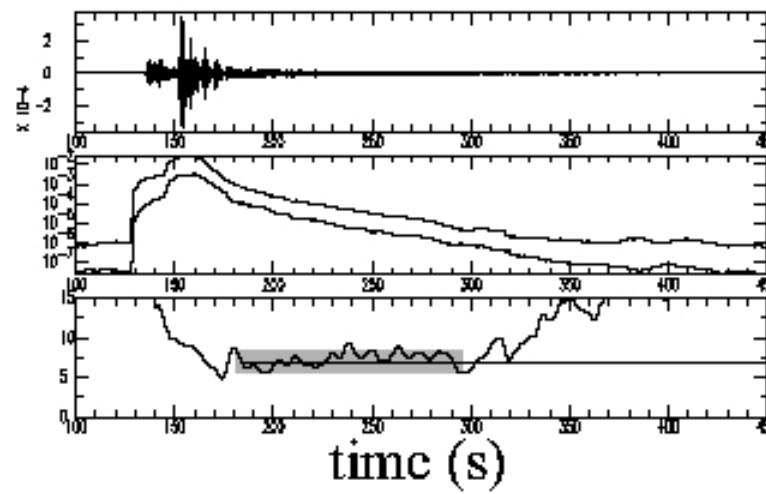
$$\delta k = \frac{\delta\omega}{c}$$



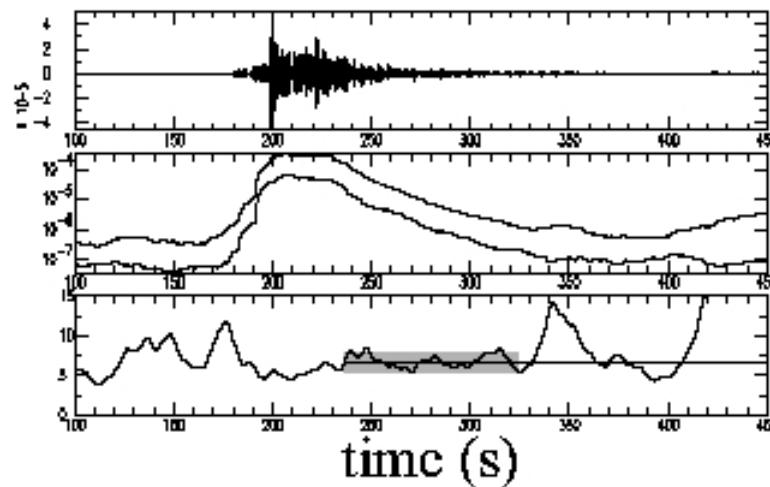
event 5



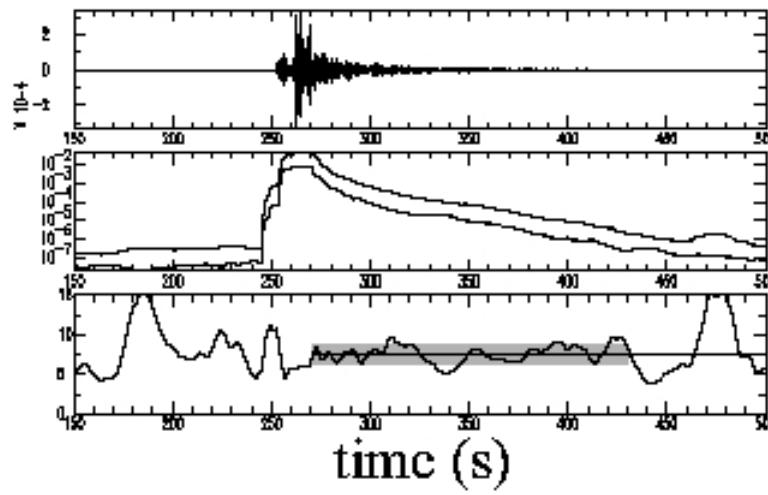
event 8



cvcnt 9

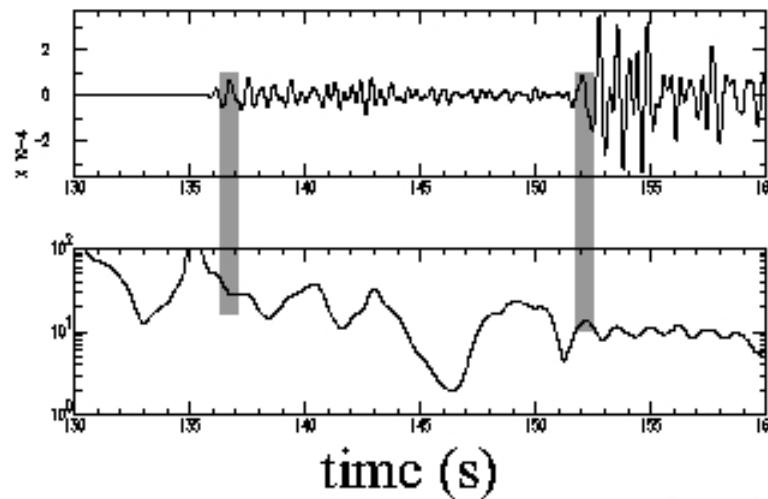


cvcent 12

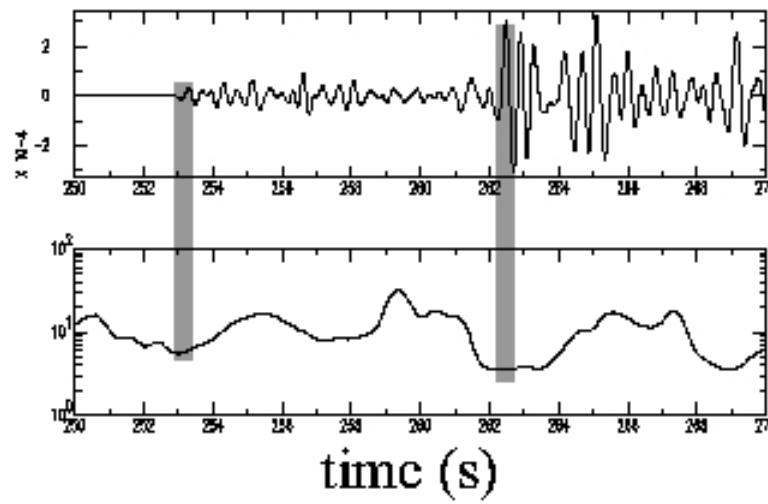


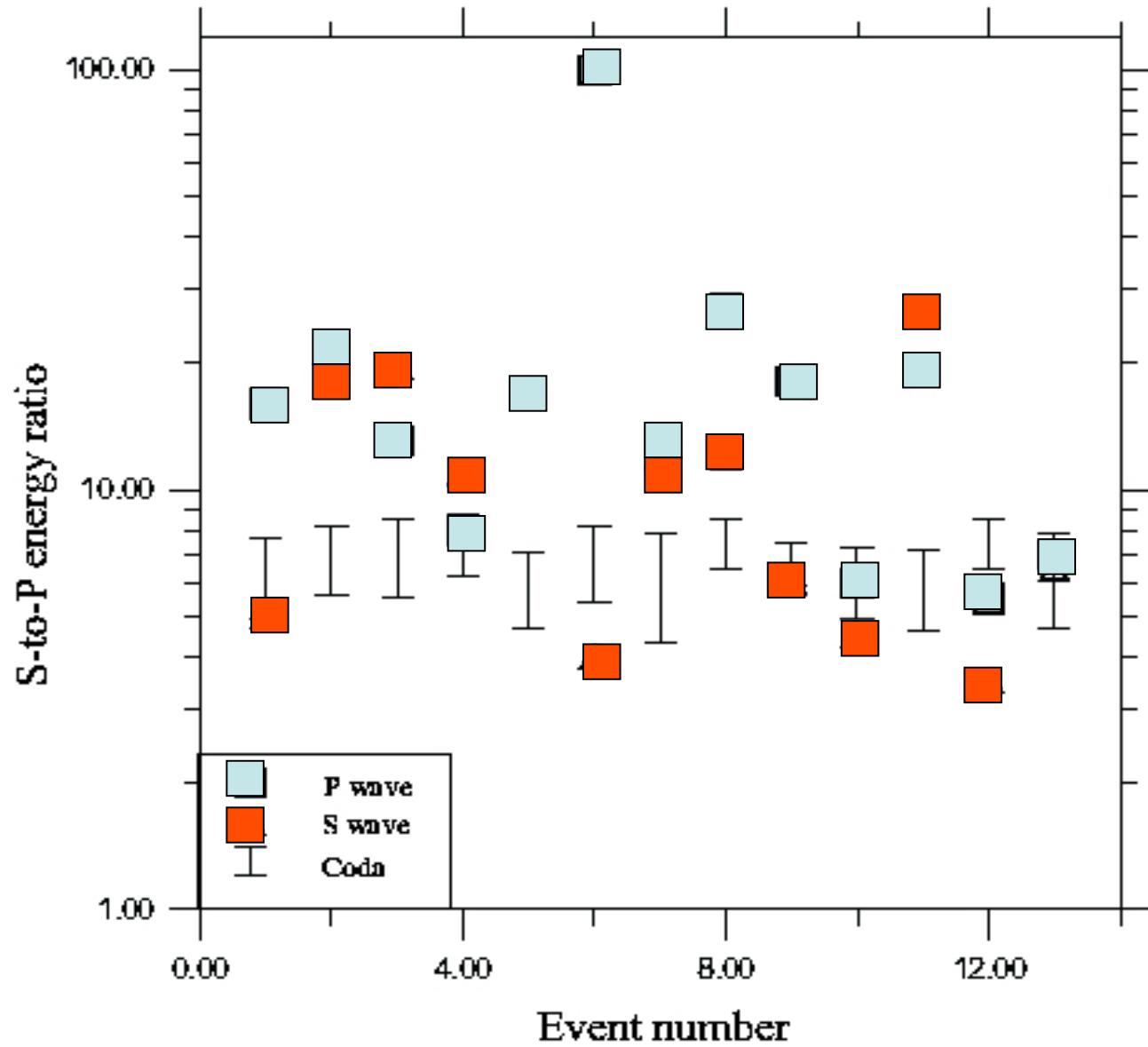
Direct waves

event 8



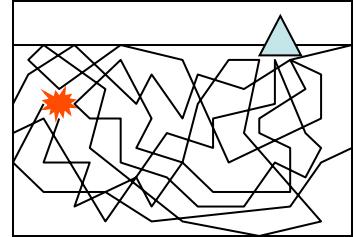
event 12





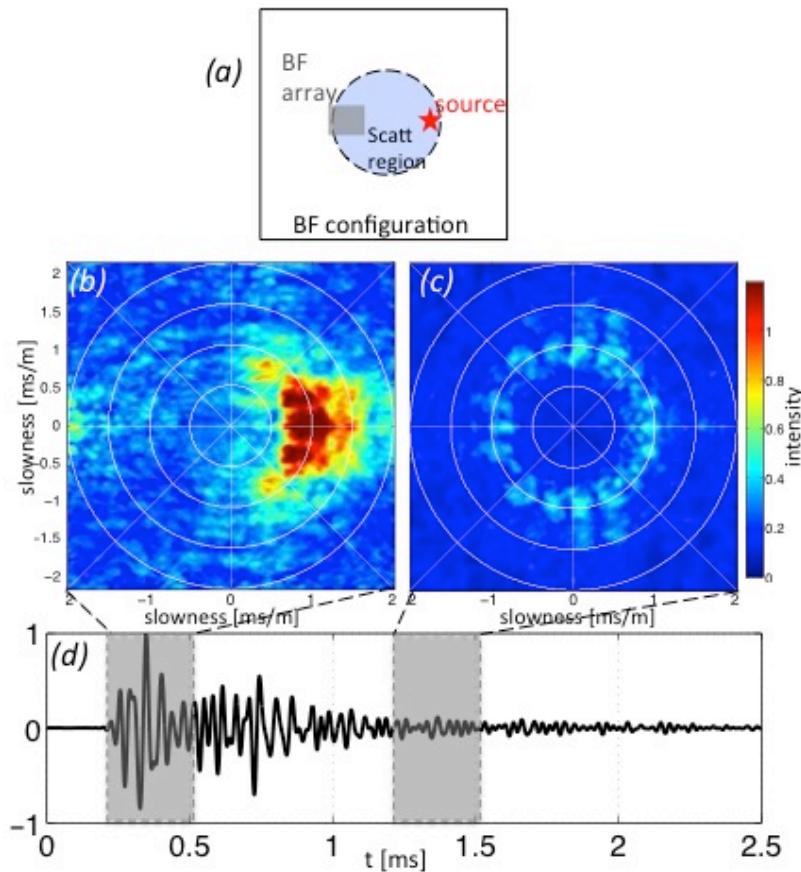
ENERGY RATIO	DATA	THEORY FULL SPACE	THEORY HALF SPACE BULK WAVES	THEORY HALF SPACE with RAYLEIGH WAVES
S/P	7.3	10.39	9.76	7.19
K/(S+P)	0.65	1	1.19	0.534
V(S+P)	-0.62	0	-0.336	-0.617

For asymptotically long lapse time (diffusion), the disorder produces a completely randomized wave-field. such that , all the modes of propagation are excited in average to equal energy (the equipartition principle).

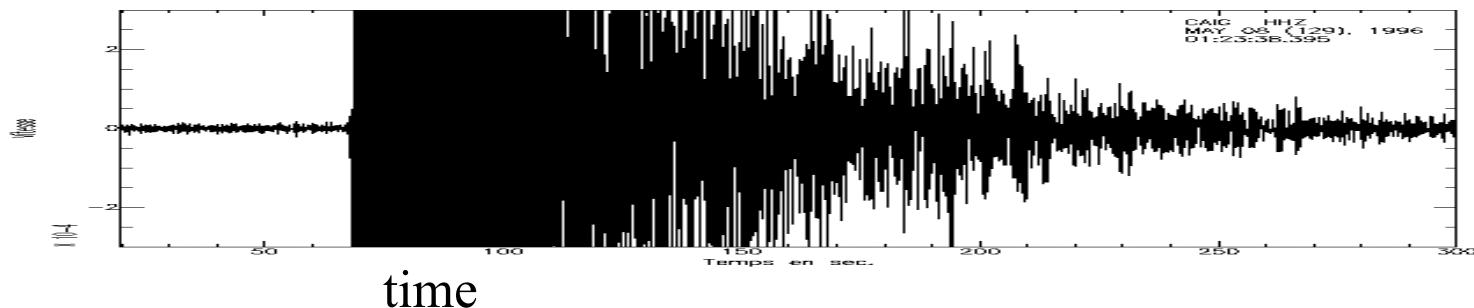
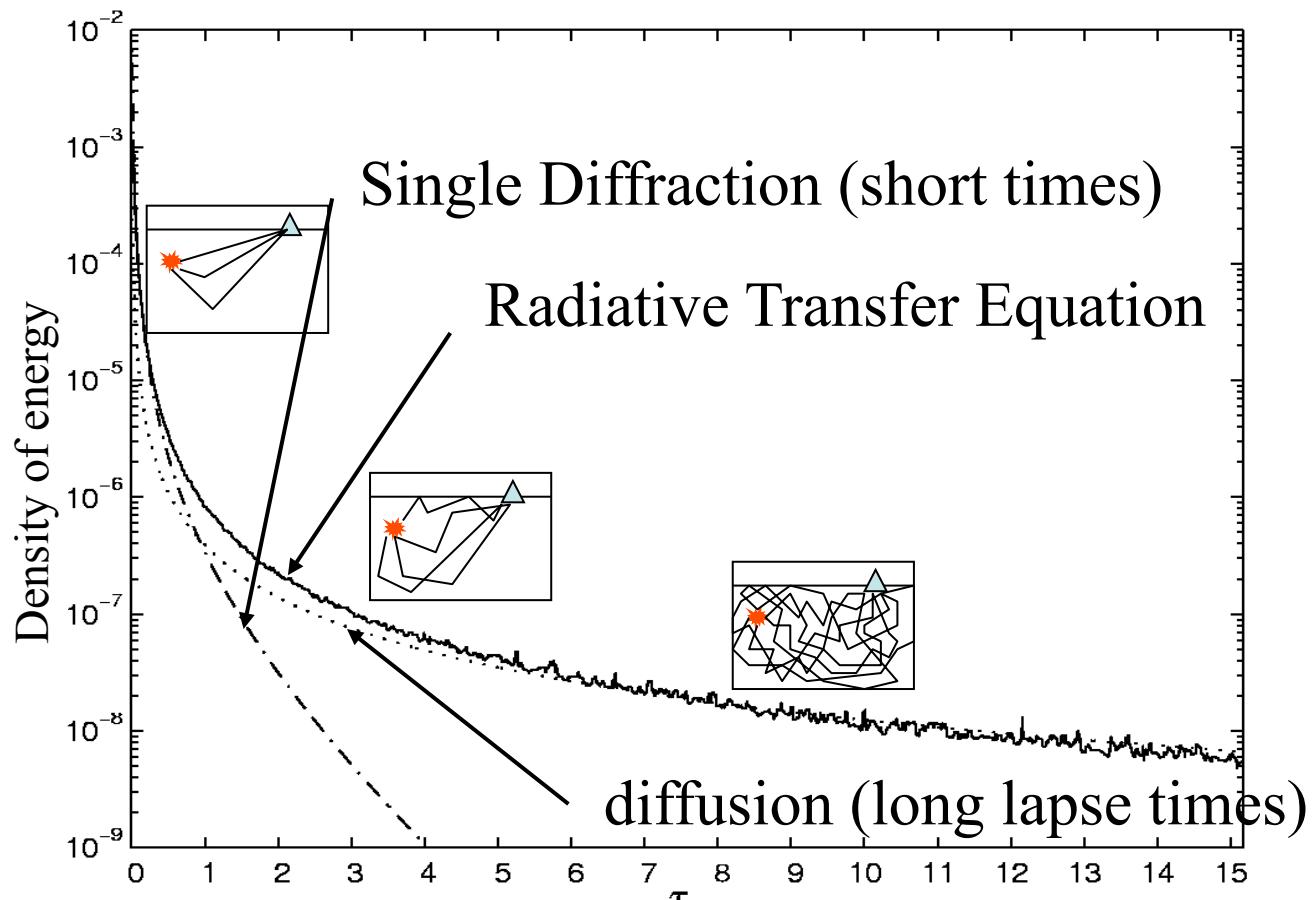


Numerical example:
2D scalar waves

→intensity isotropy



Propagation regimes and description of energy



The Diffusion Approximation

General Idea:

- Each scattering distributes energy over all space directions
- After several scatterings the intensity becomes almost isotropic

$$I(t, \vec{r}, \vec{\Omega}) = \text{Angularly Averaged Intensity} + \\ \text{constant} \times \vec{J}(t, \vec{r}) \cdot \vec{\Omega}$$

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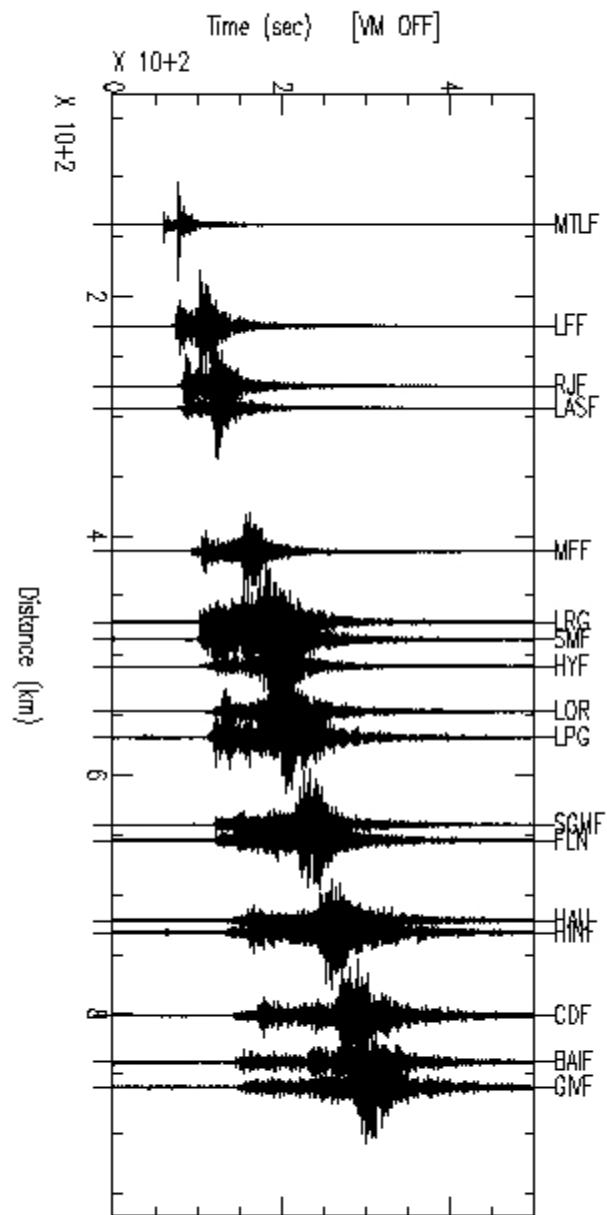
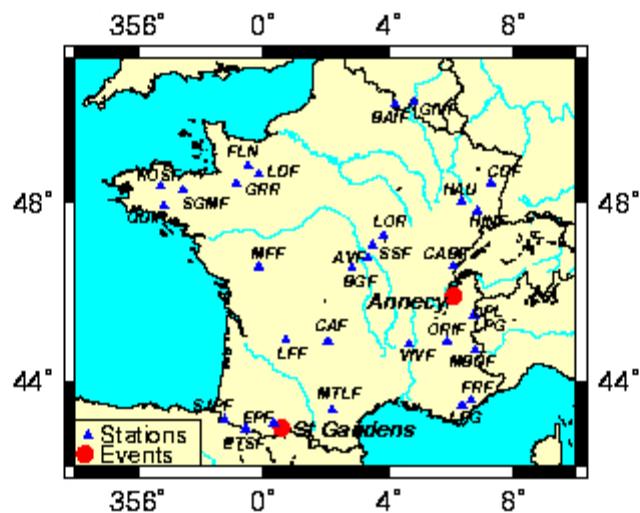
$$\partial_t \rho(t, \vec{r}) - D \nabla^2 \rho(t, \vec{r}) = S(t, \vec{r})$$

where ρ is the local energy density.

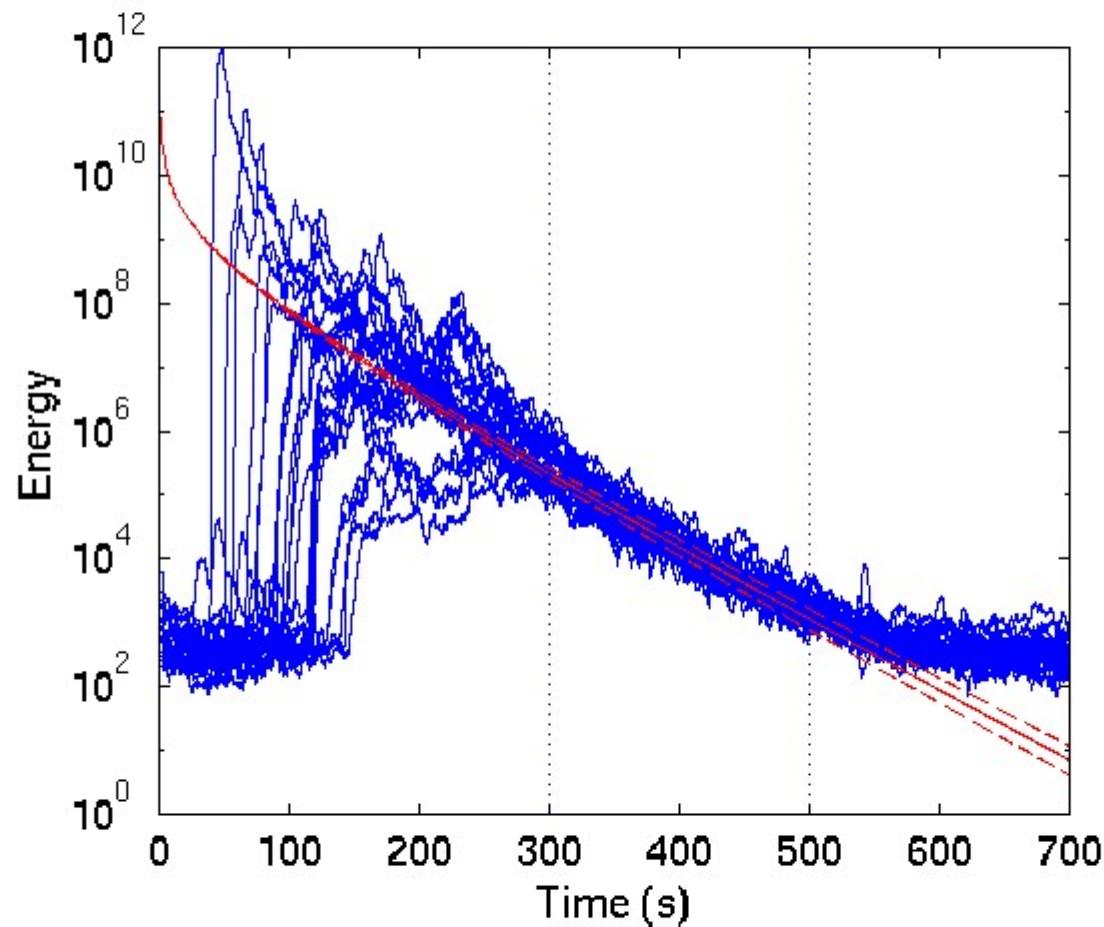
$$\rho(t, \vec{r}) \sim \frac{1}{(Dt)^{3/2}} \text{ for large } t.$$

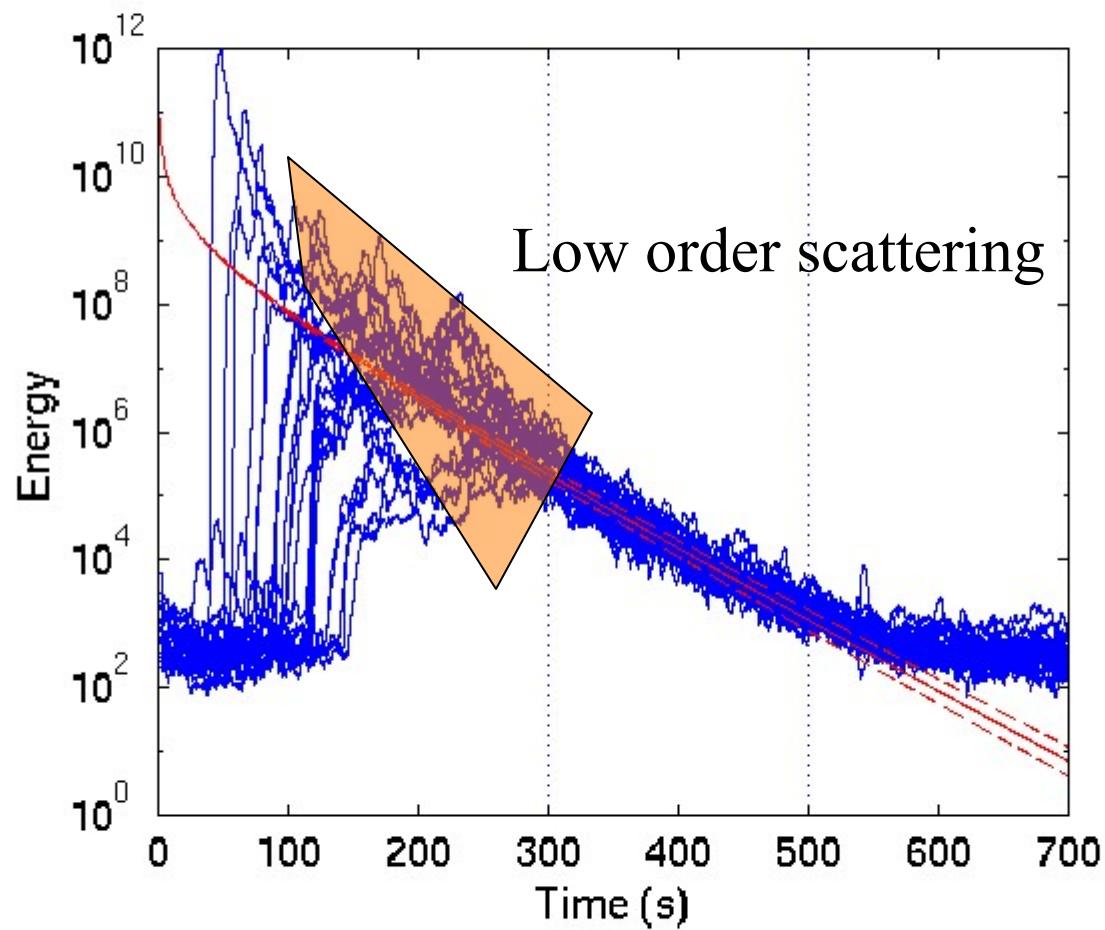
$D = v l / 3$ is the diffusion constant of the waves.

Coda of regional seismograms



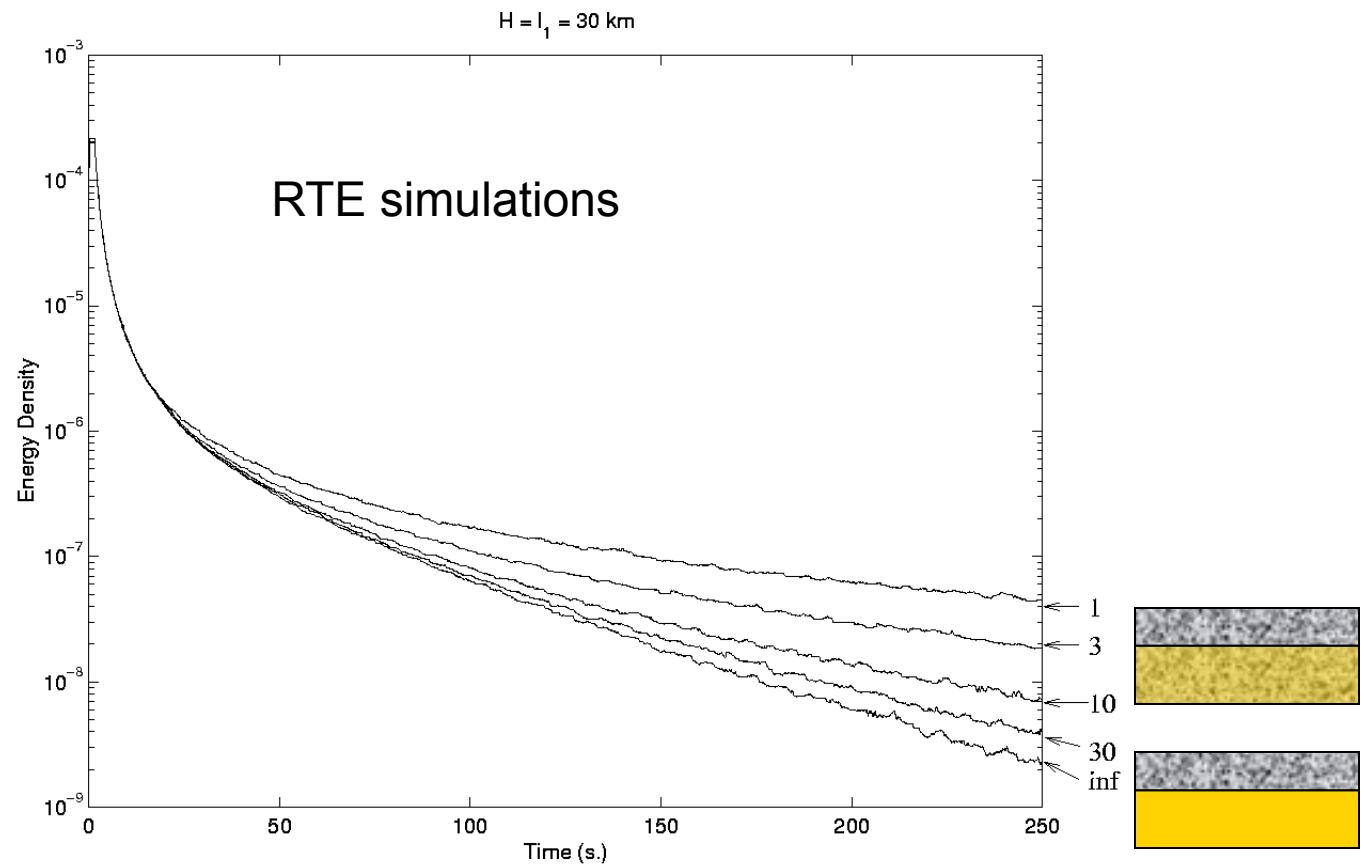
Observations at distances between 150 and 800 km!!



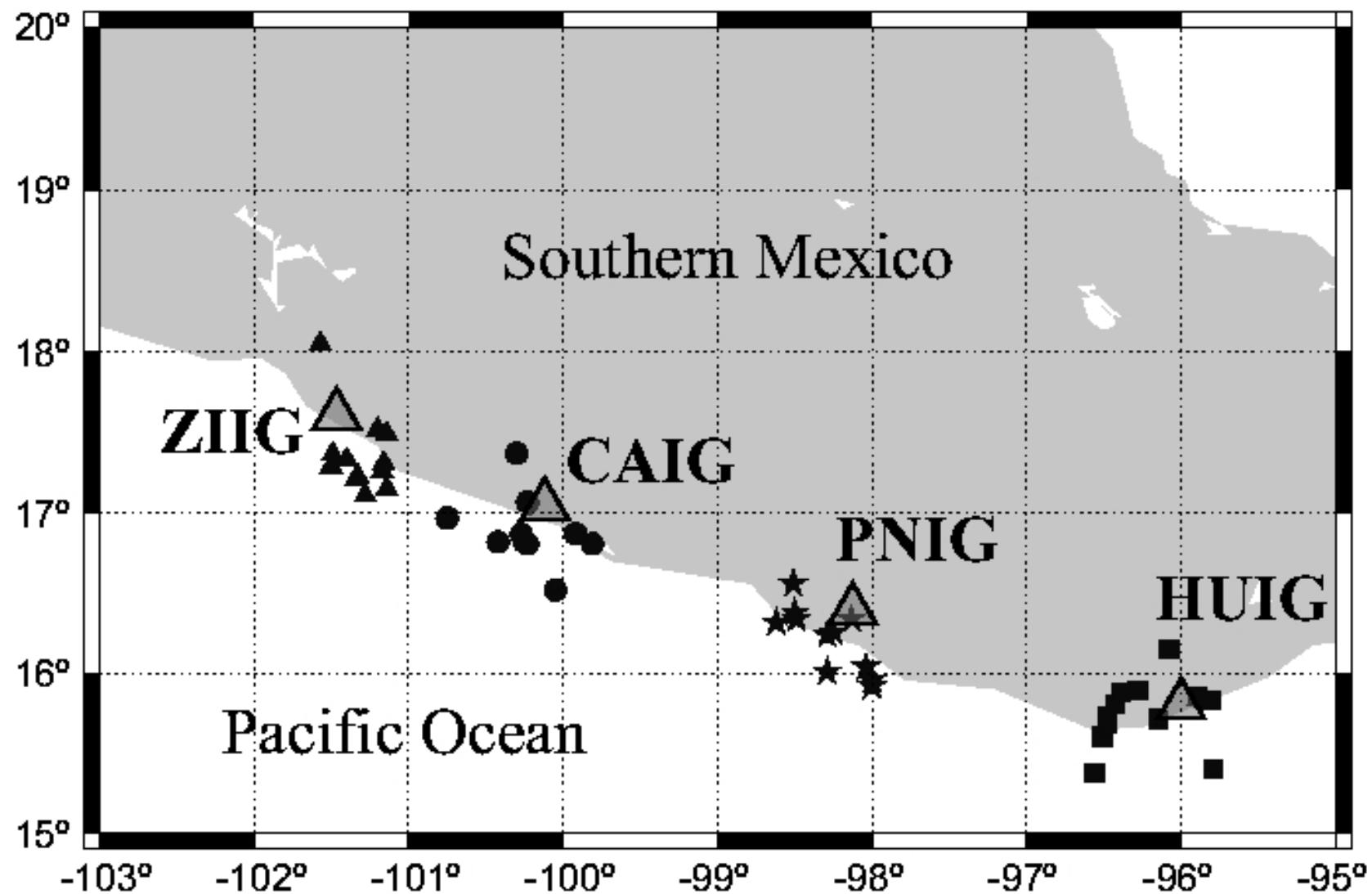


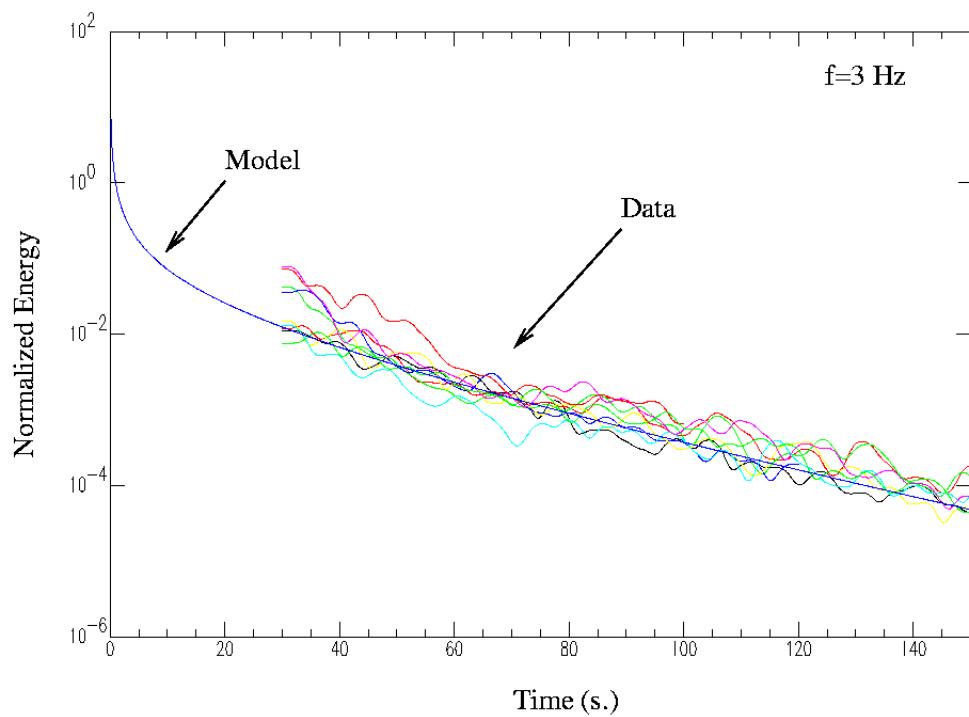
requires radiative transfer equation

Influence of the value of mantle mean free path



Leakage of energy in the mantle

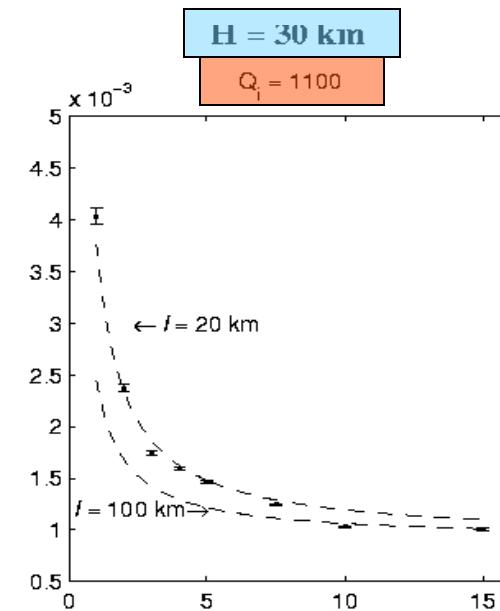
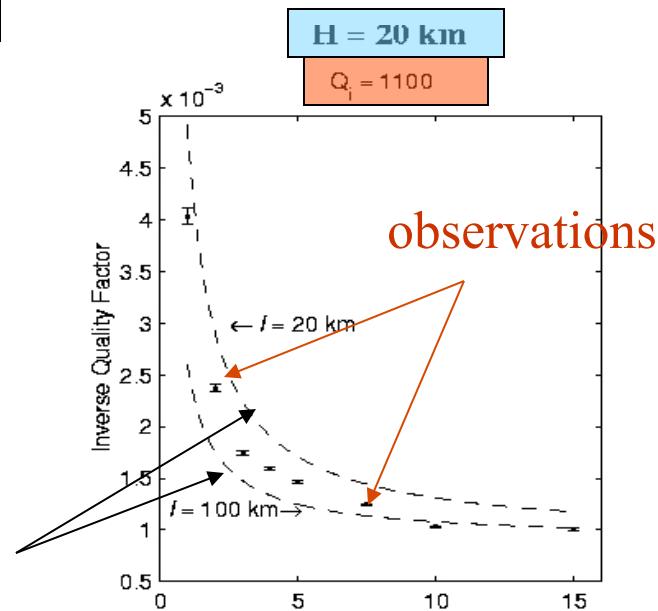




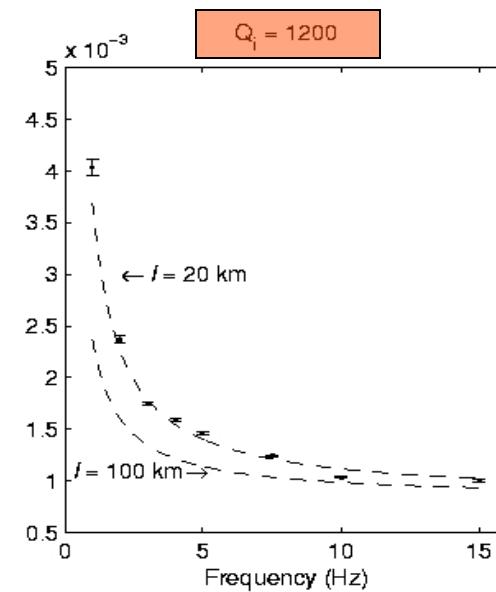
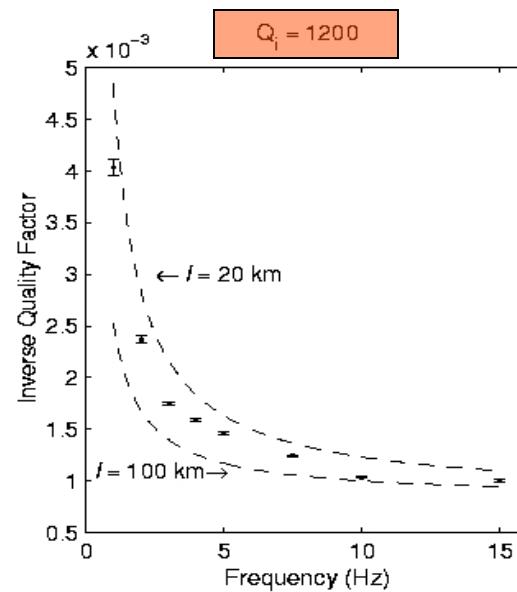
With dissipation

$Q > 1000$

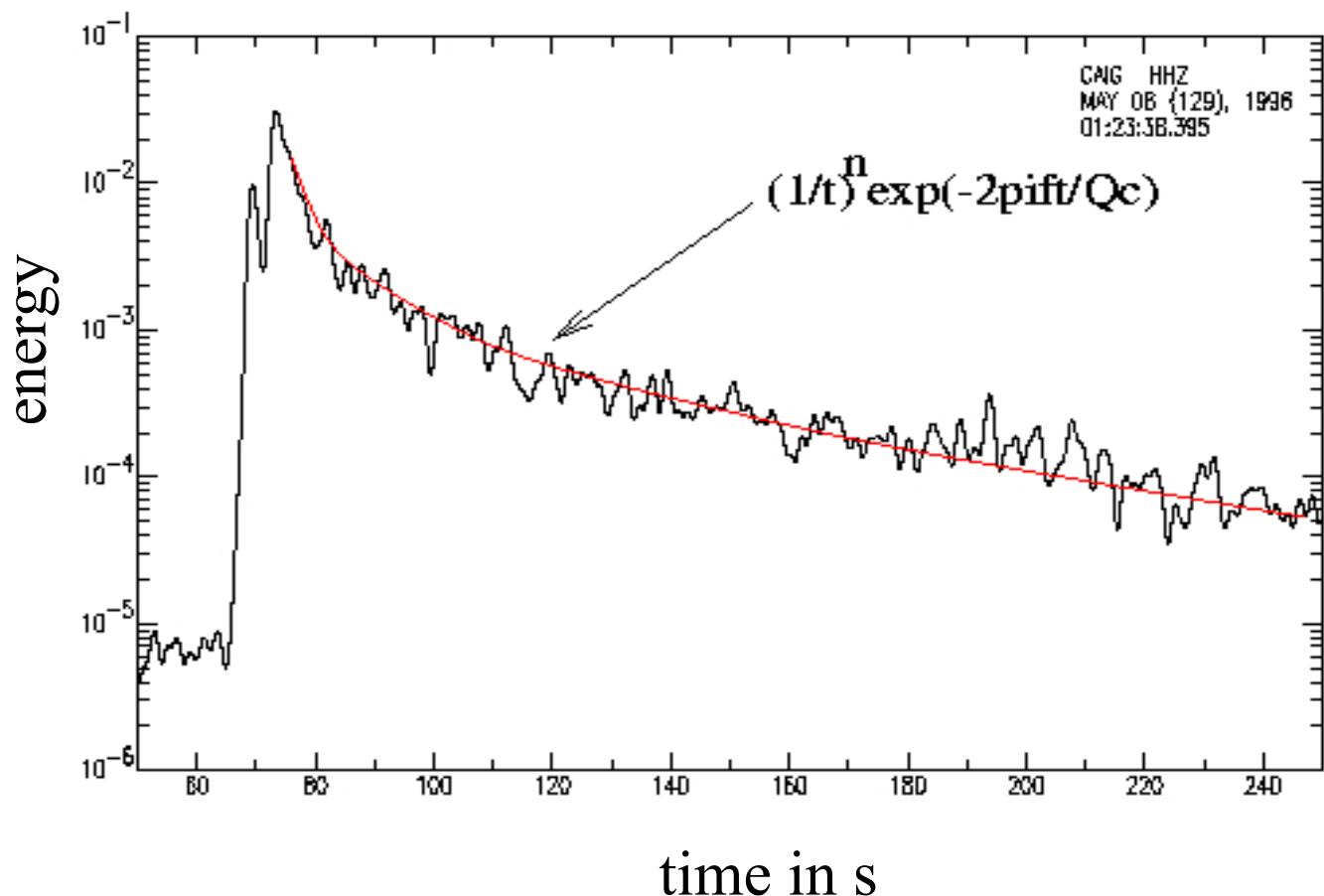
RTE solutions



$l(f)$ in the range
20-80 km

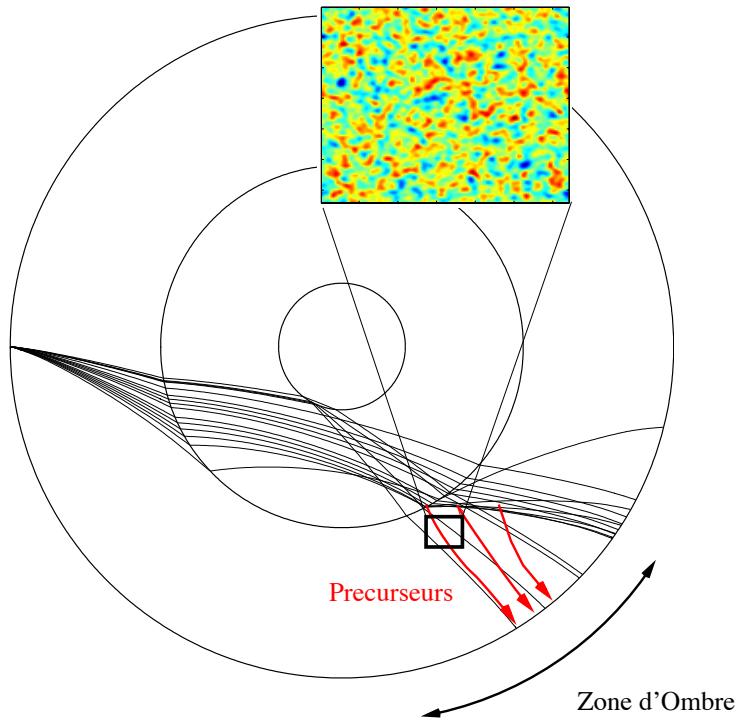


Energy decay in the coda (Aki and Chouet, 1975)

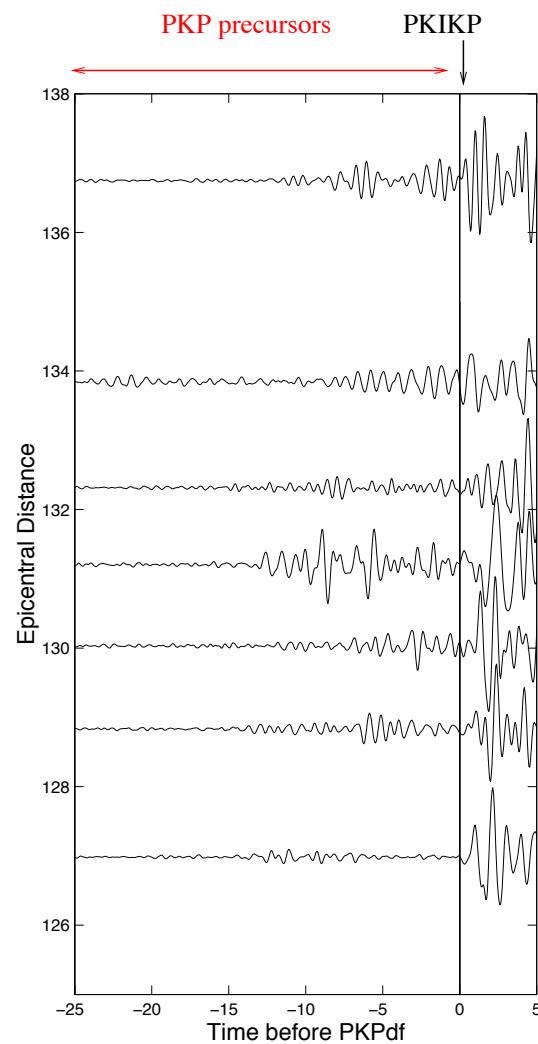


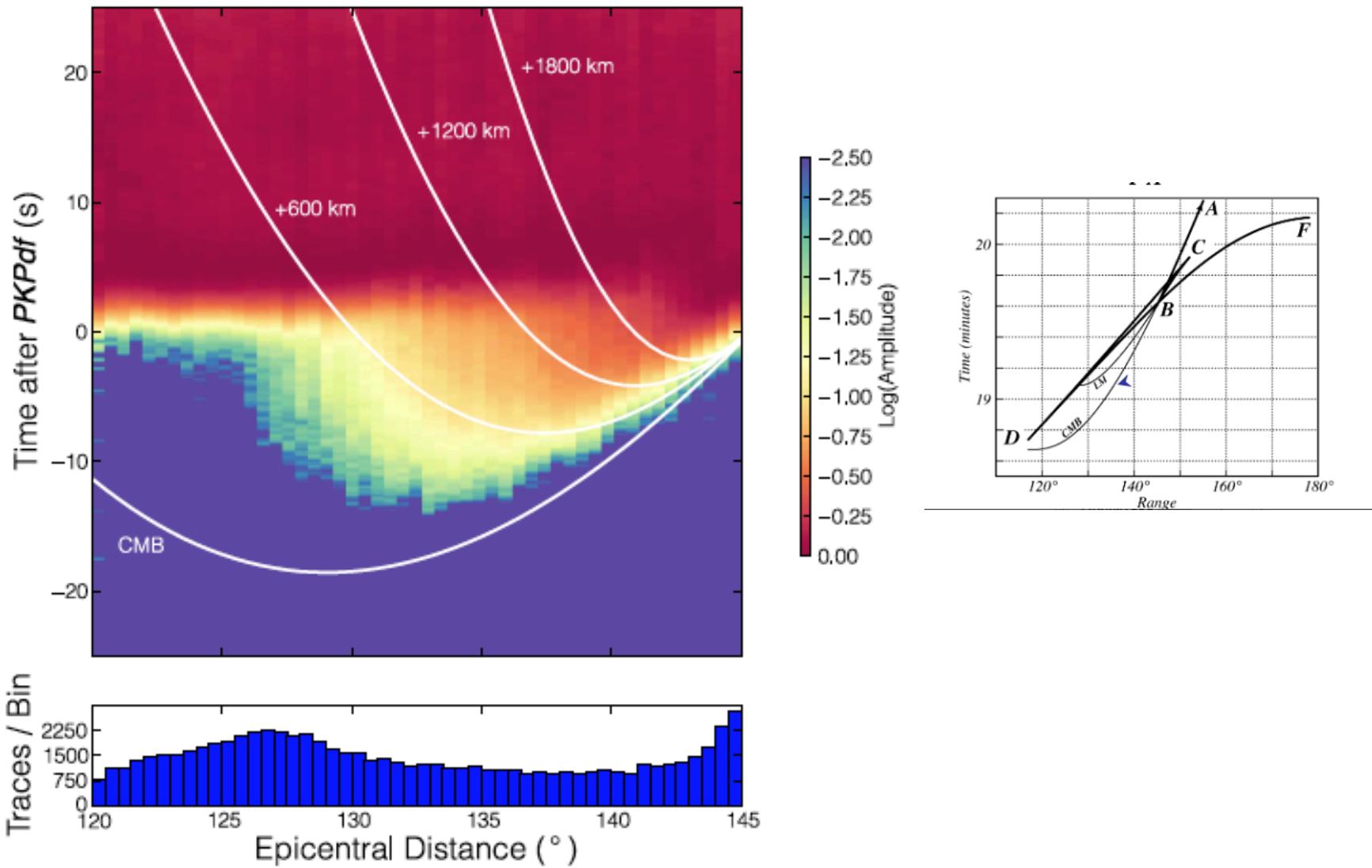
Hétérogénéité à grande profondeur

Précurseurs des PKIKP



Margerin & Nolet, 2003



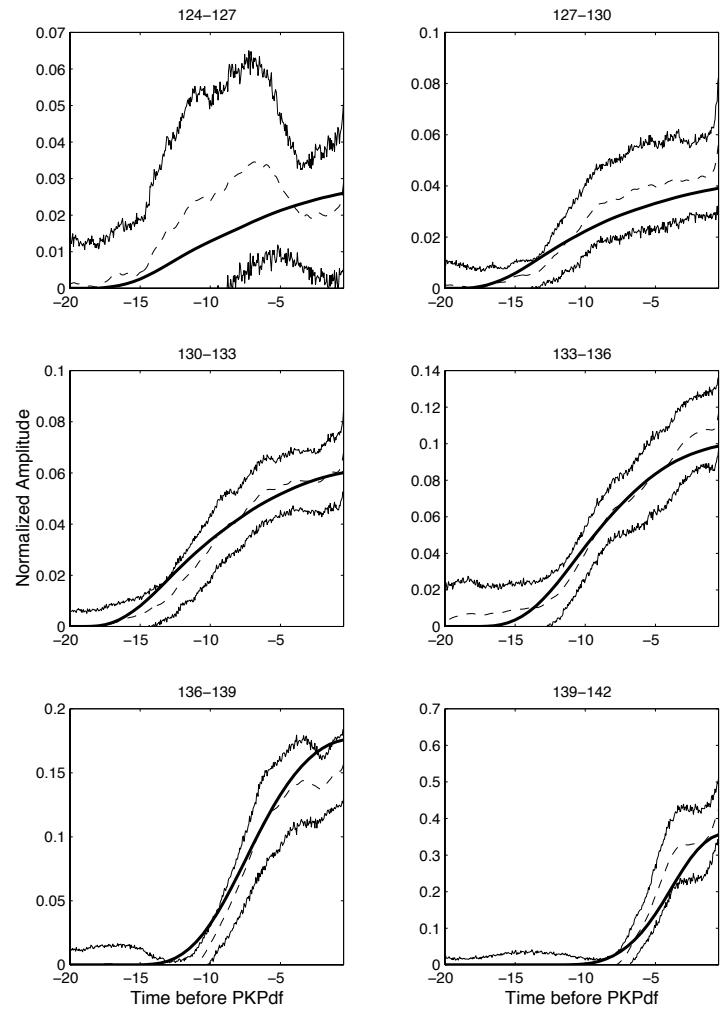


N. J. Mancinelli and P. M. Shearer

Fit of the observations with RTE and heterogeneity in the whole mantle described by a Henyey and Greenstein (H-G) correlation function.

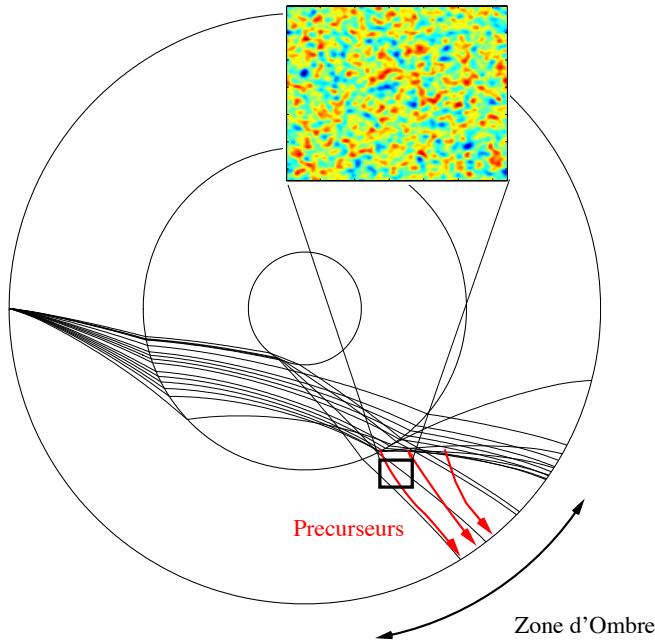
Perturbations are of the order of 0.1%, correlation length is model dependent but the heterogeneity is rich in small scales (0.1-10km).

mft of the order of 10^3 s



Transfert radiatif et hétérogénéité du manteau inférieur

Précurseurs des PKIKP

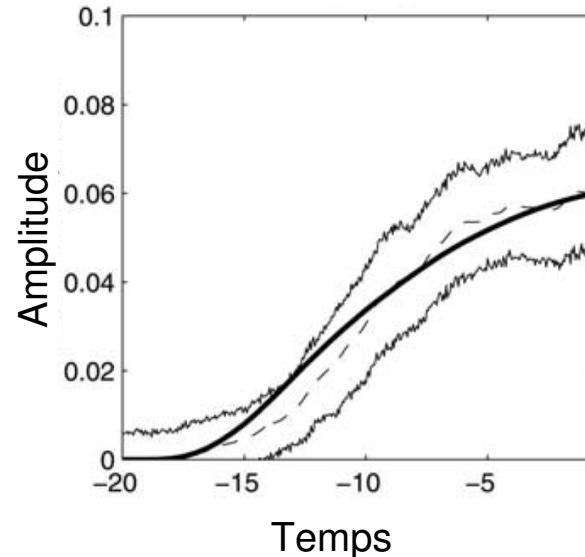


Margerin & Nolet, 2003

Observations vs Modèle

Précurseurs

PKIKP



- - - - - : Moy. des données ——— : Transfert radiatif

- Première application du **Transfert Radiatif** à l'échelle du globe
- Distribution d'hétérogénéité dans tout le manteau Inf
- Faibles fluctuations (0.1%) riches en petites échelles : 0.1km–10km, spectre k^{-3}